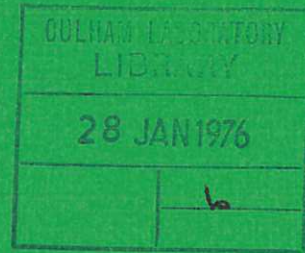


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Memorandum

ETNA
A FUEL-COOLANT INTERACTION CODE

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Abingdon Oxfordshire

1975

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ETNA
A FUEL-COOLANT INTERACTION CODE

D J Buchanan *

ABSTRACT

ETNA is a computer code for solving
the equations of a recently proposed
model for fuel-coolant interactions.

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February 1975

Title of Program : ETNA

Computer : ICL 4/70 Installation: Culham Laboratory

Operating System : ICL Multijob

Programming Languages used : Standard FORTRAN

High Speed store required : 19000 words

No. of bits in a word : 32

Overlay Structure : None

No. of Magnetic Tapes required : None

Other Peripherals used : Line printer

No. of cards in program deck : 2463

Card Punching Code : EBCDIC

Typical Running Time : On the ICL 4/70 at Culham Laboratory 5 cycles for
one value of the external pressure takes 10 seconds

Unusual features of the program : Standard FORTRAN is used except for the
NAMELIST facility and the abnormal entry statement,
ENTRY, into the subroutine ELIPSE (2.4)

1. INTRODUCTION

The computer code ETNA is a program for fuel-coolant interactions, FCI, calculations. The basic version of the code described here simulates the FCI model equations given by Buchanan (1974); however the structure of the code is sufficiently flexible to allow for substantial modifications and improvements in the basic equations. The structure of ETNA is based on the OLYMPUS package of Christiansen and Roberts (1974). MKS units are used throughout.

§ 2 contains a brief description of the physical model; § 3 is on numerical methods; and § 4 describes the structure of the code.

2. THE PHYSICAL MODEL

In this section we describe briefly the physics which the code ETNA simulates. A complete description has been given elsewhere (Buchanan 1974). The interaction is divided into five stages, the last four of which occur cyclically.

Stage 1. As a result of some triggering mechanism the liquids come into intimate contact and a vapour bubble is formed. This stage is solely a means of supplying the initial perturbation which precipitates the interaction. All our calculations start with the assumption that the initial bubble has been formed and thus the initial condition for the code is a bubble adjacent to the fuel surface.

Stage 2. The vapour bubble expands and then collapses as a result of condensation in the subcooled coolant. The collapse is asymmetric and a high velocity jet of liquid coolant, directed towards the fuel, is formed. The expansion phase is governed by the Rayleigh equation:

$$R \ddot{R} + \frac{3}{2} \dot{R}^2 = (P - P_0) / \rho_c \quad \dots (2.1)$$

R and P are the radius and pressure of the bubble at time t, measured from the start of the expansion, and P₀ and ρ_c are the ambient pressure and

coolant liquid density. If it is assumed that the expansion is adiabatic and that the value of γ is $4/3$ (the value for water vapour), then the maximum radius of the bubble is

$$R_m = R_i (1 + Z_r) \quad \dots (2.2)$$

where Z_r is the real solution of the cubic equation

$$Z^3 + 4Z^2 + 6Z - 3(P_i/P_o - 1) = 0 \quad \dots (2.3)$$

There is only one real solution of Eq (2.3). R_i and P_i are the initial radius and pressure of the bubble for the i^{th} cycle. Note the double role of the subscript i to denote 'initial' and ' i^{th} '. The time taken to expand to R_m is

$$t_g = R_i \left(\frac{3\rho_c}{2P_o} \right)^{\frac{1}{2}} \int_0^{Z_r} \frac{(1+z)^2 dz}{[3(P_i/P_o - 1)Z - Z^4 - 4Z^3 - 6Z^2]^{\frac{1}{2}}} \quad \dots (2.4)$$

When the maximum radius is reached the vapour condenses and the bubble collapses under the external pressure P_o . The equation of motion is:

$$R \ddot{R} + \frac{3}{2} \dot{R}^2 = - P_o / \rho_c \quad \dots (2.5)$$

and the time for collapse is

$$t_c = R_m \left(\frac{\rho_c}{3P_o} \right)^{\frac{1}{2}} \frac{\Gamma(5/6)\Gamma(1/2)}{\Gamma(4/3)} = 0.915 R_m \left(\frac{\rho_c}{P_o} \right)^{\frac{1}{2}} \quad \dots (2.6)$$

where Γ is the gamma function. This collapse results in the formation of a jet of liquid coolant. The results of Plesset and Chapman (1971) indicate that the velocity, length and diameter of the jet when it strikes the fuel surface are

$$\left. \begin{aligned} V_o &= V_c (P_o / \rho_c)^{\frac{1}{2}} \\ L_o &= L_c R_m \\ d_o &= d_c R_m \end{aligned} \right\} \dots (2.7)$$

The constants V_c , L_c and d_c are determined by the degree of departure from spherical symmetry. For bubble collapse adjacent to a solid wall Plesset and Chapman's results indicate that:

$$\left. \begin{aligned} V_c &= 13 \\ L_c &= 0.493 \\ d_c &= 0.237 \end{aligned} \right\} \dots (2.8)$$

and these values have been used in ETNA.

Stage 3. The jet of coolant (diameter d_o and length L_o) enters the fuel with velocity V_o . As the jet penetrates it disintegrates and mixes with the surrounding fuel. The surface area of contact between the fuel and the jet is given by:

$$A = A_o \exp(t/\tau) \dots (2.9)$$

$$\tau = \frac{11}{4} \left(\frac{\rho_f}{\rho_c} \right)^{\frac{1}{2}} \frac{d_o}{V_o} \dots (2.10)$$

A_o is the initial surface area of the jet and t is measured from the start of penetration. Whilst the jet remains liquid its volume is almost constant. Consequently, since the area increases exponentially the average distance S between material surfaces must decrease exponentially.

$$S = S_o \exp(-t/\tau) \dots (2.11)$$

S of course cannot decrease indefinitely and we denote the minimum value that is physically attainable by S_m .

Stage 4. As the jet penetrates, heat transfer occurs. We assume that heat transfer takes place one-dimensionally across each element of fuel-coolant (see Buchanan 1974, Fig 4.). The thickness of the element of coolant that is effectively heated, x , is given initially by

$$x = 4\sigma/t \dots (2.12)$$

When $4\sigma/t$ is equal to the thickness of the element, S , the coolant is heated throughout and the value of x is then given by Eqn (2.11). The temperature

of the heated region satisfies the equation (Buchanan 1973)

$$c \rho_c \times \frac{dT}{dt} = r_1 - r_2 T \quad \dots (2.13)$$

where r_1 is a parameter specifying the rate of heating of the jet per unit contact area due to heat transfer from the surrounding hot fluid; $r_2 T$ is inserted to take account of the fact that the heat transfer rate per unit contact area decreases as the temperature of the jet increases; and c is the specific heat of the coolant. Thus the temperature of the jet is given by:

i) When Eqn (2.12) applies

$$T = \frac{r_1}{r_2} - \left(\frac{r_1}{r_2} - T_0 \right) \exp\left(- \frac{r_2}{2\rho_c c \sigma} \sqrt{t} \right) \quad \dots (2.14)$$

ii) When Eqn (2.11) applies

$$T = \frac{r_1}{r_2} + F \exp\left\{ - \frac{r_2 \tau}{\rho_c c S_0} \exp(t/\tau) \right\} \quad \dots (2.15)$$

where F is a constant such that at time t' Eqns (2.14) and (2.15) give the same value for T . t' is the solution of

$$4\sigma\sqrt{t} = S_0 \exp(-t/\tau) \quad \dots (2.16)$$

iii) When S_m is attained

$$T = \frac{r_1}{r_2} - \left(\frac{r_1}{r_2} - T_1 \right) \exp\left(- \frac{r_2}{\rho_c c S_m} t \right) \quad \dots (2.17)$$

T_1 is the temperature given by Eqn (2.16) with $t = t''$ where

$$S_m = S_0 \exp(-t''/\tau) \quad \dots (2.18)$$

T_0 is the initial temperature of the coolant.

Stage 5. The jet continues to be heated until it is vaporized, which happens at either the saturation temperature, T_{sat} , if nucleation sites are available (heterogeneous nucleation), or at the homogeneous nucleation temperature, T_{hn} , if no nucleation sites are available. The fraction of the jet that is vaporized, β , depends on the process, being unity for heterogeneous nucleation and 0.33 for

homogeneous nucleation at 1 bar.

Since a finite mass of liquid is vaporized simultaneously and since the density cannot change instantaneously from ρ_l , the saturated liquid density, to ρ_v , the saturated vapour density, we assume that the vapour is initially formed at ρ_l and then expands adiabatically. The initial high pressure P_i is given by:

$$P_i \rho_l^{-\gamma} = P_v \rho_v^{-\gamma} \quad \dots (2.19)$$

where P_v is the vapour pressure. We also assume that the vapour bubble is spherical with radius R_i given by:

$$\frac{4}{3}\pi R_i^3 \rho_l = \beta M_j \quad \dots (2.20)$$

M_j is the mass of the jet. The pressure P_i in Eqn (2.19) is the same as the pressure P_i in Eqn (2.4).

The pressure in the liquid coolant at time t from vaporization and at a point r , measured from the centre of the bubble but outside the bubble, is given by:

$$P(r,t) - P_0 = [P(a, \tau') - P_0] a/r \quad \dots (2.21)$$

where

$$\tau' = t - \frac{r-a}{c_0} \quad \dots (2.22)$$

c_0 is the speed of sound in the liquid coolant while a is arbitrary ($a < r$) but is conveniently chosen to be the bubble's radius at time τ' .

From this model a number of results can be derived analytically which provide useful checks on the computer code. The ratio of the peak pressures at r due to successive cycles is:

$$\frac{P_i(r) - P_0}{P_{i-1}(r) - P_0} = \left(\frac{3}{16} \beta d_c^2 L_c \frac{\rho_c}{\rho_l} \right)^{1/3} (1 + Z_r) \quad \dots (2.23)$$

which for $P_0 = 1$ bar gives

$$\frac{P_i(r) - P_0}{P_{i-1}(r) - P_0} = \begin{cases} 6.673 & \text{(heterogeneous)} \\ 2.899 & \text{(homogeneous)} \end{cases} \quad \dots (2.24)$$

As P_0 is increased the ratio (2.23) decreases and becomes equal to unity when

$$P_0 = \begin{cases} 67.5 \text{ bar} & \text{(heterogeneous)} \\ 13.0 \text{ " } & \text{(homogeneous)} \end{cases} \quad \dots (2.25)$$

After N cycles the total mass of coolant heated during jet penetration is:

$$M_c = \frac{\pi}{4} \rho_c d_c^2 L_c \frac{\xi^N - 1}{\xi - 1} R_{m1}^3 \quad \dots (2.26)$$

where R_{m1} is the radius just prior to collapse of the initial perturbing bubble and

$$\xi = \frac{3}{16} \beta d_c^2 L_c \frac{\rho_c}{\rho_l} (1 + Z_r)^3 \quad \dots (2.27)$$

The ratio of the kinetic energy of the jet E_j just as it starts to penetrate, to the energy of the bubble E_b calculated as the work done by the bubble expanding to its maximum radius, is independent of cycle number and to order $O(1/P_i)$ is given by

$$\frac{E_j}{E_b} = \frac{3}{32} d_c^2 L_c V_c^2 = 0.439 \quad \dots (2.28)$$

3. NUMERICAL ANALYSIS

Most of the numerical work in the code is straight-forward function evaluation. The only parts that merit attention here are; 1) evaluation of Z_r , the real root of the cubic equation; 2) evaluation of the integral (2.4), and 3) the solution of Eqn (2.16).

1. Evaluation of Z_r : It is a trivial matter to show that Eqn (2.3) has only one real root. Cubic equations may be solved exactly by Cardan's method (Turnbull 1952) and this method has been used here. After the root is found the solution is checked for accuracy.

2. Evaluation of the integral (2.4): This is an elliptic integral and can, after some tedious algebra, be reduced to a combination of the three canonical forms; the elliptic integrals of the first, second and third kinds respectively (Whittaker and Watson 1927). However, these basic forms must then be evaluated numerically. In this case, it is just as adequate and easier to evaluate the integral numerically, the only minor problem being that the integrand diverges at both limits. This problem is easily overcome (Buchanan 1974) by observing that

$$\text{Integral (2.4)} = \int_0^{Z_r} F(Z) dZ + 2 \left(\frac{Z_r}{a} \right)^{\frac{1}{2}} + \frac{2(1 + Z_r)^2}{(3Z_r^2 + 8Z_r + 6)^{\frac{1}{2}}} \quad \dots (3.1)$$

where

$$F(Z) = \frac{(1 + Z)^2}{(aZ - Z^4 - 4Z^3 - 6Z^2)^{\frac{1}{2}}} - \frac{1}{(aZ)^{\frac{1}{2}}} - \frac{(1 + Z_r)^2 (Z_r - Z)^{-\frac{1}{2}}}{(3Z_r^3 + 8Z_r^2 + 6Z_r)^{\frac{1}{2}}} \quad \dots (3.2)$$

$$a = 3(P_i/P_o - 1)$$

$F(Z)$ is a perfectly well-behaved analytic function in the region $[0, Z_r]$. At $Z=0$ and $Z=Z_r$, $F(Z)$ has the values

$$\left. \begin{aligned} F(0) &= \frac{-(1 + Z_r)^2}{Z_r(3Z_r^2 + 8Z_r + 6)^{\frac{1}{2}}} \\ F(Z_r) &= -(aZ_r)^{-\frac{1}{2}} \end{aligned} \right\} \dots (3.3)$$

To effect the numerical integration of $F(Z)$ a Gaussian ten point formula is used (Abramowitz and Stegun 1965). Initially ten points are used over the whole range. The range is then divided into two equal parts and ten point integration used within each range. If the answers to the whole range and divided range integrations are not equal to within some specified limit, each range is further subdivided and ten point integration used again in each range. This process is continued until the required accuracy is achieved. Gaussian integration is performed over the interval $[-1, 1]$ and thus each integral must be between these limits. This is effected each time by a simple linear change of

variable. Near the limits 0 and Z_r the function $F(Z)$ is evaluated by a series expansion, otherwise its evaluation is straight forward.

To determine the pressure history at a point r the incomplete form of the integral (2.4) is required with some $Z \neq Z_r$ for the upper limit. This is easily done using the same routines as before with the upper limit changed.

3. Solution of Eqn (2.16). This is equivalent to finding the zeros of the function

$$G(Z) = 4\sigma\sqrt{Z} - S_0 e^{-Z/\tau} \quad \dots (3.4)$$

In fact it is obvious that this function has only one real zero for positive Z . The rule of false position is used (Hochstrasser 1962).

4. STRUCTURE OF THE COMPUTER CODE

The code ETNA consists of a main program plus subprograms. Of the subprograms, some are part of the ICL 4-70 Multijob Operating System and hence need not concern us, some are part of the OLYMPUS package (Roberts 1974, Christiansen and Roberts 1974) and the remainder are divided into five classes as described by Christiansen and Roberts (1974). The classes 1, 3, 4 and 5 almost identical to those of Christiansen and Roberts (1974). Table 1 contains a list of all subprograms in the classes 1-5 together with their purpose; Table 2 contains a list of variables in the labelled common blocks together with their purpose; and Fig 4.1 is a flow diagram of class 2 subprograms. Those subprograms marked with an asterisk are described in more detail in the text. Each of the subprograms of ETNA is decimally numbered according to OLYMPIAN conventions. References of the type <2.5> can be correlated with the list in Table 1. The source deck also contains this index. On the ICL 4-70 the corresponding file name is denoted by C2S5 (class 2, subprogram 5) followed by a type code defining the file eg FORTRAN source code F, object module Y.

CUBERT (2.3): This is a double precision subroutine which calculates the real root of the cubic equation. It calls the function DSIGN(X,Y) (2.16). DSIGN is a double precision function which transfers the sign of Y to the absolute value of X. The only reason for supplying this subprogram is that the system-supplied version does not work.

ELIPSE (2.4): This routine evaluates the elliptic integral. Normally the complete elliptic integral is evaluated; however the routine contains the statement ENTRY MOVE(TYM,MOT) which allows abnormal entry and the calculation of the incomplete elliptic integral.

HEAT (2.8): This is a rather lengthy routine whose control is as follows:

1. Decide which of the PDE and ODE methods is used to solve the heat conduction problem (Buchanan 1973). At present the ODE method is always used.
2. Solve equation (2.16) for t' .
3. Determine if Eqn (2.11) is always operable or if a specified minimum particle size S_m is to be included. If the latter is the case then Eqn (2.18) is solved for t'' .
4. Compare t' and t'' . If $t' > t''$ then Eqn (2.15) is never used. Equations (2.14) and (2.18) (possibly) are used to calculate the time required to heat to the vaporization temperature, TBRAK. The fraction of the jet vaporized is then calculated. If $t' < t''$ the solution may involve Eqn (2.15). If the temperature using Eqn (2.14) and t' is greater than TBRAK then Eqn (2.15) need not be used. If this is not the case Eqn (2.15) must be used and the constant F found. If Eqn (2.15) with $t=t''$ does not give a temperature greater than TBRAK then Eqn (2.18) must also be used. The fraction of the jet that is vaporized is found.
5. The pressure and radius of the new bubble is calculated.

- SOLRTT (2.17): This routine solves Eqn (2.14) for t with the temperature T set equal to TBRAK.
- HISTOR (2.17): This routine calculates the pressure as a function of time at the point r using Eqn (2.21). As this equation stands t is the independent variable. The pressure $P(r,t)$ is related to the pressure in the bubble at the retarded time τ' when the bubble radius is a . To avoid having to iterate to find consistent values of a and τ' it is much more convenient to regard t as the dependent variable. So an a is chosen and the time for the bubble to expand to radius a is found by evaluating the incomplete elliptic integral. This involves the abnormal entry into ELIPSE (2.4). Since a and τ' are known t can be found from Eqn (2.22) and of course $P(a,\tau')$ is just the pressure of the bubble when it has a radius a .
- SOLEXP (2.19): This routine solves Eqn (2.15) for t with T set equal to TBRAK. The constant F is such that Eqn (2.14) and (2.15) give the same value for T when $t = t'$.
- INFORM (5.1): This is a diagnostic routine which prints out selected common variables if any errors occur in the calculation. A modified version of the system DEBUG also calls this routine so that if DEBUG is called, the current state of the common variables is available to the user.

The variable TBRAK is used solely to calculate the time required for vaporization. The saturated density values are read as data. The default option of TBRAK is 373.15 corresponding to heterogeneous nucleation at 1 bar. The time to heat to TBRAK takes no account of the additional time required to supply the latent heat; however, an estimate of the time to heat to the saturation temperature plus the time to supply the latent can be obtained by using

the equivalent temperature approximation, ie $TBRAK=373.15 + L/C$ where L is the latent heat and C the specific heat. This modified value of $TBRAK$ must of course be read in as data.

The variable $GASDEN$ is the density of the vapour when it is formed. By assumption this density is the saturated liquid density corresponding to the particular vaporization temperature used (cf Eqn 2.19). The variables $RATE1$ and $RATE2$ must satisfy the condition $RATE1/RATE2 < TBRAK$.

As far as possible the program has been tested against analytical solutions. In this respect the expressions 2.23-2.28 are most useful. The only calculation that cannot be tested against an analytical formula is the total elapsed time. Obviously time solutions of Eqns (2.14)-(2.18) are easily checked; the only part that is not easily checked is the evaluation of the elliptic integral. The Gaussian integration scheme has been checked against a different integral which can be evaluated analytically.

The default option for all common variables is given in $PRESET \langle 1.3 \rangle$. These options can be overwritten by using the $NAMelist$ facility in $DATA \langle 1.4 \rangle$. All variables supplied by the user must be in MKS units. The default values will give an FCI calculation for tin and water with vaporization by heterogeneous nucleation. The calculation is started by an initial bubble collapsing and the radius of this bubble is specified as $RINIT$. A cycle is defined to be that period between successive formation of two bubbles, so that a cycle ends when a bubble is about to expand. The expansion time of the first cycle is of course zero since the calculation is started with bubble collapse.

Output from two test runs is shown in Appendices 1 and 2. The first run is the default run (Appendix 1) when no data is supplied. This corresponds to an FCI at an external pressure of 1 bar with vaporization by heterogeneous nucleation. The second run (Appendix 2) corresponds to an

FCI at 5 bar with homogeneous nucleation. The variables reset in DATA <1.4> are shown below:

```
PRESS= 5.0E 05  
PINIT= 3.614E 08  
DELP= 5.0E 05  
GASDEN=728.2  
TBRAK= 565.  
FRAC1= 0.2547
```

Some of the relevant thermodynamic properties for water can be obtained from UK Steam Tables (1970), others such as T_{hn} must be calculated by the user from one of the theories of homogeneous nucleation (Frenkel 1946 and Cooper 1952, provide useful reviews).

14 February 1975
dg

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TABLE 1 - INDEX OF SUBPROGRAMS

NB. Subprograms marked * are described in more detail in the text

Name No Dummy Arguments Purpose

Class 1 Prologue

LABRUN	1.1		Label the run
CLEAR	1.2		Clear the common blocks
PRESET	1.3		Define default values
DATA	1.4		Define data specific to the run
AUXVAL	1.5		Define auxiliary variables
INITAL	1.6		Define initial conditions
RESUME	1.7		Restart the calculation
START	1.8		Start the calculation

Class 2 Calculation

STEPON	2.1		Advance the calculation by one cycle
BUBBLE	2.2		Solve the bubble dynamics problem
*CUBERT	2.3		Find the real root of the cubic equation
*ELIPSE	2.4		Evaluate the elliptic integral
GAUSS	2.5	ANS,RLIM1,RLIM2,D E,C,MOTION	10 Point Gaussian integration scheme
FUNC	2.6	Y1,D,E,C,MOTION	Evaluate the integrand
JET	2.7		Evaluate the jet characteristics
*HEAT	2.8		Solve the heat transfer problem
THICK	2.9		Solve the thickness equation
YTHICK	2.10	XX	Evaluate the thickness function
TSQRTT	2.11	T	Find temperature on the assumption that the square root approximation is valid.
*SQLRTT	2.12		Solution of square root approximation
ARIA	2.13	T	Evaluates the area at time T.
KONST	2.14		Integration constant in exponential mass approximation.
TEMEXP	2.15	T	Temperature at time T according to the exponential mass approximation.
DSIGN	2.16	XX,YY	Transfer of sign.
*HISTOR	2.17		Calculate pressure at point AR.
PTIME	2.18	AA,PTT	Calculate retarded time
*SOLEXP	2.19		Solution of exponential approximation.

Class 3 Output

OUTPUT	3.1	K	Output routine
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Class 4 Epilogue

TESEND	4.1		Test for completion of run
ENDRUN	4.2		Terminate the run.

Class 5 Diagnostics

*INFORM	5.1		Diagnostic routine.
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The terms 'square root' or 'square root mass' approximation and 'exponential' or 'exponential mass' approximation above and in the program listing refer to Eqns (2.14) and (2.15) respectively.

Variable Default Value Purpose

Common Block FC11.

RFINAL	0.	Final radius of bubble after expansion, R_m
RINIT	1.E-04	Initial radius of bubble at start of expansion, R_i This variable is used to start the calculation.
ZCUBE	0.	Real root of cubic equation, Z_r
TCOLL	0.	Collapse time of bubble, t_c
RHO	1000.	Density of bulk coolant, ρ_c
PRESS	1.01325E 05	Ambient pressure, P_0
TIME	0.	Total time for expansion and collapse of bubble, $t_c + t_g$
TXPAND	0.	Time for bubble to expand, t_g
PINIT	1.90083E 09	Pressure at which bubble is formed, P_i
X(I)I=1,5	0.9739065 0.8650634 0.6794096 0.4333954 0.1488743	Nodes for Gaussian ten point integration
W(I)I=1,5	0.0666713 0.1494513 0.2190864 0.2692667 0.2955242	Weights for Gaussian ten point integration
VJET	0.	Velocity of jet, V_0
RLJET	0.	Length of jet, L_0
DJET	0.	Diameter of jet, d_0
RJETV	13.	Proportionality constant for jet velocity, V_c
RJETL	4.78/9.70	Proportionality constant for jet length, L_c
RJETD	2.30/9.70	Proportionality constant for jet diameter, d_c
DELP	1.01325E 05	Pressure difference under which the bubble collapses, P_0
TAU	0.	Time constant for area increase, τ .
RJETTA	11./4.	Proportionality constant for area increase, $\frac{11}{4}(\rho_f/\rho_c)^{\frac{1}{2}}$
AREAO	0.	Initial area of jet, A_0
TCRIT	647.3	Critical temperature of coolant (not used)
TIMHET	0.	Time to heat to vaporization temperature
TWID	0.	Time solution of Eqn (2.16), t'
TEXP	0.	Temperature from exponential approximation, Eqn (2.15)
TEMP	0.	Temperature from square root approximation, Eqn (2.14)
GAMMA	21.E-05	Coefficient of thermal expansion of coolant (not used)
RKAPPA	4.45 E-10	Compressibility of coolant (not used)
TZERO	300.	Ambient coolant temperature, T_0 .
GASDEN	958.13	Density of vapour at temperature and pressure at which vapour is formed (<u>ie</u> the saturated liquid density, see (Eqn 2.19))

SIGMA	3.879126E-04	Thermal diffusivity of coolant, σ
FACTOR	0.5	Proportionality constant for Eqn (2.16). This constant allows for the fact that S_0 in Eqn (2.11) may not equal the diameter of the jet, d_0 . For example, in Buchanan (1974) Fig 4, $S_0=d_0/2$ seems more appropriate.
RATE1	1.E 06	Rate of heating, r_1
RATE2	1.E 03	Rate of cooling, r_2
SPHT	4200.	Specific heat of coolant, c
AREA	0.	Area of jet at time t , A .
CONST	0.	Integration constant in Eqn (2.15), F .
TBRAK	373.15	Temperature at which the jet is vaporized. For example, T_{sat} or T_{hn} .
GAM	4./3.	Ratio of specific heats of coolant vapour, γ .
ENBUB	0.	Energy of bubble, E_b
ENJET	0.	Energy of jet, E_j
ELTIME	0.	Total elapsed time.
RELENTH	1.	Factor giving thickness in square root approximation, Eqn (2.14). The heated mass is defined somewhat arbitrarily to be $2\sigma/t$. This constant allows any multiple of $2\sigma/t$ to be used.
ACC1	1.E-05	Accuracy of solution of cubic equation
ACC2	1.E-05	Accuracy of evaluation of integral
ACC3	1.E-08	Accuracy of solution of Eqn (2.16)
SRI	0.	Saves previous value of RINIT
SPARE (I) I=1,10		Spare vector used as follows:
SPARE (1)	0.	Controls number of runs
SPARE (2)	2.	If GE 0.5 then Eqns (2.17) and (2.18) are included in calculation.
SPARE (3)	3.E-07	Minimum particle size for cut-off, S_m .
SPARE (4)	0.	Thickness of jet at time of vaporization
SPARE (5)	0.	Not used
SPARE (6)	0.	Not used
SPARE (7)	0.	Stores ZCUBE
SPARE (8)	0.	Stores RINIT
SPARE (10)	0.	Total mass of coolant, M_c
NFREQ	1	Frequency of output
NMAX	5	Maximum number of cycles
NSCAN	4	Not used
METHOD	0	Switch to determine which of the ODE and PDE methods are used. Only the ODE method is used in the current version.

Common Block FCI2

PHIST(I) I=1,10	0.	Pressure at point r as a function of time, P(r,t)
PT(I) I=1,10	0.	Times at which P(r,t) is calculated
PTIME(I) I=1,10	0.	Total times at which P(r,t) is calculated
AR	0.1	Point where pressure is calculated, r .
ZMAX	0.	Upper limit of incomplete elliptic integral
SONCOL	1500.	Velocity of sound in coolant, c_0
FRAC1	1.	Maximum fraction of jet that can be vaporized, β
RIMP	0.	Impulse at point r .

Common Block FCI3

FREE(I) I=1,20		Spare vector used as follows
FREE(1)	1.	Factor modifying RJETV
FREE(2)	1.	Factor modifying RJETL
FREE(3)	1.	Factor modifying RJETD. These constants are intended for future use to take account of possible modifications to Eqn (2.8) as a result of non-condensable gas, etc.
FREE(4)	7000.	Density of fuel, ρ_f
FREE(I) I=5,20	0.	Not used.

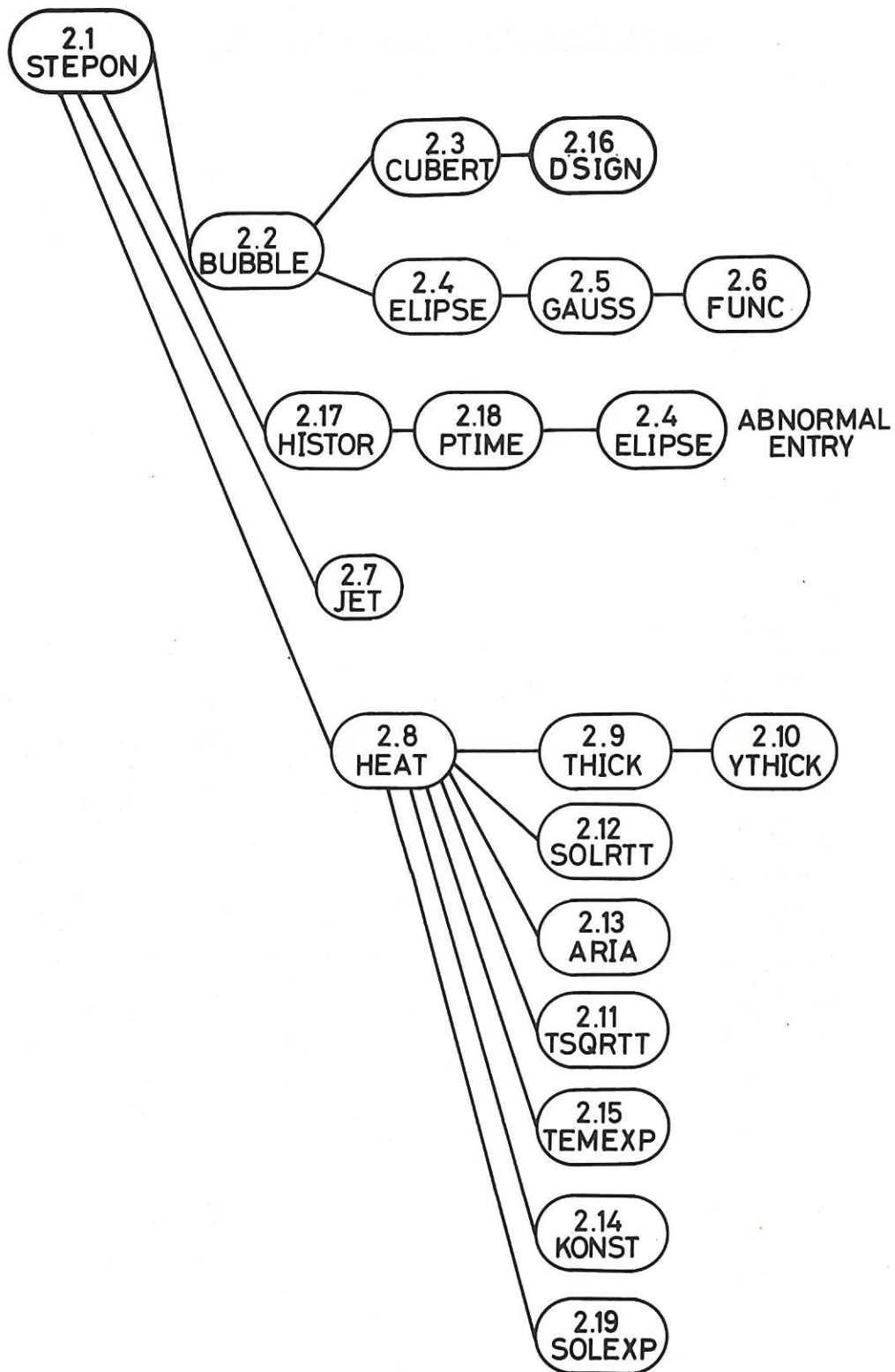
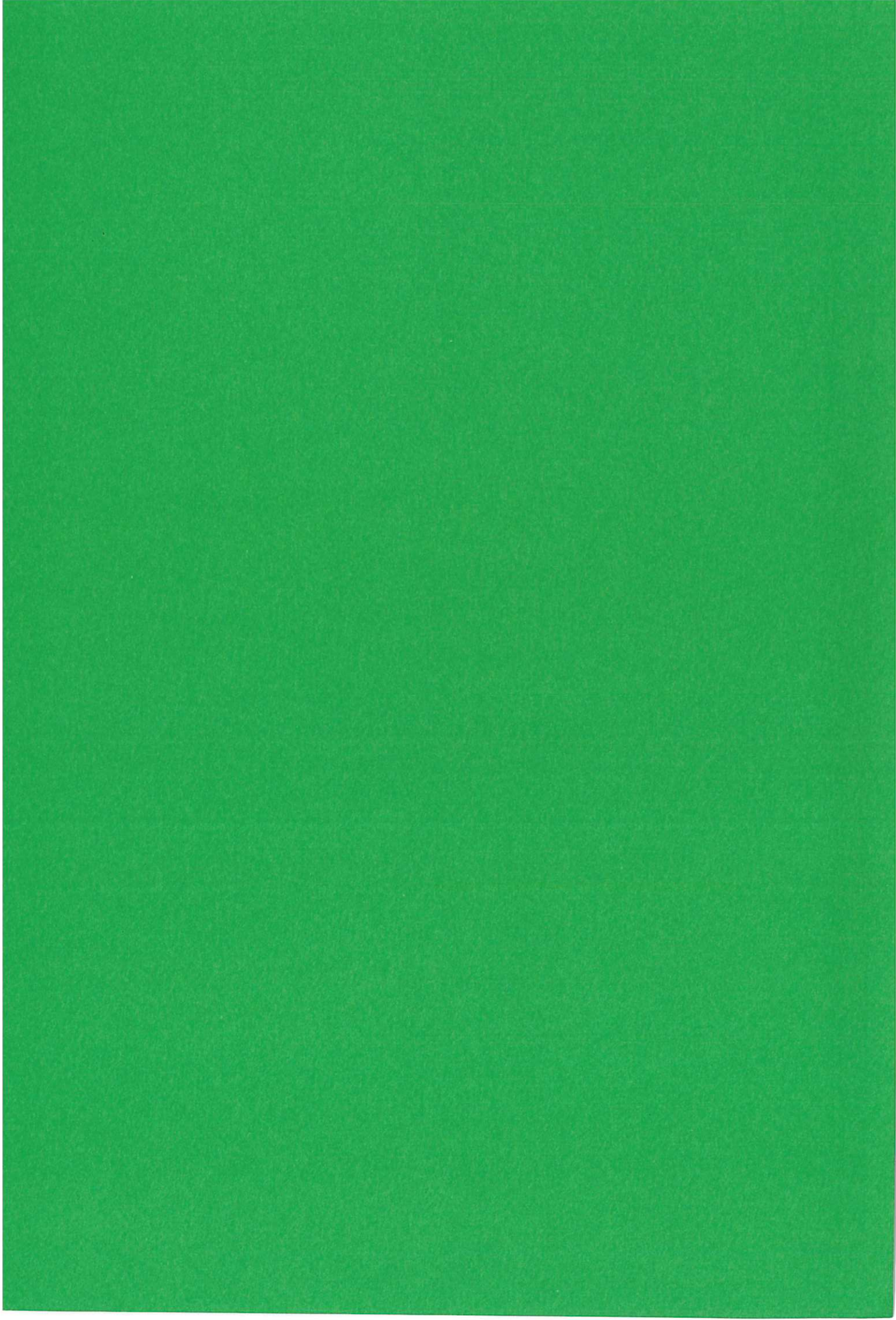


Fig.4.1

Flow Diagram of Class Two Subprograms



The first part of the document discusses the importance of maintaining accurate records of all transactions. It emphasizes that every entry should be supported by a valid receipt or invoice. This not only helps in tracking expenses but also ensures compliance with tax regulations.

In the second section, the author provides a detailed breakdown of the company's revenue streams. This includes sales from various product lines and services. The data shows a steady increase in revenue over the past year, which is attributed to strategic marketing efforts and improved operational efficiency.

The third section focuses on the company's financial health and liquidity. It highlights the strong cash flow and the ability to meet all financial obligations. The author notes that the company's debt-to-equity ratio remains low, indicating a solid financial foundation.

Finally, the document concludes with a summary of the overall performance and a look ahead at future goals. The author expresses confidence in the company's ability to continue its growth trajectory and achieve its long-term objectives.