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A Simple Theory of Relativistic Damping of Electromagnetic Cyclotron and Electron Bernstein Waves in Tokamaks

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Abstract. An approximation to the relativistic broadening of an electron cyclotron resonance is obtained by including only the effect of the momentum parallel to the magnetic field. This approximation is then applied to the calculation of the absorption of the electromagnetic X-mode and compared to the full relativistic analysis with which it agrees very well. The propagation and damping of electron Bernstein waves is calculated using the approximate electrostatic dispersion relation and compared with the result obtained from the full electromagnetic dispersion relation.

INTRODUCTION

Electron cyclotron waves are important for heating and current drive in tokamaks. Spherical tokamaks offer a special challenge to wave heating methods and electron Bernstein waves appear to be a promising candidate [1]. The damping of electron Bernstein waves is very strong and for heating and current drive applications in tokamaks the location of the absorption region is very important. The width of the electron cyclotron resonance is due to the Doppler effect and to the relativistic dependence of mass on energy. The relativistic resonance condition is

$$\omega - k_{\parallel} v_{\parallel} - n\Omega_e \left(1 - \frac{v_{\parallel}^2}{c^2} - \frac{v_{\perp}^2}{c^2} \right)^{1/2} = 0 \quad (1)$$

where Ω_e is the electron cyclotron frequency calculated with the rest mass and $\Omega_e > 0$, ω , k_{\parallel} , v_{\parallel} , n , v_{\perp} and c are the wave frequency, the wave number parallel to the magnetic field, the velocity parallel to the magnetic field, the cyclotron harmonic number, the velocity perpendicular to the magnetic field and the velocity of light respectively. For weakly relativistic conditions, $v_{\parallel}^2 \ll c^2$, $v_{\perp}^2 \ll c^2$, the resonant values of v_{\parallel}/c given by Eq (1) are

$$\frac{v_{\parallel}}{c} = \frac{n_{\parallel}\omega}{n\Omega_e} \pm \frac{n_{\parallel}\omega}{n\Omega_e} \left[1 - \frac{2(\omega/n\Omega_e - 1)n^2\Omega_e^2}{n_{\parallel}^2\omega^2} - \frac{n^2\Omega_e^2 v_{\perp}^2}{n_{\parallel}^2\omega^2 c^2} \right]^{1/2} \quad (2)$$

The complications of the relativistic theory arise because of the dependence of v_{\parallel} on v_{\perp} . The theory is greatly simplified if v_{\perp}/c is neglected in Eq (2) giving

$$\frac{v_{\parallel}}{c} \cong \frac{n_{\parallel}\omega}{n\Omega_e} \left[1 \pm \left(1 + \frac{2Xn^2\Omega_e^2}{n_{\parallel}^2\omega^2} \right)^{1/2} \right] \quad (3)$$

where $X = 1 - \omega/n\Omega_e$. The approximation given by Eq (3) is compared with the exact result of Eq (2) in Figures 1(a) and 1(b) where the resonant value of v_{\parallel}/c is plotted against the parameter X which represents the spatial co-ordinate across the resonance. The full line is the approximate value from Eq (3) and the dotted and dashed dotted lines correspond to Eq (2) where v_{\perp}^2 corresponds to 5keV and 10keV respectively. The non-relativistic result, the straight broken line, is shown for comparison. The approximation given in Eq (3) agrees well and improves further from the cold resonance.

RELATIVISTIC DAMPING

The above approximation to the resonant values of v_{\parallel} can now be used to obtain a much simpler dispersion relation. This is illustrated by obtaining an approximation to the weakly relativistic expression for the xx component of the dielectric tensor, ϵ_{xx} . For weakly relativistic conditions the main effect of the relativistic correction occurs in the resonance condition. By neglecting the correction due to v_{\perp}/c the v_{\perp} -integration in the calculation of ϵ_{xx} is the same as in the non-relativistic case. Since there are now two values of the resonant v_{\parallel} , ϵ_{xx} is approximated by

$$\epsilon_{xx} = 1 + \sum_{n=-\infty}^{\infty} \frac{\omega_{pe}^2}{\omega^2} \frac{n^2 e^{-\lambda^2}}{\lambda} I_n(\lambda) \zeta_0 \left[\text{Re} Z(\zeta_n) + i\pi^{1/2} (e^{-v_1^2/v_{Te}^2} + e^{-v_2^2/v_{Te}^2}) \right] \quad (4)$$

where $v_{1,2}$ are the two roots of Eq (3), $\lambda = k_{\perp}^2 v_T^2 / 2\Omega_e^2$, $\zeta_0 = \omega / k_{\parallel} v_{Te}$, $\zeta_n = (\omega - n\Omega_e) / k_{\parallel} v_{Te}$ and $Z(x)$ is the plasma dispersion function. The other elements of $\underline{\epsilon}$ can be approximated in a similar manner. This approximation to the weakly relativistic theory is now used to calculate $\text{Im } n_{\perp}$ for the X-mode at the second harmonic resonance and the results are compared with the corresponding calculation of Airoldi-Crescentini et al [2] who used the standard relativistic theory. The comparison is illustrated in Figures 2(a) and 2(b) in which the variation of $\text{Im } n_{\perp}$ across the resonance region is shown for different values of n_{\parallel} . It can be seen that the simple

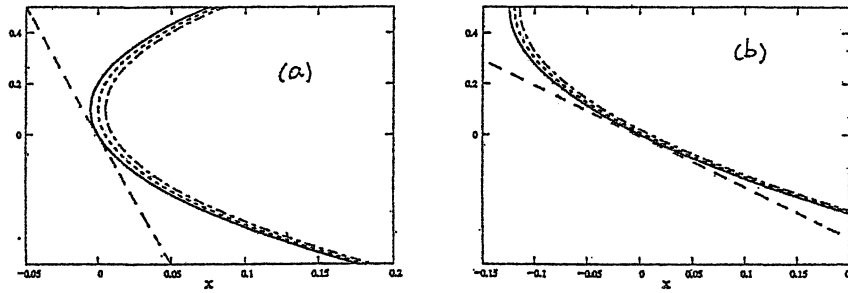


FIGURE 1 The resonant values of $v_{||}/c$ plotted as a function of $X = 1 - \omega/n\Omega_e$. The solid line corresponds to Eq (3) ($v_{\perp}^2 = 0$), the dotted and dashed dotted lines to Eq (2) for values of $v_{\perp}^2 = 5\text{keV}$ and 10keV respectively. The straight dashed line is the non-relativistic value (a) $n_{||} = 0.1$, (b) $n_{||} = 0.5$.

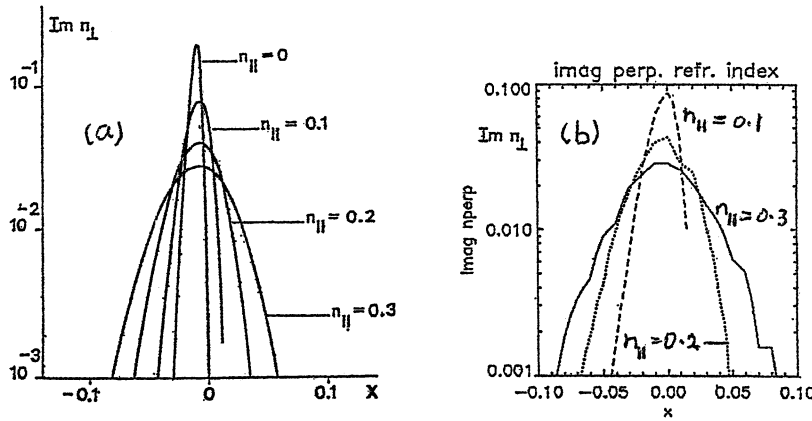


FIGURE 2 $\text{Im } n_{\perp}$ plotted as a function of distance across the resonance region for the X-mode at $\omega = 2\Omega_e$ with $\omega_{pe}^2/\omega^2 = 0.2$, $T_e = 0.42\text{keV}$ and for various values of $n_{||}$, (a) relativistic results of Airoidi-Crescentini et al (1980), (b) relativistic approximation.

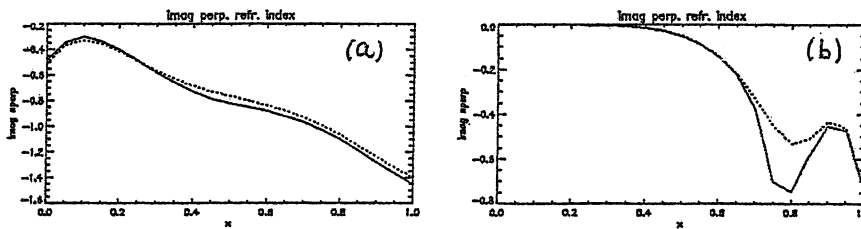


FIGURE 3 $\text{Im } n_{\perp}$ as a function of Ω_e/ω which varies linearly with x from 0.3 to 0.25 for electron Bernstein waves calculated from electromagnetic dispersion relation (solid line) and from $\epsilon_{xx} - n_{\perp}^2 = 0$ (dotted line) both using the relativistic approximation (a) $n_{||} = 0.5$, (b) $n_{||} = 0.1$ and $\omega_{pe}^2/\omega^2 = 5$ and $T_e = 5\text{keV}$ in both cases.

approximation to the relativistic damping reproduces the results of the more exact theory very well, both qualitatively and quantitatively.

The approximate relativistic theory is now used to calculate the damping of electron Bernstein waves. For this case we compare the values of $\text{Re } n_{\perp}$ and $\text{Im } n_{\perp}$ obtained from the full electromagnetic theory with those obtained from the approximate dispersion relation

$$\varepsilon_{xx} - n_{\parallel} = 0 \quad (5)$$

In both cases, the relativistic damping is represented by the approximation in which the v_{\perp}^2/c^2 correction is neglected. The real parts of n_{\perp} agree very closely and the imaginary parts are shown in Figure 3(a) for $n_{\parallel} = 0.5$ and in Figure 3(b) for $n_{\parallel} = 0.1$. The solid line corresponds to the full electromagnetic dispersion relation and the broken line is obtained from the solution of Eq (5). The values for $\text{Im } n_{\perp}$ agree well for $n_{\parallel} = 0.5$ but not so well for $n_{\parallel} = 0.1$.

CONCLUSIONS

A simple approximate method of calculating relativistic electron cyclotron damping has been proposed. It is demonstrated that the approximation agrees well with the standard, weakly relativistic analysis for the second harmonic X-mode even for $n_{\parallel} = 0.1$. This relativistic approximation is then used to calculate the real and imaginary parts of n_{\perp} for electron Bernstein waves using both the full electromagnetic and approximate electrostatic ($\varepsilon_{xx} - n_{\parallel}^2 = 0$) dispersion relations. The simplified electron Bernstein dispersion relation agrees closely with the electromagnetic dispersion relation. The development of a simple approximation to the relativistic damping and propagation of electron Bernstein waves should be useful for current drive and ray tracing calculations. It could also be useful for similar calculations for the O- and X-modes for conventional ECRH.

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