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Impact of nitrogen seeding on confinement and power load control of a high-triangularity JET ELMy H-mode plasma with a metal wall

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Abstract

This paper reports the impact on confinement and power load of the high-shape 2.5 MA ELMy H-mode scenario at JET of a change from all carbon plasma-facing components to an all metal wall. In preparation to this change, systematic studies of power load reduction and impact on confinement as a result of fuelling in combination with nitrogen seeding were carried out in JET-C and are compared with their counterpart in JET with a metallic wall. An unexpected and significant change is reported on the decrease in the pedestal confinement but is partially recovered with the injection of nitrogen.

(Some figures may appear in colour only in the online journal)

^a See the appendix of Romanelli F. *et al* 2012 24th IAEA Fusion Energy Conf. on Fusion Energy 2012 (San Diego, CA, 2012) (Vienna: IAEA) www.naweb.iaea.org/naweb/physics/FEC/FEC2012/papers/197_OV13.pdf.

1. Introduction

The reference scenario for achieving $Q = 10$ in ITER is the type-I ELMy H-mode at 15 MA, with a deuterium–tritium mixture, sufficient energy normalized confinement ($H_{98(y,2)}$), at high density ($H_{98} \sim 1.0$, Greenwald fraction $f_{\text{GDL}} \sim 0.85$) and compatible with its plasma-facing components (PFCs). The expected power to the divertor makes high divertor radiation mandatory and requires the injection of extrinsic impurities. The challenge of achieving the scenario requirements with the material selection of the DT phase of ITER is being addressed at JET with its new wall (JET-ILW) including main-chamber limiters made of bulk beryllium (Be) and a full tungsten (W) divertor [1, 2]. Three aspects require investigation: first the capacity to reduce the inter-ELM power load to the W divertor without significant degradation of the energy confinement; second the limitation of impurity production from the new PFCs to maintain performance, in terms of dilution (Be) and core radiation (W), the latter by reducing the energy of impurity ions in the divertor, which governs the W sputtering yield [3]; third the compatibility of the extrinsic impurities with the ILW materials in particular in the case of a species such as N_2 .

The choice of an extrinsic impurity for ITER is not only motivated by its impact on the plasma confinement with a high divertor radiation but also by its interaction with the PFCs and its consequences for the tritium processing plant. Being chemically reactive, N_2 can interact with the PFCs to form nitrides but the more important issue is its chemical reaction with hydrogen isotopes to form ammonia (ND_3) [4]. The produced amount of ammonia is potentially considerable and would have implications for the design of the ITER tritium plant. The ITER team plans to have the flexibility to use Ne, Ar and N_2 (or a mixture of) as seed gases, but it remains the case that chemically reactive species will be more of a safety challenge for the plant than a noble gas. Nevertheless, nitrogen as an extrinsic radiator has a role to play in the development of radiative scenario in current devices for several reasons: it plays the role that neon will have under ITER divertor conditions where the pedestal will be hotter $\sim 3\text{--}4$ keV [5]; it also isolates the effect that high divertor radiation for power load reduction has on an integrated scenario. At JET, Ne and Ar also radiate in the pedestal and main plasma (respectively) adding an additional degree of difficulty in achieving stationary conditions due to their ability to reduce the power entering the edge region and as a result reduce the pedestal confinement [6, 7].

This paper investigates the potential of N_2 seeding for achieving a radiative divertor for power load control at JET with the new metal wall. It focuses on the ELMy H-mode scenario at high triangularity (2.7 T/2.5 MA, $q_{95} \sim 3.5$, $P_{\text{in}} \sim 16$ MW, $\delta \sim 0.4$). The high-triangularity scenario was chosen as it had the following characteristics in JET-C: its global confinement enhancement factor $H_{98(y,2)}$ is higher than the low-triangularity scenario at the required Greenwald fraction for ITER of $f_{\text{GDL}} \sim 0.85$; it can be fuelled up to the Greenwald density ($f_{\text{GDL}} \sim 1.0$) without confinement loss ($H_{98} \sim 1.0$) at JET, the so-called type-I/III ELM regime [8].

This paper is organized as follows: a summary of the seeded experiments in JET with carbon-fibre composite PFCs

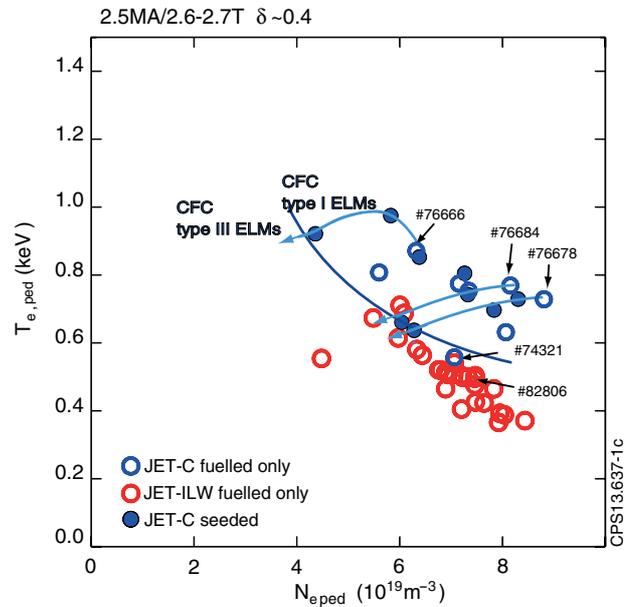


Figure 1. Pedestal n_e – T_e diagram for high- δ ELMy H-mode discharges for JET-C and JET-ILW discharges with 2.5 MA and 2.6–2.7 T: each symbol corresponds to a discharge. JET-C discharges are in blue symbols, open if only fuelled and filled with the addition of N_2 seeding. The dark blue line indicates the boundary in n_e – T_e between type-I and type-III ELMy H-modes obtained experimentally, as explained in [6]. Discharges with increasing D_2 -fuelling rate of 0.3×10^{22} el s^{-1} , 1.8×10^{22} el s^{-1} , 2×10^{22} el s^{-1} , 7×10^{22} el s^{-1} and 6×10^{22} el s^{-1} are indicated with, respectively, #76666, #76684, #76678 and #74321. The light blue arrows indicate the trajectory of pedestal parameters of discharges at the same D_2 -fuelling and increased N_2 -seeding rate from JET-C discharges [6].

(JET-C) is given in section 2; differences in confinement between JET-C and JET-ILW are reported in section 3; the recovery of confinement with N_2 seeding is presented in section 4; the inter-ELM power load reduction resulting from N_2 seeding is shown in section 5; the lack of stationarity of the plasma discharges and the averaged erosion yield are discussed in section 6; the lack of improved confinement with N_2 seeding in the low-triangularity plasma is briefly discussed in section 7 together with the possible effect of dilution in the confinement improvement; a concluding summary is given in section 8.

2. Description of experiments and main results from JET-C seeding studies

Experiments carried out in JET-C provide the necessary set of discharges to document the change of wall and its impact on power load to the divertor. This reference set of discharges explored the reduction of the inter-ELM power load in an ELMy H-mode scenario (2.7 T/2.5 MA, $q_{95} \sim 3.5$, $P_{\text{in}} \sim 16$ MW, $\delta \sim 0.4$, $H_{98(y,2)} \sim 1.0$, $f_{\text{rad}} \sim 0.45$) at a high density ($f_{\text{GDL}} \sim 0.8$) with a fuelling mix of D_2 and either Ne or N_2 as seed impurity, constant input power and proximity to the transition from type-I to type-III ELMs [6, 9], see figure 1. A deuterium fuelling scan from 0.3 to 6.0×10^{22} el s^{-1} was performed during part of which the ELM regime changed from type I, with a frequency of 20–25 Hz (at the lowest fuelling

level), to the so-called type I/II with an ELM frequency of 7–10 Hz [6].

It was shown that with either N₂ or Ne, conditions of partial divertor detachment were achieved with less than 10% energy confinement reduction in the main plasma with respect to the low-fuelled reference case ($\sim 0.3 \times 10^{22} \text{ el s}^{-1}$, #76666). A difference in the radiation localization was observed between the two seeding gases. N₂ seeding increased the divertor radiation while Ne seeding decreased it. Although reduced with the injection of extrinsic impurity, C concentration remained above 1% and contributed to the divertor radiation in addition to the extrinsic impurity [6]. Also it was observed that the H-mode pedestal at the plasma edge of Ne-seeded discharges often crossed the boundary between type-I and type-III ELM regimes, leading to compound ELMs and unsteady edge conditions. No condition was observed in this configuration for which N₂ seeding increased the plasma confinement, as was observed in ASDEX Upgrade with a full W wall [10].

The transition from type-I to type-III ELM regime turned out to be the limitation in achieving a high value of divertor over main plasma radiation with an H -factor close to 1.0. The main plasma, $P_{\text{rad}}^{\text{bulk}}$, and divertor, $P_{\text{rad}}^{\text{div}}$, radiated power were determined from tomographic reconstruction by integrating the 2D reconstruction for the z -abscissa value greater than -1.2 m for $P_{\text{rad}}^{\text{bulk}}$ and for $P_{\text{rad}}^{\text{div}}$ with z less than -1.2 m , as described in [6]. Figure 1 shows that when N₂ is seeded in JET-C, at a rate from 0.0 to $2.4 \times 10^{22} \text{ el s}^{-1}$, in the two plasmas with the highest density (#76684 and #76678), the pedestal density decreases followed by a decrease in both temperature and density at the highest seeding rate of $\sim 4 \times 10^{22} \text{ el s}^{-1}$ for which a transition from type-I to type-III ELM regime occurred [6].

3. Differences in plasma confinement between JET-C and JET-ILW high-shape ELMY H-mode discharges at 2.5 MA

Similar discharges were repeated in JET-ILW with practically identical high-shape (2.5 MA/2.7 T) with the aim to characterize the difference between JET-C and JET-ILW. Studies have shown that C concentration has dropped in the divertor and the main plasma by a factor of 10 [11]. Although dedicated studies address specific issue related to a change of material (such as divertor detachment [12], W erosion [13, 14], impurity composition [15]), the experiments presented here integrate all the aspects and provide a comparison with JET-C of the overall impact of the new PFCs on the chosen plasma H-mode scenario. As expected, W contamination in the main plasma makes operation at low fuelling difficult and high-shape ELMY H-mode at 2.5 MA is restricted to fuelling levels higher than $0.9 \times 10^{22} \text{ el s}^{-1}$ in contrast to the lowest fuelling of $\sim 0.3 \times 10^{22} \text{ el s}^{-1}$ accessible in JET-C related studies. In any case, studies made in JET-C in preparation for the change of PFCs provide a set of reference discharges for higher fuelling levels (up to $6 \times 10^{22} \text{ el s}^{-1}$).

A comparison of two plasmas with similar starting parameters between JET-C (#76678) and JET-ILW (#82806) plasmas at 2.5 MA at a D₂-fuelling rate of $\sim 2.7 \times 10^{22} \text{ el s}^{-1}$ is shown in figure 2. At similar average electron density, and similar input power, the stored energy of the JET-ILW

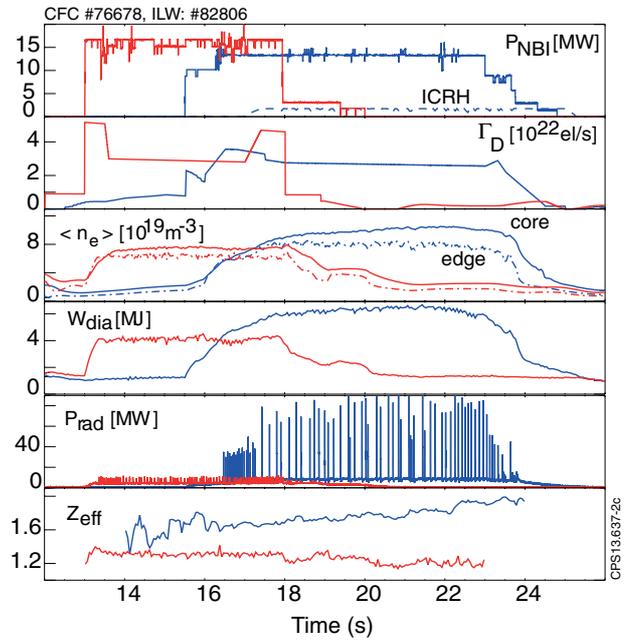


Figure 2. Time traces of unseeded ELMY H-mode discharges in JET-C wall #76678 (in blue), and JET-ILW #82806 (in red): (from top to bottom) neutral beam injection (NBI) heating and ion cyclotron resonance heating (ICRH) (dashed) power, D₂-fuelling waveform, line-integrated core and edge (dashed) electron densities measured by an interferometer, diamagnetic stored energy, total radiated power and Z_{eff} .

discharge is reduced by 40% in comparison with JET-C with a stored energy of 3.8 MJ instead of 6.3 MJ. The total radiative power has clearly decreased and the radiative fraction has dropped from $P_{\text{rad}}/P_{\text{in}} = 0.45\text{--}0.30$. The radiative power from the main and divertor plasmas is reduced in comparison with the JET-C reference discharge, by 1.4 MW and 2.1 MW, respectively. In JET-ILW, the plasma is cleaner with the effective charge Z_{eff} reduced from 1.7 to 1.3 owing to the main impurity being now Be instead of C. The neutral pressure in the sub-divertor is similar in both discharges. The drop in confinement in the JET-ILW plasma compared with that for JET-C cannot be attributed to an increased main plasma radiation due to W. In fact, it stems from a drop in the pedestal pressure and mostly temperature. As shown in figure 1, the pedestal electron temperature has dropped from a value of 0.6 keV down to 0.4 keV. More generally, similar plasmas within a fuelling range $(0.9\text{--}3.2) \times 10^{22} \text{ el s}^{-1}$ in the ILW cover a similar range of pedestal electron density as in JET-C, as shown in figure 1, but the pedestal electron temperature is on average 30–40% lower. Although a reduction of the operational window was foreseen if the main plasma radiation increased due to W contamination, this impact of the ILW on the confinement in these high-triangularity plasmas was not anticipated.

Even more surprising is the fact that JET-ILW plasmas with type-I ELMs, such as #82806, exist below the critical electron temperature for type-III ELMs in JET-C at high densities, see figures 1 and 3 [6]. These plasmas have benign ELMs with a slower drop in pedestal pressure than anything noted before [1, 16] and seem to exist in the pedestal n_e and T_e space which in JET-C would have been occupied by the

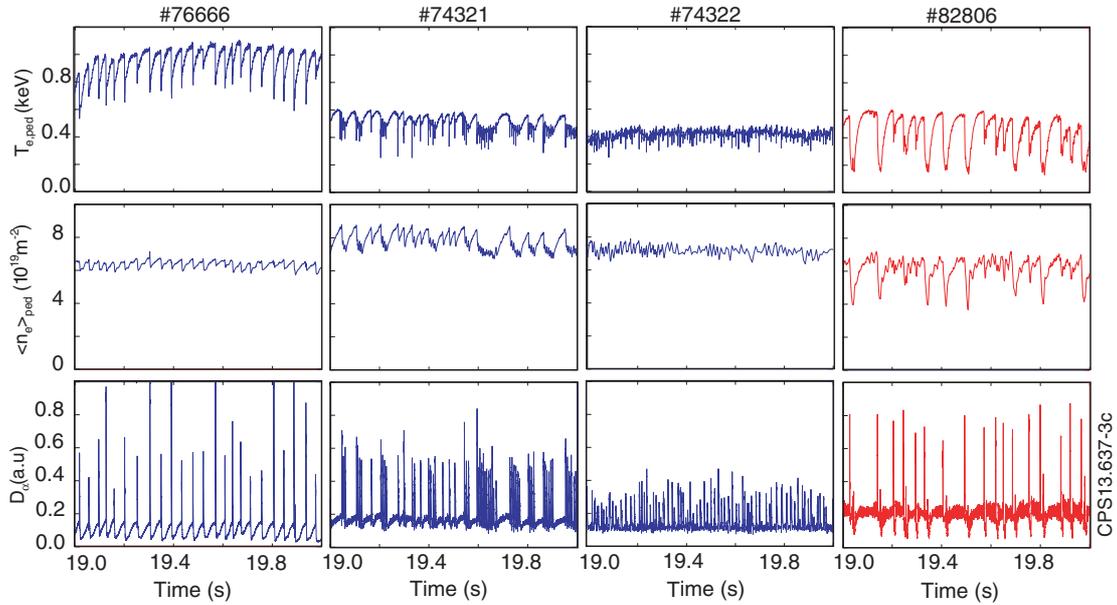


Figure 3. Time trace of pedestal temperature, edge line-integrated density and D_α intensity: discharges in JET-C (in blue) are shown with typical type-I ELM (#76666), in between type-I and type-III ELM regimes (#74321) and in the type-III ELM regime (#74322, $T_{e,ped} \sim 0.55$ keV and $n_{e,ped} \sim 7.1 \times 10^{19} \text{ m}^{-3}$). JET-ILW discharge #82806 (in red) shows a slow ELM collapse of electron temperature and density. The electron and temperature pedestals of the discharges as determined by high-resolution Thomson scattering are shown in figure 1 with the exception of discharge #74322.

type-III ELM regime. The ELM stability regime seems to have changed significantly. The slow ELMs are discussed in more detail in [1, 16]. The boundary between type-I and type-III ELM regimes, which was such a key player in both ELM dynamics and drop in confinement in JET-C discharges as a result of seeding, is not yet fully characterized for JET-ILW, but lies far below its previous position.

4. Recovery of plasma confinement with N_2 seeding in high-shape ELMy H-mode discharges

Although nitrogen was first injected to increase the divertor radiation and to reduce the power loads, nitrogen seeding has been found to improve plasma energy confinement in the high-shape ELMy H-mode. Figure 4 shows time traces of comparable plasmas, seeded with N_2 in JET-ILW and an unseeded JET-C reference case at a similar pedestal density. When nitrogen is injected into deuterium-fuelled discharges in JET-ILW, it raises the pedestal density and temperature leading to an increase in stored energy to 5.5 MJ close to the JET-C D_2 -fuelled counterpart, i.e. #76684 with 5.8 MJ. The ELM frequency drops from ~ 16 to 4 Hz as a result of seeding (was 9 Hz for #76684). Without seeding, ELM crashes are slow and lead to benign power load on the divertor. With N_2 seeding, however, the ELM crashes speed up again and make the ELMs resemble more those seen with the carbon wall [16]. For the range of fuelling level of $(0.8\text{--}3.0) \times 10^{22} \text{ el s}^{-1}$, it can be seen in figure 5 that nitrogen seeding allows access to higher electron pedestal temperatures, only slightly lower than deuterium-fuelled JET-C counterparts. This leads to stored energies and H -factors only slightly below their deuterium-fuelled counterparts as seen in figures 6 and 7, and a very good match to JET-C N_2 -seeded pulses at similar nitrogen seeding rates, particularly when the density dependence is

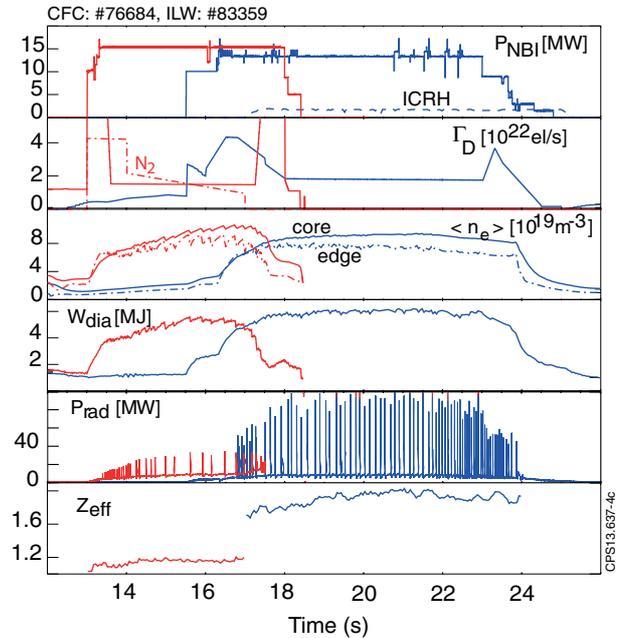


Figure 4. Time traces of unseeded ELMy H-mode discharges in JET-C wall #76684 (in blue), and seeded JET-ILW #83359 (in red): (from top to bottom) NBI and ICRH (dashed) heating power, D_2 -fuelling and N_2 -seeding (dashed) waveforms, line-integrated core and edge (dashed) electron densities measured by an interferometer, diamagnetic stored energy, total radiated power (without the disruption phase for #83359) and Z_{eff} .

removed from the H_{98} -factor scaling to allow the energy confinement time to be compared more directly with plasmas of different densities, as shown in figure 7. The electron density of seeded discharges tends to be higher in JET-ILW than JET-C for the same seeding level. The best N_2 -seeded

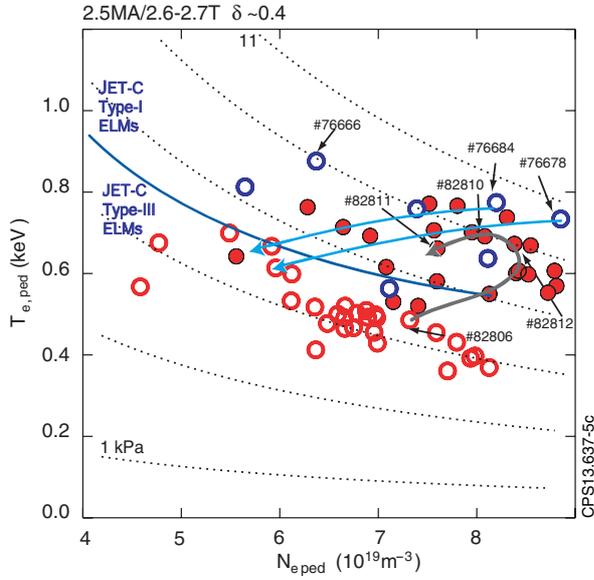


Figure 5. Pedestal n_e - T_e diagram for high- δ ELMy H-mode discharges for fuelled-only JET-C and fuelled and N_2 -seeded JET-ILW discharges with 2.5 MA and 2.6–2.7 T: each symbol corresponds to a discharge. JET-C discharges are in blue symbols. For the sake of clarity the N_2 -seeded JET-C discharges are removed and the trajectory in the n_e - T_e diagram, as N_2 seeding is raised in successive discharges, is shown in light blue for the two highest fuellings shown in figure 3. JET-ILW discharges are shown in red, open symbol for fuelled-only discharges and filled symbol with N_2 seeding. The grey arrow is indicative of the trajectory in pedestal density and temperature of discharges with increasing N_2 -seeding rate from 0 to 3.7×10^{22} el s^{-1} at a D_2 -fuelling rate of 2.8×10^{22} el s^{-1} .

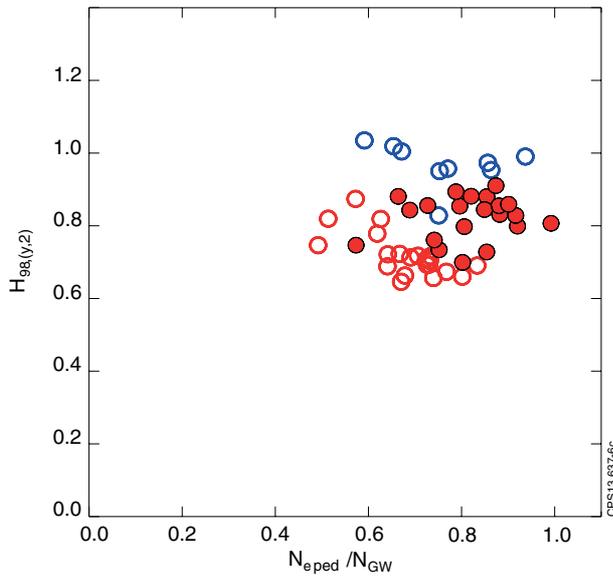


Figure 6. Normalized energy confinement $H_{98(y,2)}$ versus pedestal density normalized to the Greenwald density: symbols and colours are the same as in figure 5.

pulse with the JET-ILW gives $H_{98(y,2)} \sim 0.9$, $f_{\text{GDL}} \sim 1.0$ with $Z_{\text{eff}} \sim 1.5$. The plasma stationarity is not achieved however, and is addressed in section 7.

This is the first time in JET that injection of impurities leads to an increase in global energy confinement. Such a

strong correlation between nitrogen seeding rate and stored energy has already been observed in ASDEX Upgrade and is reported there to be linked to a positive correlation between H -factor and Z_{eff} [10]. In JET-ILW, such a clear dependence with Z_{eff} has not been demonstrated and is discussed in more detail in section 7.

It is important to note that the increase in confinement as a result of the nitrogen injection in the ILW stems from the pedestal and not from an improved energy and particle confinement in the core plasma. The peaking of the electron density and temperature has been calculated taking the ratio of the averaged quantity at $r/a \sim 0.4$ over its value at $r/a \sim 0.8$ and is shown in figure 8 (see figure 9 for the peaking of electron pressure). No change in the electron density peaking is observed between the fuelled-only discharges in JET-C and JET-ILW and the seeded discharges in JET-ILW. The electron temperature peaking for the fuelled JET-C discharges varies but cover a similar range to the JET-ILW discharges whether fuelled only or with the addition of seeding. This indicates that the confinement in these plasmas is set by the achieved pedestal pressure in combination with profile stiffness.

The JET-ILW results strongly suggest that the carbon impurities played a role in the performance of the high-shape plasma scenario in the JET-C era. This is all the more striking if a closer look is taken at the plasma trajectory in the pedestal T_e - n_e diagram following N_2 injection together with the ratio of divertor to main plasma radiative power (as described in section 2), $P_{\text{rad}}^{\text{div}}/P_{\text{rad}}^{\text{bulk}}$, in the inter-ELM phase, as shown in figures 5 and 10. Focusing on the two highest fuelling levels in the JET-C plasmas, one can see that as the N_2 -seeding rate is increased, the ratio $P_{\text{rad}}^{\text{div}}/P_{\text{rad}}^{\text{bulk}}$ rises as expected and the pedestal confinement is slightly reduced up to the transition to the type-III ELM regime where the pedestal confinement drops by 20% at the highest divertor radiation. In comparison, for JET-ILW plasmas as the nitrogen seeding rate is increased, the ratio $P_{\text{rad}}^{\text{div}}/P_{\text{rad}}^{\text{bulk}}$ is raised and the pedestal confinement increases up to values similar to those in the deuterium-fuelled JET-C counterparts. Once the maximum pedestal pressure has been achieved with nitrogen seeding, a further rise in seeding level leads to a decrease in density and a weaker decrease in temperature, as shown in figure 5, very like the trajectories of plasmas in JET-C with N_2 seeding. Similarly, the ELM frequency then starts to increase in JET-ILW as was the case for the JET-C plasmas. This is discussed in more detail in a follow-up paper [17]. For completeness, the effective charge Z_{eff} is shown in figure 11 and shows that with increased N_2 -seeding rate, Z_{eff} is increased in JET-ILW but remains at the highest seeding rate lower than the fuelled JET-C counterpart. These observations indicate that C concentration and the associated radiation may have been hidden parameters in the JET-C confinement behaviour in high-shape fuelled discharges, and that the injection of N_2 in JET-ILW partially recovers this effect. The exact mechanism is unknown. It remains to be identified whether the effect is a result from the ion dilution, which is addressed in section 7, a change in resistivity or local current density or possible turbulence suppression.

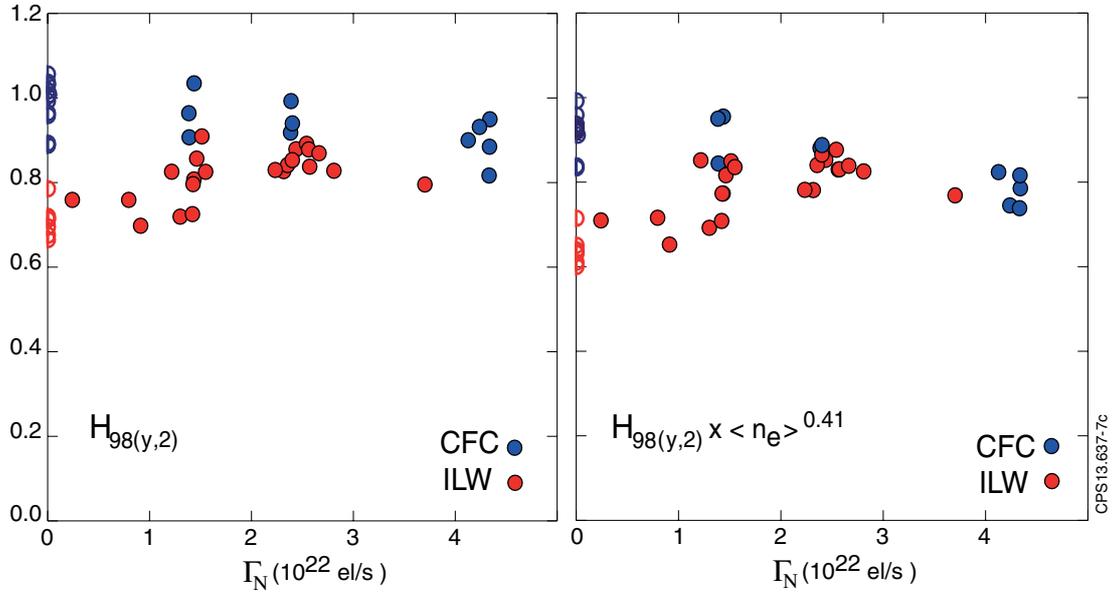


Figure 7. Dependence of normalized confinement with N_2 -seeding rate: (on the left) with the $H_{98(y,2)}$ density scaling and (on the right) with the density scaling removed. Symbols and colours are the same as in figure 5. Filled blue symbols correspond to N_2 -seeded discharges in JET-C.

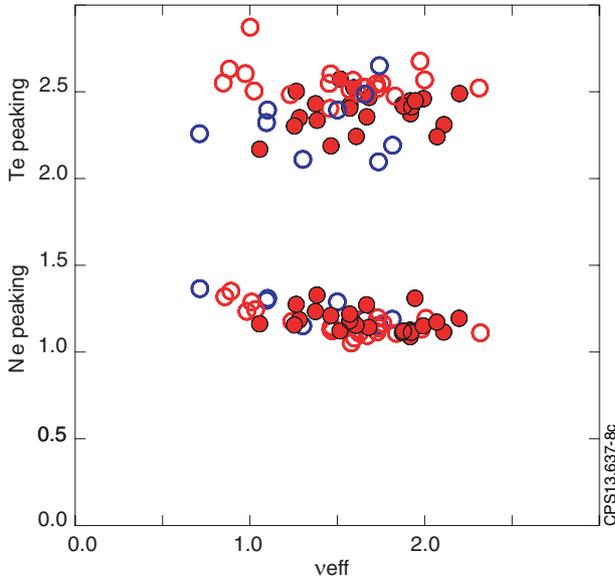


Figure 8. Dependence of temperature peaking and density peaking with ν_{eff} : symbols and colours are the same as in figure 5. The peaking is calculated by taking the average quantity at $r/a \sim 0.4$ over its value at $r/a \sim 0.8$. ν_{eff} is defined as the ratio of the electron collision frequency to the curvature drift frequency and is evaluated by the simple relation $\nu_{\text{eff}} = 0.1 \cdot R Z_{\text{eff}} \cdot n_e \cdot T_e^2$ where the major radius R is in m, the local density n_e is in 10^{19} m^{-3} and the electron temperature T_e is in keV.

5. Inter-ELM power load reduction in JET-ILW high-triangularity discharges

Although the increase in pedestal stored energy is a welcome benefit of seeding, nitrogen injection was first envisaged as a way to increase the divertor radiation and thereby reduce the inter-ELM power load to the divertor. With a factor 10 reduction of C in the ILW, the resulting reduced divertor radiation is expected to lead to higher power loads and

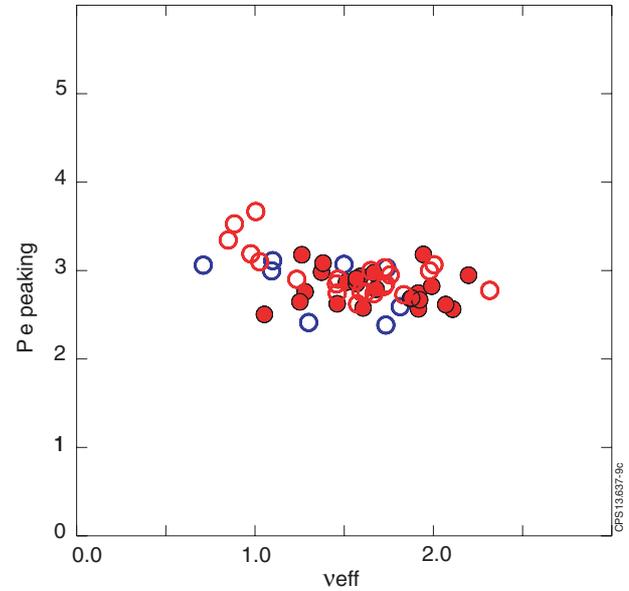


Figure 9. Dependence of the electron pressure peaking with ν_{eff} : symbols and colours are the same as in figure 5.

higher electron temperatures at the outer target in the inter-ELM phases.

In JET-C, for the D_2 -fuelling rate of $2.7 \times 10^{22} \text{ el s}^{-1}$ (#76678), the radiative fraction is ~ 0.5 (see figure 12) with a ratio of divertor to main plasma radiative power ($P_{\text{rad}}^{\text{div}}/P_{\text{rad}}^{\text{bulk}}$) of ~ 0.7 in the inter-ELM phase, as shown in figure 10. In JET-ILW at the same D_2 -fuelling rate as above (#82806), the radiative fraction is down to ~ 0.3 with a ratio ($P_{\text{rad}}^{\text{div}}/P_{\text{rad}}^{\text{bulk}}$) also reduced down to ~ 0.3 . As the N_2 -seeding rate is increased in successive discharges in JET-C, the radiative fraction is increased to the highest value of ~ 0.55 (N_2 -seeding rate of $\sim 2.4 \times 10^{22} \text{ el s}^{-1}$) whilst still in the type-I ELM regime with a ratio $P_{\text{rad}}^{\text{div}}/P_{\text{rad}}^{\text{bulk}}$ of ~ 0.78 . Similarly for the JET-ILW

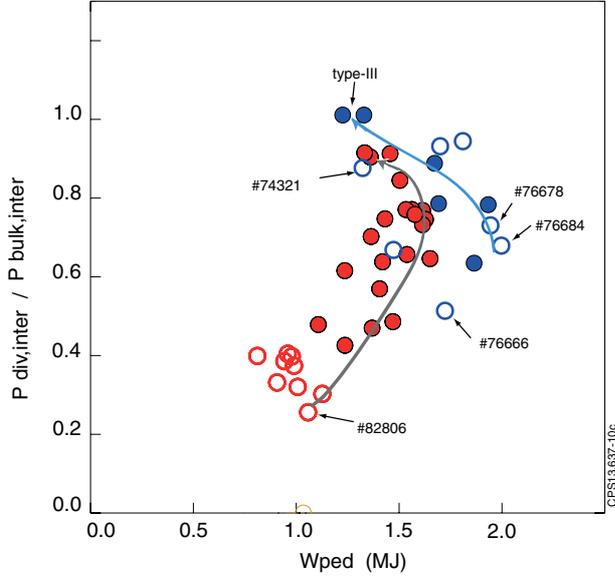


Figure 10. Ratio of inter-ELM phase divertor radiation over main plasma radiation versus pedestal stored energy: for JET-ILW (in blue) and JET-C (in red), with fuelling only (in open circles) and with additional seeding (in full circles). The light blue and grey arrows are the same as in figure 5. W_{ped} is the pedestal stored energy and is calculated from density and temperature profiles, averaged over the time window of interest. W_{ped} is equal to $3/2k_B(n_e \cdot T_e + n_i \cdot T_i) \cdot V_{plasma}$ with electron density n_e , electron temperature T_e , main ion density n_i and ion temperature T_i evaluated at the pedestal top and V_{plasma} the total plasma volume.

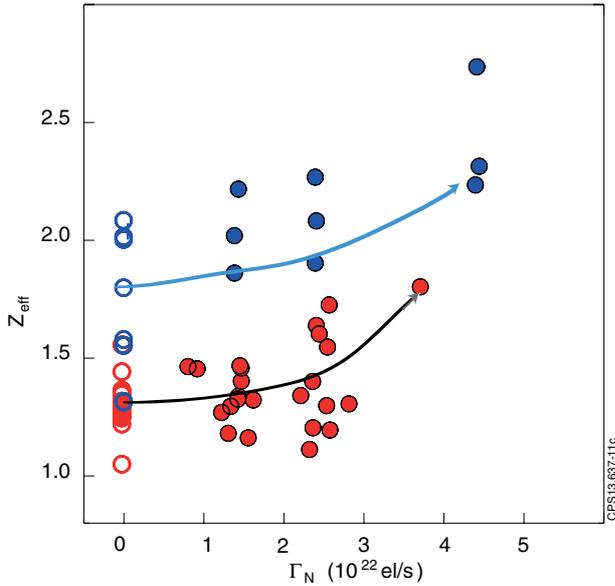


Figure 11. Dependence of line-averaged Z_{eff} versus N_2 -seeding rate: Z_{eff} is measured from line-integrated bremsstrahlung radiation. Symbols and colours are the same as in figure 10. Arrows follow the same trajectory as in figure 10.

discharges with the same fuelling rate and at nitrogen seeding rate $\sim 2.35 \times 10^{22} \text{ el s}^{-1}$ at the highest pedestal stored energy, the radiative fraction reached ~ 0.52 (#82810) with $P_{rad}^{div}/P_{rad}^{bulk}$ of ~ 0.78 , similar conditions in terms of radiative fraction and distribution to the JET-C counterpart. At the higher seeding rate of $\sim 3.7 \times 10^{22} \text{ el s}^{-1}$, the f_{rad} reached ~ 0.54 (#82811)

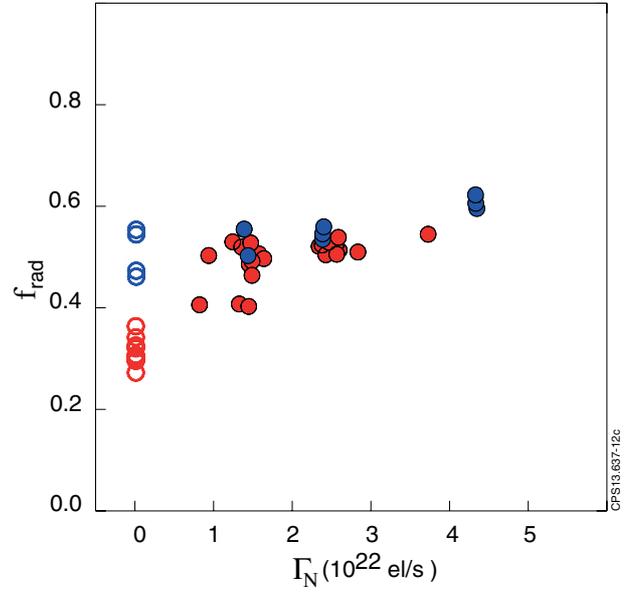


Figure 12. Total radiative fraction versus the N_2 -seeding rate: symbols and colours are the same as in figure 10.

with $P_{rad}^{div}/P_{rad}^{bulk}$ of ~ 0.90 with a somewhat degraded pedestal stored energy. In fact, independently of the fuelling rate, with a high enough nitrogen seeding rate in the JET-ILW, similar conditions of radiative fraction and value of the ratio $P_{rad}^{div}/P_{rad}^{bulk}$ can be obtained in JET-ILW as was obtained in the JET-C N_2 -seeded studies, see figures 10 and 12. It is then expected that the power load would be the most different between JET-C and JET-ILW at no or low N_2 seeding and become similar at a high level of seeding.

If the focus now is on the measured power landing on the outer target in the inter-ELM phase, we need to study the infra-red (IR) or Langmuir probe (LP) measurements. In the JET-C discharges, the power loads were monitored by IR. The IR measurements were validated against thermocouple data and an energy balance was carried out. It was found that the measured energy loss was 80% of the input energy over the N_2 seeding scan [6]. In other words, 20% of the input energy was unaccounted for in the energy balance. For the JET-ILW discharges, technical difficulties linked to the high- δ configuration meant that IR measurements were not available. The thermocouples, embedded in the tile on which the outer strike point (OSP) was, were unreliable. The power load then has to be monitored by LPs [18]. The power load in the inter-ELM period can now be compared between the JET-C and JET-ILW discharges. A higher power load is expected for JET-ILW than JET-C for a similar fuelling level. The power P_{sep}^{inter} flowing through the separatrix in the inter-ELM phase can be calculated with $P_{sep}^{inter} = P_{tot}^{inter} - P_{rad,main}^{inter} - dW^{inter}/dt$ with P_{tot}^{inter} the total input power (taking into account the shine-through losses but not the fast-ion losses), $P_{rad,main}^{inter}$ the main plasma radiated power and finally dW^{inter}/dt the inter-ELM stored energy build-up rate. Relying on results presented in [6] for JET-C discharges, at the D_2 -fuelling level of $2.7 \times 10^{22} \text{ el s}^{-1}$ (#76678), the power flowing through the separatrix in the inter-ELM phase, P_{sep}^{inter} , is $\sim 8.5 \text{ MW}$ calculated from a total input power of 15.8 MW

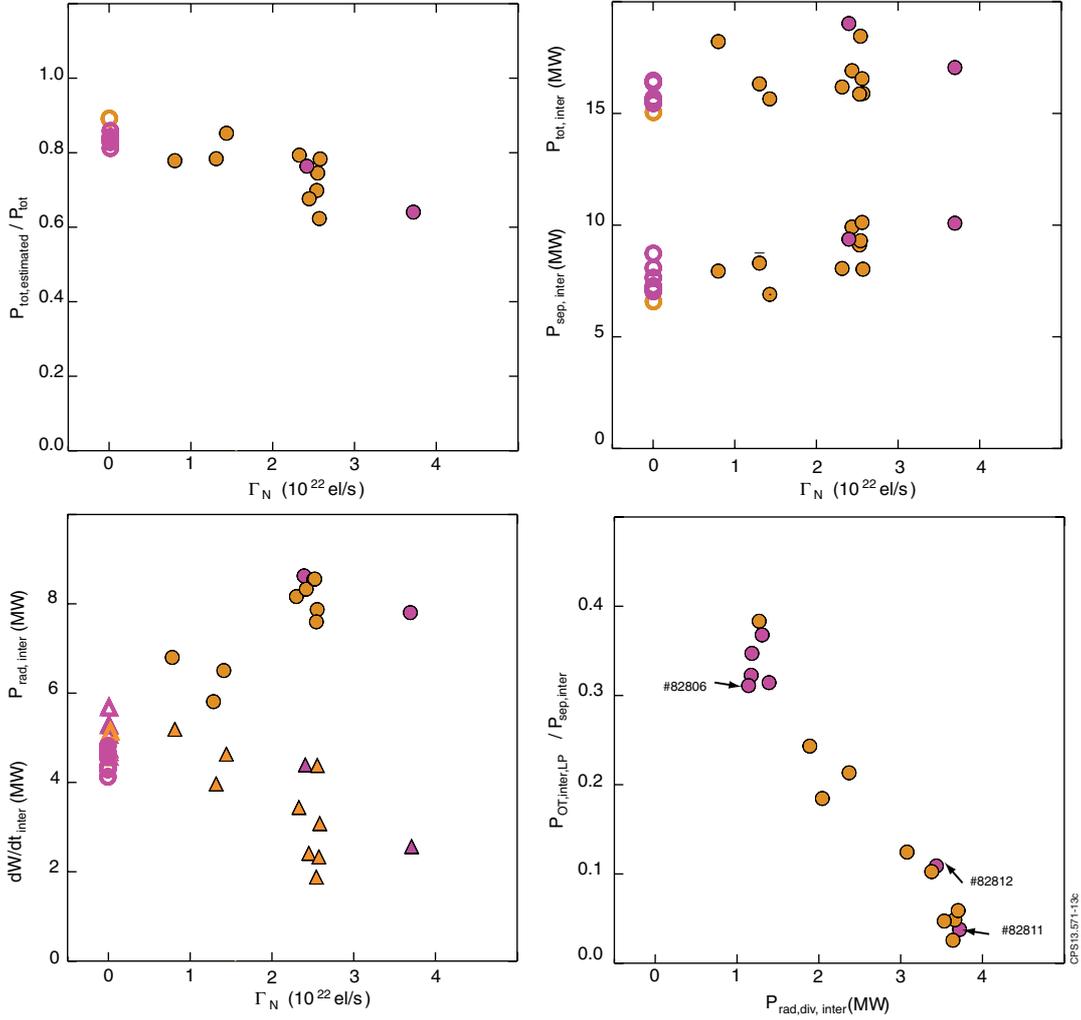


Figure 13. From left to right, top row: ratio of estimated loss power over total input power, total input power and power flowing through the separatrix, $P_{\text{sep,inter}}$; bottom row: total radiated power and the stored energy build-up rate, the ratio of power measured by the LP at the outer target over $P_{\text{sep,inter}}$. All graphs show inter-ELM averaged quantities and are versus the N_2 -seeding rate. Only JET-ILW discharges are shown here. Pink discharges correspond to discharges with the D_2 -fuelling rate between 2.6 and $3.0 \times 10^{22} \text{ el s}^{-1}$ and the orange one between 0.8 and $2.0 \times 10^{22} \text{ el s}^{-1}$. Triangles indicate the stored energy build-up rate dW^{inter}/dt .

from which $P_{\text{rad,main}}^{\text{inter}}$ of $\sim 4.8 \text{ MW}$ and dW^{inter}/dt of $\sim 2.3 \text{ MW}$ were subtracted. The power landing on the outer target plate is 1.9 MW and therefore a fraction $f_{\text{div,OT}}^{\text{inter}} \sim 0.22$ of $P_{\text{sep}}^{\text{inter}}$ reaches the outer divertor. The JET-ILW discharge with the same fuelling rate has a similar power $P_{\text{sep}}^{\text{inter}}$ of $\sim 7.6 \text{ MW}$ calculated with an input power of 16 MW , with $P_{\text{rad,main}}^{\text{inter}} \sim 3.5 \text{ MW}$ and $dW^{\text{inter}}/dt \sim 5.6 \text{ MW}$, as shown in figure 13. The power reaching the outer target is 2.4 MW and a fraction 0.3 of $P_{\text{sep}}^{\text{inter}}$. The electron temperature at the OSP is $\sim 26 \text{ eV}$ in JET-ILW compared with $\sim 5 \text{ eV}$ for its JET-C counterpart [6]. As expected, the power load and electron temperature are higher in JET-ILW than in the JET-C plasmas for a similar fuelling but weaker increase as could have been expected from a reduction of the C content by a factor of 10. Related studies are currently investigating this question [12].

When the N_2 -seeding rate is increased in the JET-C discharges up to $\sim 2 \times 10^{22} \text{ el s}^{-1}$ (focusing here on a discharge with type-I ELM regime at a D_2 -fuelling rate of $2.2 \times 10^{22} \text{ el s}^{-1}$), the power $P_{\text{sep}}^{\text{inter}}$ is slightly decreased due to the increase in the dW^{inter}/dt term to 7.2 MW . Partial

detachment was reached at this seeding level, as shown in [6]. The peak power load is reduced down from 2.8 MW m^{-2} (#76678) to below 1 MW m^{-2} (#76680) at the N_2 -seeding rate of $2.4 \times 10^{22} \text{ el s}^{-1}$ and within the error of the IR measurements. For the JET-ILW counterpart with the same fuelling rate, as the N_2 -seeding rate is raised, the power $P_{\text{sep}}^{\text{inter}}$ is in fact increased slightly from 7.6 to 9.4 MW as dW^{inter}/dt decreases with a reduction of the ELM frequency. The power measured on the outer divertor is 1.0 MW (#82812) and a decrease from 0.3 to 0.1 of $P_{\text{sep}}^{\text{inter}}$. With additional seeding of $3.7 \times 10^{22} \text{ el s}^{-1}$, the ratio $f_{\text{div,OT}}^{\text{inter}}$ can be reduced down to 0.03 bearing in mind that the inter-ELM phase become difficult to measure for discharge #82811 as the ELM frequency rises again. In any case, the reduction in the power at the outer target is achievable at high enough N_2 seeding and a wide D_2 -fuelling rate of $(0.8\text{--}2.5) \times 10^{22} \text{ el s}^{-1}$. The electron temperature is reduced to $\sim 7 \text{ eV}$ (#82812) and reaches similar conditions to those in discharge #76680 with partial detachment and in type-I ELM regime. More detailed studies on detachment will be presented in a follow-up paper [17].

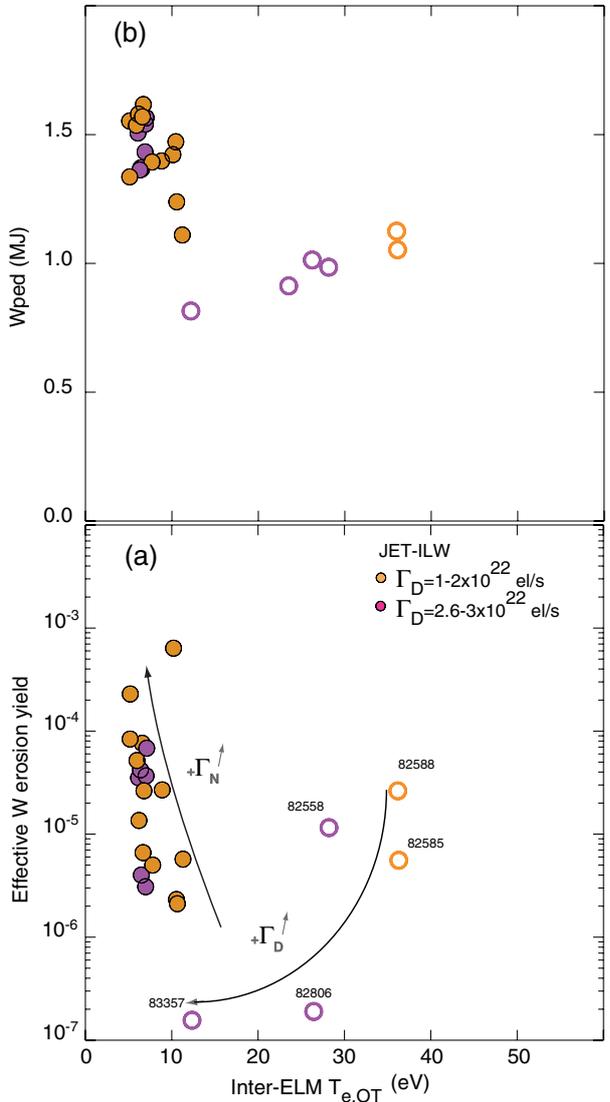


Figure 14. Effective W erosion yield (a) and average pedestal stored energy versus target electron temperature in an inter-ELM period (b).

6. Stationarity of the plasma and average W erosion yield

The W release is governed by impurity ions due to their threshold energy for W sputtering being an order of magnitude lower than that for deuterium. The sputtering yield for impurity ions such as C^{4+} , Be^{2+} and N^{4+} rises dramatically for plasma temperatures above ~ 5 eV [19, 20]. At ~ 5 eV, it is expected that no significant W will be sputtered by low-Z impurities in between ELMs. Tungsten will, however, still be sputtered due to ELMs.

In these experiments, W erosion was measured by means of passive emission spectroscopy. The W_I line radiation at 400.9 nm is monitored by a mirror-link system [21] with a time resolution of 40 ms, mostly able to provide an ELM-averaged measurement. The measured light intensities were transformed into W particle flux densities using the number of ionizations per emitted photons, i.e. S/XB value as in [22]. The effective erosion yield is then obtained by normalization of the

ELM-averaged peak W particle flux to the saturation current measured by the LP. Results are shown in figure 14. It can be seen that for fuelled-only plasmas (open symbol) where the pedestal pressure is weak and ELM benign, the erosion does decrease with target temperature as would be expected for averaged erosion dominated by the inter-ELM phase. When N_2 is seeded, the effective erosion increases even though the electron temperature is decreased to a value of ~ 5 eV, thus all sputtering is due to high energetic ions hitting the target plate during ELMs [13]. The ELM energy density responsible for material erosion is directly related to the pedestal pressure [23]. In fact, the increased erosion yield at a low target temperature is directly linked to an increase in pedestal confinement and therefore pressure, as shown in figure 14. In other words, as nitrogen is seeded and the pedestal confinement increases, the W erosion yield is governed by the ELMs and not the inter-ELM phase. This is compounded by the fact that transport of W into the core through the pedestal region is a critical step in a chain, which makes it a dominant factor affecting the build-up of W in normal H-modes [24, 25]. Previous studies in D_2 -fuelled discharges (in which Be is the main impurity responsible for W sputtering) have shown that the intra-ELM W sputtering dominated by a factor of 5 over the inter-ELM phase [13]. In our discharges, N is present, and being heavier than Be is expected to sputter W more than Be during an ELM.

The N_2 -seeded high- δ ELMy H-mode considered provides the first case in JET-ILW of a scenario with power load control though high ELM energy density loads together with long enough plasma duration to probe the impact of W on plasma confinement and stationarity. Time traces of a typical N_2 -seeded discharge with $H_{98} \sim 0.9$, $n/n_{GW} \sim 1$ and $Z_{eff} \sim 1.5$ are shown in figure 15. This figure illustrates that N_2 -seeded high- δ ELMy H-mode plasmas with higher confinement so far have an unstationary behaviour. Although the total radiative power and energy confinement are constant over most of the plasma duration, the core electron temperature decreases, while the plasma radiation within the sawtooth (ST) inversion radius increases. Once the sawteeth disappears, the core radiation increases exponentially. W accumulates in the very plasma centre (within $\rho \sim 0.3$) and when the plasma enters its termination phase leads to high radiative power and finally a disruption. It is important to note that high- δ ELMy H-modes were not stationary in JET-C with Z_{eff} constantly increasing, as shown in figure 2. In fact, ICRH was added to neutral-beam heating as an integral part of the scenario for control of density peaking and sawteeth [8, 26]. In the JET-ILW discharge, the first test with ICRH heating with a rather limited maximum power of ~ 2 MW (due to technical problems) did not show much benefit, but further investigations will be carried out in the future. The density peaking outside the ST radius is the same in ILW as in JET-C. As a result, it could be that the impurity peaking is similar to what was observed in JET-C, just made more evident by the presence of an additional and much higher Z impurity such as W.

To develop a long discharge of about 20 s flat-top it is essential to improve the plasma stationarity by reducing first W penetration through the separatrix and second W peaking inside the main plasma. The first point needs to be addressed by a reduction of the W source or transport from the source to the pedestal both of which can be addressed with different plasma

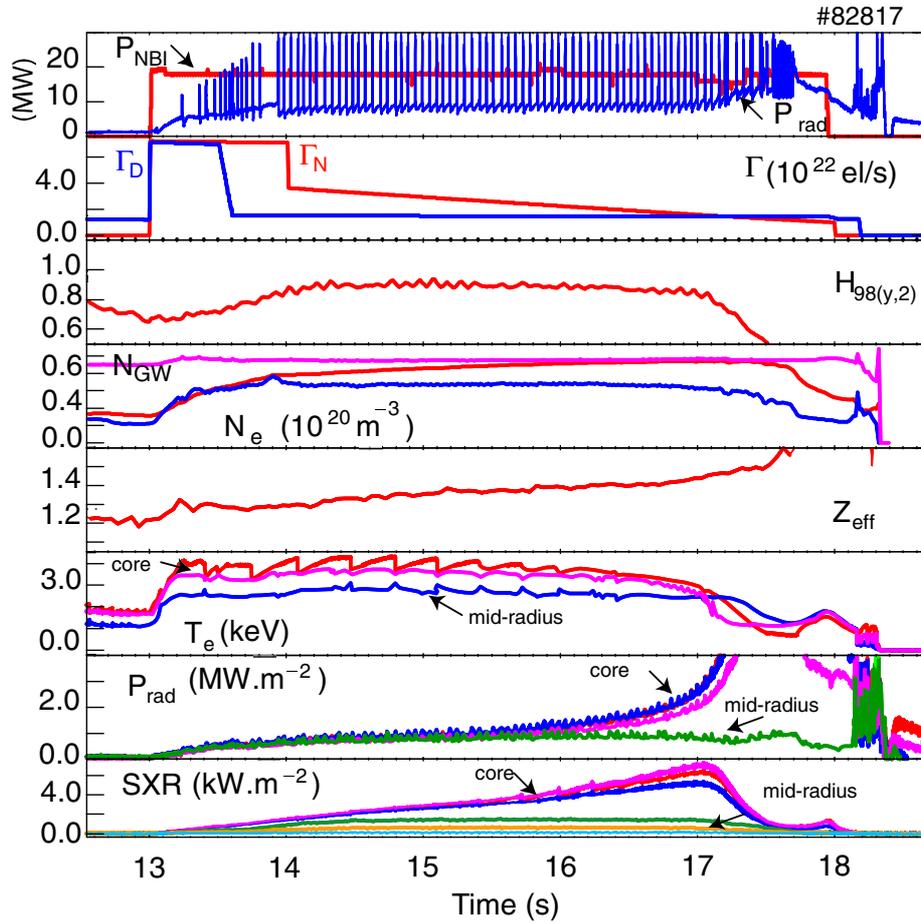


Figure 15. Time traces of N_2 -seeded ELMy H-mode discharges in JET-ILW wall #82817 (from top to bottom) NBI heating and total radiated power, D_2 and N_2 waveforms, normalized confinement factor H_{98} , core (red) edge (in blue) and Greenwald electron density (in pink), line-integrated Z_{eff} , electron temperature from the core (in red and pink) and mid-radius (in blue), radiated power from the core (in blue and pink) and mid-radius (in green) channels, and finally soft x-ray radiated power.

divertor configurations. The second point is likely to require the use of core heating for either controlling the ST behaviour or some possible beneficial outward turbulent convection.

7. Lack of improved confinement with N_2 seeding in other configurations and discussion of the possibility of dilution being the reason for improvement of confinement

Observations presented so far have indicated that the C level and associated radiation may have been hidden parameters in the JET-C confinement behaviour in high-shape fuelled discharges. This could be why such discharges in JET-ILW show up to $\sim 40\%$ decrease in confinement compared with the JET-C counterpart, and why the injection of N_2 partially recovers this confinement. In contrast, it is not believed that C radiation is a hidden parameter for low- δ ELMy H-mode discharges. The energy confinement of low- δ ELMy H-mode discharges is reported to be comparable in JET-ILW and JET-C [16], the only difference being that the low fuelling operation, favourable for high normalized confinement, is no longer accessible due to W contamination in the main plasma. Nevertheless, as a means to gather further evidence on the effect of nitrogen, this seed impurity was injected into the low-

δ ELMy H-mode at 2.5 MA and a similar divertor geometry. A low- δ ELMy H-mode with vertical targets, more similar to the ITER divertor geometry, was also tested. No significant increase in confinement with N_2 seeding was observed. The range of Z_{eff} values covered with the low- δ configuration was from 1.2 to 1.5 and similar to that of the high- δ configuration (1.1–1.8) (see figure 11), demonstrating the lack of dependence of confinement on the Z_{eff} value.

The vertical target configuration proved interesting for two reasons. Firstly, the W contamination in the main plasma was significantly reduced and the operational window seemed to extend to very low fuelling close to levels used in JET-C. The very first results show signs of a higher stored energy with no additional fuelling but further work is needed to confirm this. Secondly, W radiation events (attributed to small particles of medium- and high-Z materials entering the plasma [1]) were fairly frequent with the horizontal target divertor configuration [1], but with the vertical target none have been observed so far. These results will be presented in detail in future papers.

Discussion on a possible dilution effect to explain the increase in stored energy. In this paper, we presented the increase in stored energy following the injection of nitrogen as an effect due to the increase in the pedestal stored energy by studying the

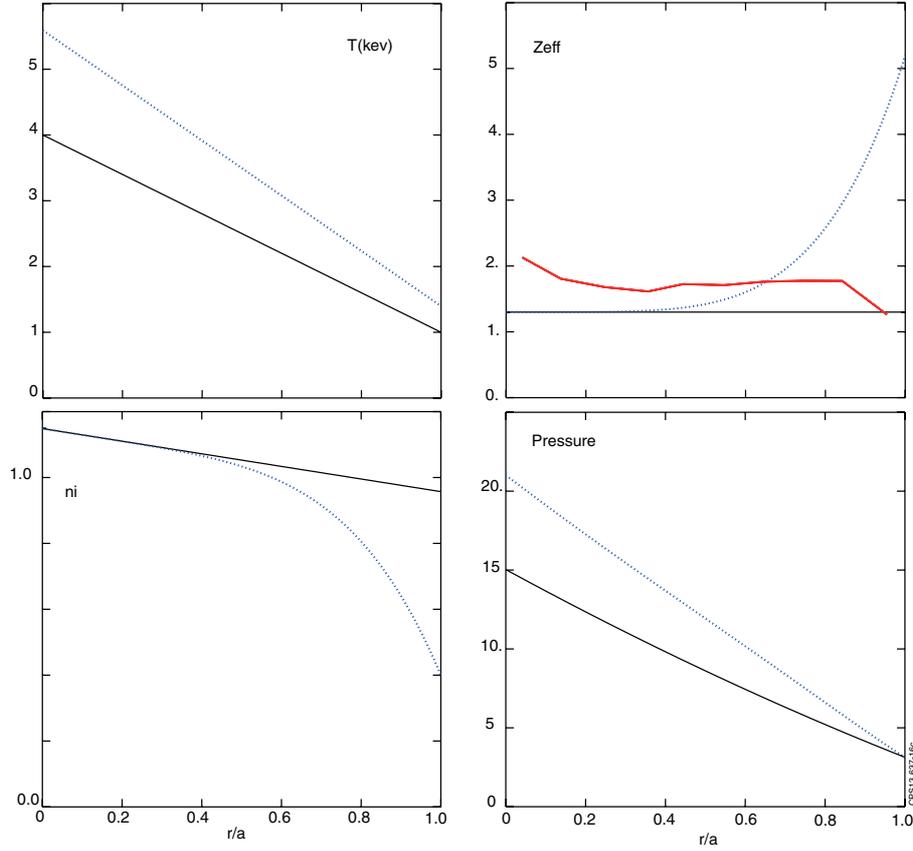


Figure 16. Schematic of increase in pedestal temperature due to a dilution effect compatible with an increase in pedestal temperature, Z_{eff} and total stored energy seen in the experiment: in black the example profiles before and in blue after N_2 seeding. The Z_{eff} profile in red is typical of those measured in JET-C seeded studies.

pedestal electron density and temperature implying tacitly that the main ion temperature and density have a similar behaviour. It is indeed a possibility that the measured increase in the electron channel could not be related to an increase in the total pedestal pressure but to a dilution effect as is advanced for ASDEX Upgrade to explain the increase in confinement following N_2 seeding [10]. For the sake of argument, if we were to make this hypothesis and state that there is no total improvement of the pedestal pressure we still need to be able to find the conditions compatible with the experiments, meaning the same Z_{eff} , $T_{e,\text{ped}}$ and $n_{e,\text{ped}}$ increase, and the total increase in the stored energy. Assuming that the total pedestal pressure P_{ped} remains unchanged with the injection of nitrogen:

$$P_{\text{ped}} = k_B \times (n_{e,\text{ped}} T_{e,\text{ped}} + n_{i,\text{ped}} T_{i,\text{ped}}) = k_B \times T_{\text{ped}} \times (n_{e,\text{ped}} + n_{i,\text{ped}}) = C$$

with $n_{i,\text{ped}}$ the total ion density and $T_{i,\text{ped}}$ the ion temperature. Here $T_{e,\text{ped}} = T_{i,\text{ped}}$ is taken as is confirmed by charge exchange recombination measurements. The ratio T_e/T_i is not measured due to the absence of core ion temperature measurements, but T_e/T_i equal to 1 can also be expected for the whole of the profile since the equipartition time between ions and electrons is an order of magnitude lower than the energy confinement time at the high densities considered here. The increase in $T_{e,\text{ped}}$ by a factor of 1.4 (ignoring the increase in density $n_{e,\text{ped}}$ from now on by a factor of 1.1) measured in the experiment (#82810 compared with #82806) would need to be compensated by a

decrease in $n_{i,\text{ped}}$ by about at least 1.4. Assuming nitrogen (with atomic number $z_k \sim 7$) to be the main impurity ions, and invoking quasi-neutrality,

$$n_{i,\text{ped}} = n_{e,\text{ped}} \times ((z_k + 1) - Z_{\text{eff,ped}}) / z_k = n_{e,\text{ped}} \times ((8 - Z_{\text{eff,ped}}) / z_k).$$

The pedestal $Z_{\text{eff,ped}}$ would have to increase from the unseeded level of ~ 1.3 (assuming a flat Z_{eff} profile) to ~ 4 in order to achieve the required reduction in $n_{i,\text{ped}}$. Bearing this in mind, we can now turn our attention to the rest of the profile. In the experiment, $T(r=0)/T_{\text{ped}} \sim 4.5$ and this ratio can be assumed to stay constant as we have seen that no change in electron temperature peaking is observed in the experiment, see figure 8. The Z_{eff} profile necessary to explain the increase in stored energy is compatible with the increase in line-integrated Z_{eff} measurement from 1.3 to 1.4 and $Z_{\text{eff}}^{\text{ped}} \sim 4$. Such a profile is shown in figure 16 together with the profile of the stored energy. It can be seen that the Z_{eff} profile would have to be very hollow, as shown in figure 16, to be compatible with a pedestal temperature increase of a factor of 1.4, with a central Z_{eff} of 1.4 and a 1.4 increase in the total stored energy. Although not measured at present in JET-ILW, the Z_{eff} profile was measured in the JET-C discharges and is slightly hollow, as shown in figure 16. As the density and temperature gradients have not significantly changed for the considered discharges between JET-ILW and JET-C, in the region $r/a \sim 0.4$ – 0.8 , it is not expected that the impurity

convection and diffusion coefficients change substantially and as a result the peaking of Z_{eff} is expected to be similar for JET-ILW and JET-C discharges.

Furthermore, if the dilution effect was the explanation for the increase in stored energy it should be independent of the plasma configuration. But we have seen that a similar rise in Z_{eff} is measured in low- and high- δ configuration but such an increase in total plasma pressure has not been observed. It is then highly unlikely that the increase in confinement observed in the high-triangularity configuration with N seeding is related to a dilution effect. Further investigations are considering a change in pedestal stability, which is outside the scope of this paper.

8. Conclusion

It was expected that the change from carbon-fibre composite (CFC) to ILW at JET would limit the operational window to higher fuelling levels both for low- and high-shape ELMy H-mode discharges, in order to keep the W contamination and induced radiation low in the main plasma. However, it was not anticipated that the new PFCs would reduce the pedestal confinement of the high-shape ELMy H-mode in comparison with the JET-C counterpart. The injection of an extrinsic impurity, N_2 , in the high- δ ELMy H-mode was shown to recover the pedestal confinement. This paper suggests that C content could have been a hidden parameter in the JET-C confinement behaviour of high-shape plasmas and that N_2 partially recovers this effect in JET-ILW. The exact mechanism has not been identified so far. It was shown that once the pedestal pressure of typical discharges in JET-C is mostly recovered, the ELM governs the average W erosion yield, preventing from achieving stationary conditions so far. Operation at high H -factor and high density is paramount for JET-ILW and ITER. The next step for the development of JET-ILW compatible scenarios will be to integrate techniques to control plasma stationarity to open up the operational space to long-pulse and high-current operation.

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