

UKAEA-CCFE-PR(19)42

Matthew J. Lloyd, Robert G. Abernethy, Mark R. Gilbert, Ian Griffiths, Paul A. J. Bagot, Duc Nguyen-Manh, Michael P. Moody, David E. J. Armstrong

# **Decoration of Voids with Rhenium and Osmium Transmutation Products in Neutron Irradiated Single Crystal Tungsten**

Enquiries about copyright and reproduction should in the first instance be addressed to the  
UKAEA  
Publications Officer, Culham Science Centre, Building K1/0/83 Abingdon, Oxfordshire,  
OX14 3DB, UK. The United Kingdom Atomic Energy Authority is the copyright holder.

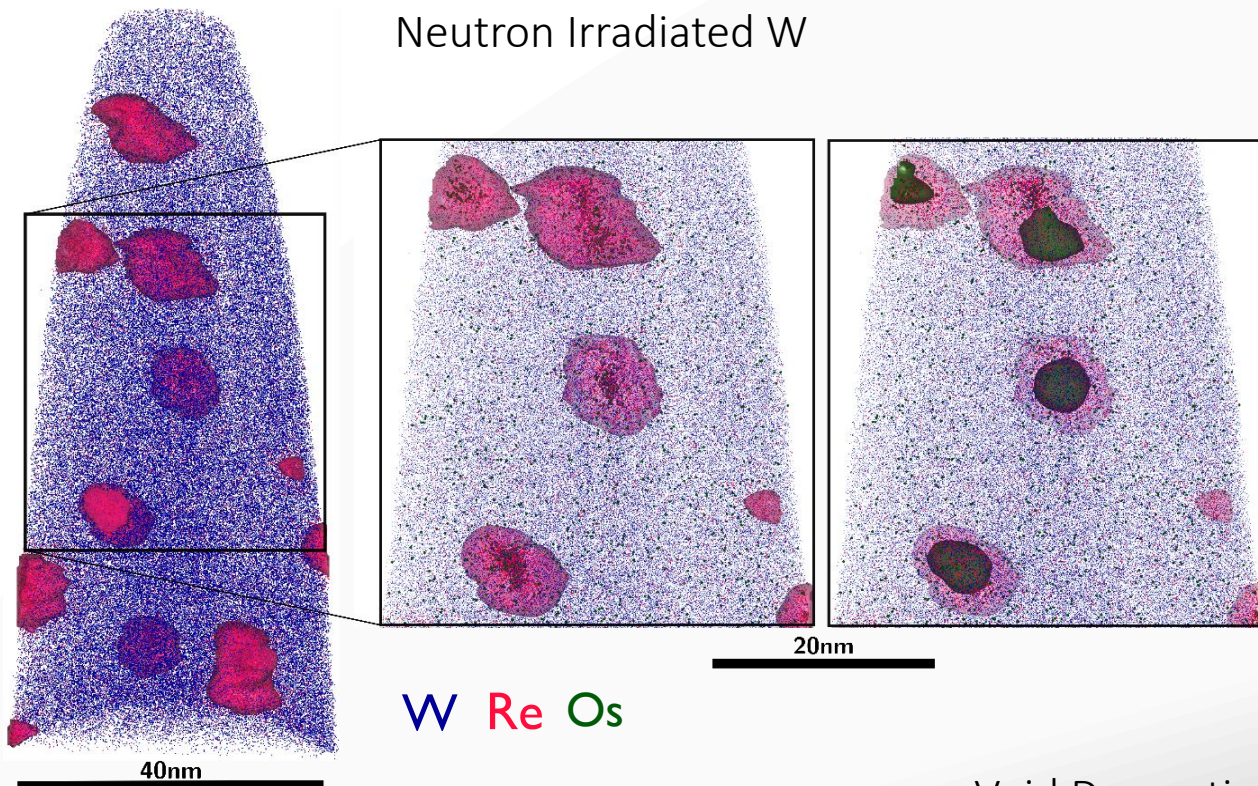
# **Decoration of Voids with Rhenium and Osmium Transmutation Products in Neutron Irradiated Single Crystal Tungsten**

Matthew J. Lloyd, Robert G. Abernethy, Mark R. Gilbert, Ian Griffiths, Paul A. J. Bagot, Duc Nguyen-Manh, Michael P. Moody, David E. J. Armstrong

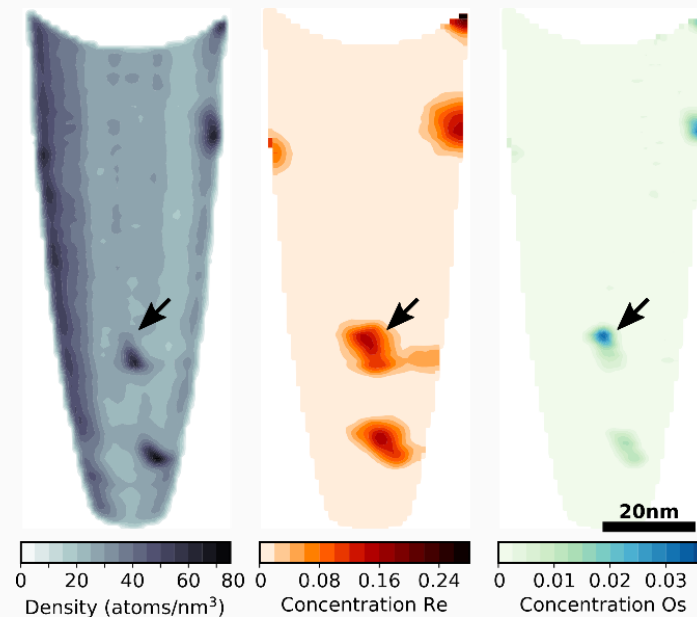


# ATOM PROBE TOMOGRAPHY

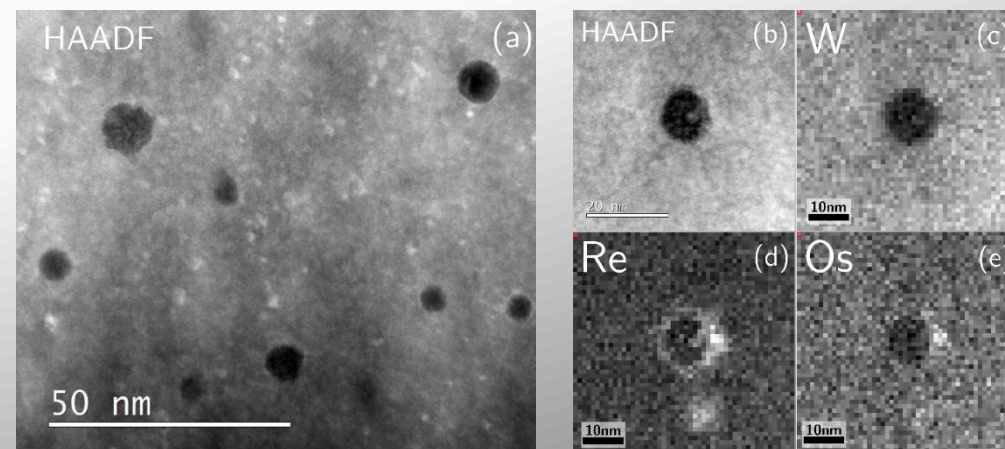
Clustering of Re/Os in High Temperature Neutron Irradiated W



Variation in the Atomic Density During APT Reconstruction



## HAADF + STEM-EDX



Void Decoration with Re/Os in STEM-EDX

## Decoration of Voids with Rhenium and Osmium Transmutation Products in Neutron Irradiated Single Crystal Tungsten

Matthew J. Lloyd<sup>a,b,\*</sup>; Robert G. Abernethy<sup>a,b</sup>; Mark R. Gilbert<sup>b</sup>; Ian Griffiths<sup>a</sup>; Paul A. J. Bagot<sup>a</sup>; Duc Nguyen-Manh<sup>b</sup>; Michael P. Moody<sup>a</sup>; David E. J. Armstrong<sup>a</sup>;

a) Department of Materials, University of Oxford, Parks Road, Oxford, UK, OX13PH

b) Culham Centre for Fusion Energy, Culham Science Park, Abingdon, Oxfordshire, UK

matthew.lloyd@materials.ox.ac.uk\*

Keywords: High Angle Annular Dark Field (HAADF), Atom Probe Tomography (APT), Tungsten Rhenium Osmium Alloys

### Abstract

High temperature, neutron irradiated single crystal tungsten, with a post irradiation composition of W-1.20±0.11at.%Re-0.11±0.05at.%Os-0.03±0.01at.%Ta was characterised using a combination of Atom Probe Tomography (APT) and Scanning Transmission Electron Microscopy (STEM). APT showed that within nanoscale clusters of Re/Os, the atomic density was above the theoretical limit. Complimentary High Angle Annular Dark Field (HAADF) imaging shows that some clusters contain voids at their centre which are leading to APT aberrations and enhancing the atomic density. High resolution Energy Dispersive X-ray (EDX) spectroscopy shows that voids are decorated with a shell of rhenium with a small osmium cluster to one side.

1 A demonstration fusion power station (DEMO) relies on the development of plasma facing materials that can  
2 reliably withstand the combined effects of a large applied heat flux ( $10\text{-}20\text{MWm}^{-2}$  [1]) and a large dose of  
3 14MeV neutrons. Tungsten is currently the leading candidate material because it has a high melting  
4 temperature (3695K [2]), a good thermal conductivity and a high resistance to sputtering through interaction  
5 with the plasma [3]. Radiation damage and transmutation at high operating temperatures ( $\sim 1300\text{K}$  [4])  
6 results in an increase of the inherent brittleness of tungsten, and causes an increase in its already high  
7 ductile to brittle transition temperature (DBTT) at (673-773K [5]). Transmutation occurs through a sequence  
8 of neutron reactions and nuclide decays within the wall material. The beta decay of W after neutron  
9 absorption produces Re, Os, while Ta is produced via the beta decay of other W isotopes that have  
10 undergone neutron loss (n,2n) reactions. As a result, the material composition varies significantly with  
11 exposure time [6,7]. Spatial variations in the neutron flux and localised moderation also result in considerable  
12 variation in the composition of components within the reactor [7]. Understanding how these transmutation  
13 products accumulate is important for efforts to predict changes in physical properties of the plasma facing  
14 component during their lifetimes. Currently there is little reported data on tungsten samples with realistic  
15 transmutation product contents produced by neutron irradiation at reactor relevant temperatures.  
16  
17

18  
19  
20  
21  
22  
23  
24  
25 In this study, neutron irradiated single crystalline W was investigated using a combination of Atom Probe  
26 Tomography (APT), High Angle Annular Dark Field (HAADF) imaging and Energy Dispersive X ray (EDX)  
27 mapping. Pure W samples were irradiated at the High Flux Reactor (HFR) at Petten for a total of 208 full  
28 power days at  $900^{\circ}\text{C}$ , as part of the Extremat II program, resulting in a total dose of 1.67dpa. For more  
29 information on the irradiation, the reader is referred to [8]. APT and TEM samples were prepared using a  
30 Focussed Ion Beam (FIB) lift out technique with a FEI Helios FIB-SEM, at the Materials Research Facility  
31 (MRF) CCFE [9]. For APT, the lifted out material was mounted onto silicon posts using Pt deposition and  
32 sharpened using a series of successively smaller annular milling patterns (beam currents 1.5nA – 40pA,  
33 acceleration voltage of 30kV) until the tip of the needle was between 50-100nm in diameter. A final polishing  
34 stage using an acceleration voltage of 2kV and a beam current of 300pA was used to remove approximately  
35 500nm of material to minimise any Ga implantation. For TEM, thinning was carried out at 30kV using  
36 successively lower beam currents, with a final thinning stage at 2kV, using a 200pA beam current at  $45^{\circ}$  until  
37 the lift out was transparent under a 10kV electron beam. Flash polishing was carried out in a 0.5%NaOH  
38 solution with a 10V, 50ms applied voltage. The samples were then cleaned in methanol and plasma cleaned  
39 for 120s using a Fischione 1020 plasma cleaner.  
40  
41  
42  
43  
44  
45  
46  
47

48  
49 APT was carried out in laser mode in a CAMECA LEAP<sup>®</sup> 5000 XR using a sample temperature of 55K, a  
50 laser pulse energy of 125pJ and a pulse frequency of 100kHz. The data was reconstructed and visualised  
51 using IVAS<sup>®</sup> 3.8.2. Figure 1a shows a portion of the mass spectrum obtained during APT with the peaks  
52 ranged as W, Ta, Re and Os highlighted, compared in figure 1b with the expected isotopic distribution of  
53 these elements, based on the natural abundances. During conventional reconstruction of data, the ratios of  
54 the peaks can be expected to follow the natural isotopic distribution, and can therefore be used to identify  
55 peak species and prevent erroneous ranging. However, in the samples studied here, the Re, Os and Ta are  
56 introduced artificially through transmutation and therefore do not follow the natural abundances. Based on  
57 comparison with FISPACT II [10] modelling (see table 1), and to EDX measurements carried out elsewhere  
58 by Klimenkov et al. [8], the assumption was made that the overlaps between Re, Os and Ta in the 61.67, 62,  
59  
60  
61  
62  
63  
64  
65

62.33 and 62.67 Da peaks were negligible and could be ranged as  $^{185}\text{Re}$ ,  $^{186}\text{W}$ ,  $^{187}\text{Re}$  and  $^{188}\text{Os}$  respectively. The entire Os signal used for this analysis came from the  $^{188}\text{Os}$  peak. This resulted in a bulk sample composition of W 1.20±0.11at.%Re 0.11±0.05at.%Os 0.03±0.01at.%Ta in good agreement with FISPACT modelling and EDX measurements which predicted W-1.4Re-0.1Os [8]. The uncertainty on the composition was estimated by comparing the result when peaks in the mass spectrum were ranged background-to-background and at the full width half maximum (FWHM), rather than taking the counting uncertainty which is much lower. **The results presented in table 1 indicate that FISPACT-II is a very reliable approach to simulating the transmutation of these alloys and provides a robust description of post irradiation composition.**

Figure 2 shows one of the APT samples analysed, reconstructed using an initial tip radius estimated by IVAS and using the voltage profile to determine the tip diameter as a function of depth. W atoms are shown in blue, Re in red and Os in green. Ta atoms are not shown as they are present only in very small numbers and were uniformly distributed throughout the specimen. Red and green iso-concentration surfaces are displayed for 5at.%Re and 1.5at.%Os respectively. There is clear segregation of Re and Os from the matrix and a substantial increase in the Re and Os concentration compared to the nominal concentration. Furthermore, Os is seen to be located at the centre of the larger Re precipitates.

**Table 1:** Comparison between the calculated distribution of isotopes and the measured distribution from APT produced by transmutation, as a percentage of the total number of atoms in the simulation and the total number of ranged ions in APT. The 11 most commonly occurring isotopes from the simulation are included, isotopes with lower concentration could not be separated from background using APT. Uncertainty on the concentration from APT is estimated by comparing different ranging approaches, background-to-background and at FWHM. A further 60+ isotopes were predicted in lower concentrations. FISPACT-II calculations were performed with appropriate self-shielding and correctly simulated neutron spectra for the EXTREMAT-II experiment and used TENDL-2015 nuclear data as input - see [7] for further details. The simulations even accounted for a change in location of the sample after 60 days of exposure. The results also account for the approximately 6 years of decay cooling between the end of irradiation and APT/EDX measurements.

Nuclide	FISPACT-II %	APT %
$^{184}\text{W}$	30.97	30.98 ± 0.20
$^{186}\text{W}$	26.99	26.78 ± 0.17
$^{182}\text{W}$	25.35	25.39 ± 0.22
$^{183}\text{W}$	15.02	14.84 ± 0.02
$^{187}\text{Re}$	1.34	1.05 ± 0.07
$^{188}\text{Os}$	0.10	0.10 ± 0.03
$^{180}\text{W}$	0.10	0.13 ± 0.02
$^{185}\text{Re}$	0.09	0.12 ± 0.04
$^{181}\text{Ta}$	0.02	0.03 ± 0.01
$^{186}\text{Os}$	0.01	Undetectable
$^{189}\text{Os}$	2.1 x10 <sup>-3</sup>	Undetectable

33dpa ion implantations at 573 and 773K have been conducted by Xu et al. in both binary W-2Re [4,11] and ternary W-2Re-1Os [11]. Their APT of these samples showed both the formation of Re rich clusters in the binary alloy and smaller Os rich clusters in the ternary. At 773K, small Os clusters were found to form



preferentially to Re, inhibiting the formation of Re clustering seen in the binary alloy, but with some Re attachment to the Os cluster.

Xu et al.'s work used ion irradiation as a surrogate for neutrons, which is a commonly used technique of introducing point defects into a material. This results in a far higher dose rate than seen in the neutron irradiated specimens studied here. In the case of the neutron irradiated samples, the defects have more time to migrate through the material and drag Os and Re. The results shown in figure 2 are consistent with the initial formation of a small Os cluster due to its higher binding energy to vacancies [12], which are known to form voids in neutron irradiated W [13]. Os segregation to grain boundaries and dislocations within tungsten was observed in early atom probe field ion microscopy work in the 1980's by Eaton and Nordén [14]. These Os clusters could act as a nucleation sites for the Re, which forms around the Os over time.

Void formation is known to occur in neutron irradiated W samples [13]. The impact of voids on the reconstruction of APT data is not fully understood, with some studies showing an apparent decrease in density that the authors interpret as the location of a void [15], whilst others indicate an apparent increase in the density due to the effects of a higher atomic evaporation at the periphery of the void and the focussing of the evaporated ion flight path [16,17]. By binning a region of the APT data into voxels using IVAS, a map of the specimen density in number of atoms per cubic nanometre can be produced, as shown in figure 3, by counting the number of atoms per voxel. A density map of the sample is shown in figure 3.a. There is a clear increase in the atomic density associated with the Re/Os clusters, the location of which are shown in figures 3.b and 3.c. Re has a lower evaporation field threshold and therefore high localised concentrations of Re could contribute to an increase in the observed density, however, this is typically lower than observed here. At the centre of the Re clusters shown in figure 3, the measured atomic density from APT is  $\sim 80$  atoms  $\text{nm}^{-3}$ . Bulk density measurements from APT are lower than the theoretical atomic density as APT does not detect 100% of the atoms. Despite this, the measured core density of the clusters is above the theoretical maximum atomic density of tungsten ( $62.84$  atoms  $\text{nm}^{-3}$ ) which is bcc and has a lattice parameter of  $3.1690\text{\AA}$ . This indicates that the clusters observed may be decorating voids which have the potential to produce this apparent increase in atomic density through trajectory aberration at the periphery of the void.

A combination of STEM HAADF and EDX mapping was carried out using a JEOL 200F ARM TEM, operating at 200kV and equipped with a JEOL Centurion EDX detector and a Gatan GIF. This was used to map the Re/Os concentration around dark spots in the HAADF image, which suggested the presence of voids. This was confirmed using Electron Energy Loss Spectroscopy (EELS) thickness measurements which showed the darker regions corresponded to a smaller sample thickness. EDX mapping was carried out on a region containing several voids which indicated the presence of Os and Re. The  $L\alpha$  peaks of Re, Os and W at 8.651eV, 8.910eV and 8.396eV were used to identify the atomic species. To reduce drift artefacts, the EDX was performed using multiple scans and a fast acquisition time.

A large population of voids was seen as well as smaller isolated bright spots in the HAADF image which were confirmed to be Re clusters using EDX mapping. Figure 4a shows a high magnification STEM HAADF image of a single void, confirmed to be a void using EELS line scans. Figures 4c, 4d and 4e show the W, Re and Os EDX maps of this area respectively. The Re appears to form a shell around the void with a precipitate offset to one side. Within this offset region is the Os precipitate. This is further evidenced by the APT data shown in figure 3c, where the central precipitate also shows an offset of the osmium rich region compared to the rhenium one.

1 The samples studied here represent the upper right portion of the structural diagram for the irradiated W  
2 microstructure as a function of dose and temperature produced by Hasegawa et al., which also have a  
3 microstructure dominated by voids [13]. Several neutron irradiation studies using fission neutron spectra  
4 have shown the formation of lenticular chi or sigma phase precipitates [18–21], however the post irradiation  
5 compositions of these samples is typically much higher than can be expected for fusion wall materials.  
6 Klimenkov et al. found Re/Os precipitates in the same material as studied here but only a small Os signal  
7 was observable. Re was shown to segregate to the grain boundaries in these samples but Os was not  
8 shown in this region [8]. By using APT here we are able to separate and map both Re and Os. Hwang et al.  
9 saw both Re and Os within small clusters in neutron irradiated (initially pure) W, transmuted to 4%Re-  
10 0.8%Os using APT. Os was present within the clusters but was not seen to form a core region of Os as seen  
11 here. The regular separation of the clusters indicated that they had formed on a void lattice, however there  
12 was no direct indication from atomic density mapping if voids had affected the internal structure of the  
13 clusters [19].  
14  
15  
16  
17

18 The mechanism through which this precipitation occurs is not clear. **The microstructure presented here is**  
19 **dominated by voids unlike in Xu et al's work [4,11] in which there was no indication of voids**  
20 **associated with the Re clusters. This could be due to the lower irradiation temperature and higher**  
21 **dose rate used in Xu et al's ion implantation, which may have suppressed void growth, making them**  
22 **more difficult to detect using APT.** Kinetic Monte Carlo (KMC) simulations of W-Re have shown the  
23 importance of the mixed interstitial for precipitation. Unlike the self-interstitial, which is confined to 1D motion,  
24 a mixed W-Re interstitial is able to move through a series of rotations and translations, to effectively migrate  
25 in 3D [12,22]. The increased mobility of the mixed interstitial facilitates the nucleation and growth of Re  
26 clusters, despite the predicted solid solution [22]. Models such as this rely on a simple Ising like model, which  
27 uses pairwise interactions to construct process rates, such as migration and rotation. The pairwise binding  
28 energy of the di-vacancy in W is close to zero, which results in an underestimation of the role of vacancies  
29 within this model. The use of a Cluster Expansion (CE) Hamiltonian can overcome this problem by  
30 considering interactions between 3 or more atoms/defects at once. At high temperatures (stage IV recovery  
31 [3]) vacancies become much more mobile and play a more significant role in microstructural evolution. MC  
32 simulations using a CE Hamiltonian have predicted the precipitation of Re in W alloys through the formation  
33 of small voids, decorated with Re atoms [23], however, the role that Os plays in this process is still being fully  
34 investigated. **More modelling work is required to understand the operative mechanisms during**  
35 **irradiation, particularly when in the presence of Os.**  
36  
37  
38  
39  
40  
41  
42  
43  
44  
45

46 In summary, neutron irradiation of pure single crystal W led to the formation of a W -  $1.20 \pm 0.11$  at.%Re -  
47  $0.11 \pm 0.05$  at.%Os -  $0.03 \pm 0.01$  at.%Ta alloy. **Comparison between FISPACT-II results and APT**  
48 **demonstrated that FISPACT-II is a very reliable way of quantifying transmutation in neutron**  
49 **irradiated material.** This sample is of particular importance for the study of the evolution of materials in a  
50 fusion environment as it has undergone a high temperature irradiation and has a transmutation induced  
51 composition similar to a W first wall component after several years of operation in a reactor. Under irradiation  
52 APT showed clusters of mainly of Re, with Os at the centre, which were on average 8nm in diameter.  
53 Analysis of the reconstruction indicated that voids within the clusters may have contributed to an artefact  
54 which could have disrupted the inner structure of the clusters seen in APT. Additionally, STEM HAADF and  
55 EDX of the sample confirmed that clustering can occur through the decoration of a void and likely contributed  
56 to the structure seen in the APT analysis. This work shows the importance of using transmutation  
57 compositions when ranging the APT data produced on samples containing non-natural abundances of  
58  
59  
60  
61  
62  
63  
64  
65

elements, and raises questions about the validity of using ion implantation which typically does not facilitate void formation. New neutron irradiation campaigns, **and more modelling studies, which include the full range of transmutation products** are needed to fully understand the impact transmutation will have in a reactor. In addition, care must be taken if only using APT to study precipitation in tungsten and other alloys that can undergo void formation during irradiation, as artefacts can be induced into data reconstruction due to trajectory aberrations around the void. Further work is need on methods used to reconstruct the APT data to include the voids seen in HAADF and account for localised magnification.

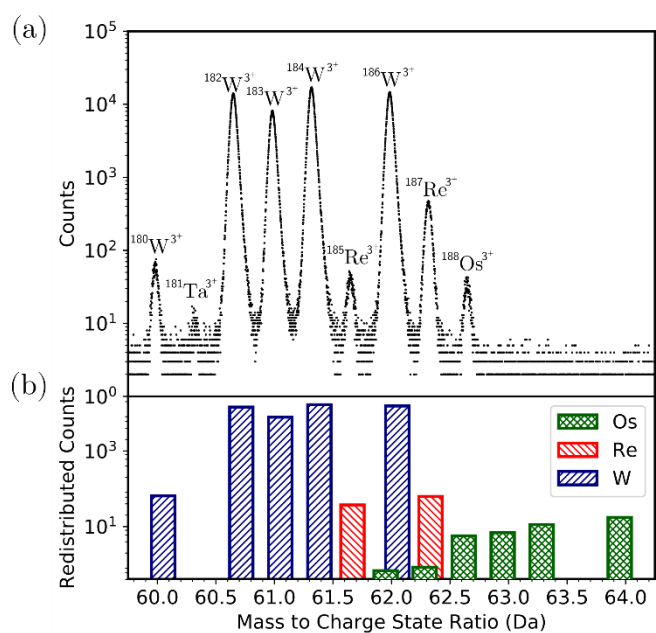
#### Acknowledgements

Funding: This work was supported by the UK Engineering and Physical Sciences Research Council [EP/N509711/1] and the Culham Centre for Fusion Energy, United Kingdom Atomic Energy Authority through an Industrial CASE scholarship, [Project Reference Number 1802461].

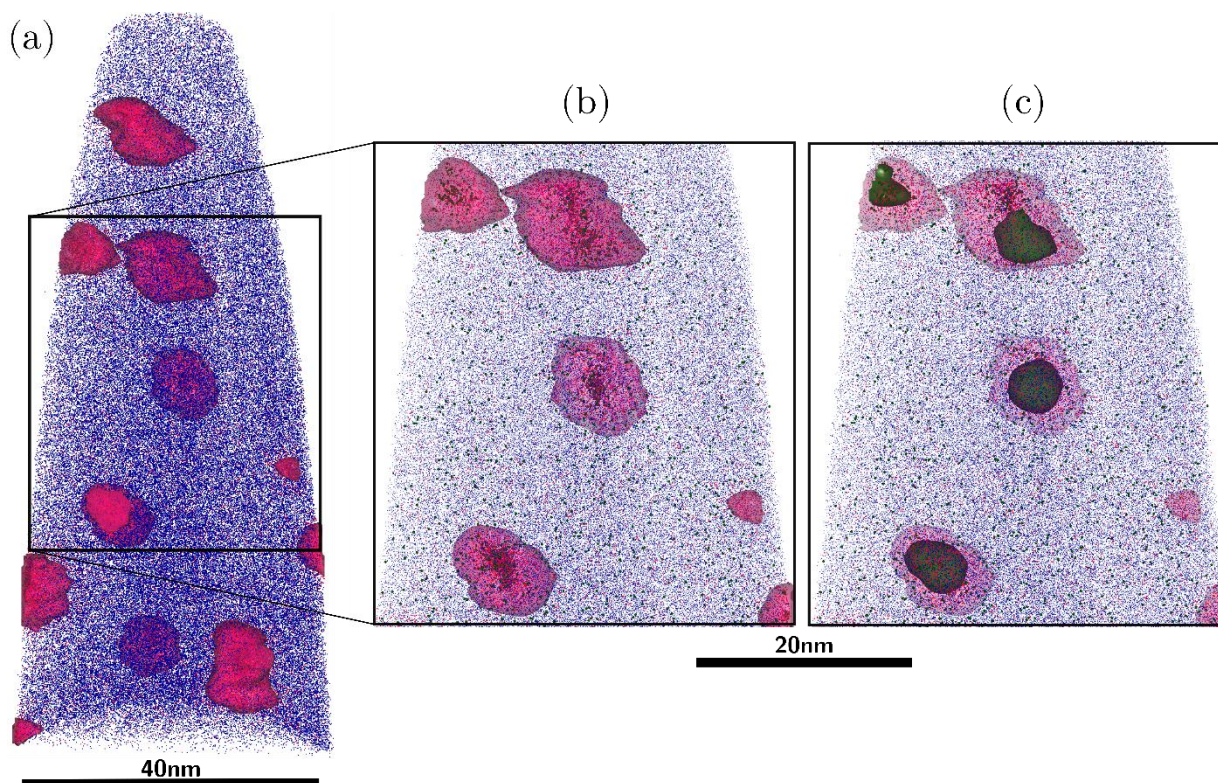
Atom Probe Tomography was carried out at the Oxford Materials Atom Probe Group. The authors would also like to thank Steve G. Roberts for supplying the samples used in this study. The research used UKAEA's Materials Research Facility, which has been funded by and is part of the UK's National Nuclear User Facility and Henry Royce Institute for Advanced Materials. M.J.L. would like that thank Andy London and Enrique Martinez for their involvement in discussions on this topic.

- [1] J. Matejicek, M. Battabyal, P. López-Ruiz, M. Celino, J. Linke, J.G. van der Laan, C.S. Becquart, H. Li, S. Wahlberg, W.W. Basuki, K. Heinola, J. Riesch, C. Linsmeier, T. Weitkamp, J. Aktaa, L. Romaner, M. Rosiński, S. Lindig, J. Brinkmann, J. Opschoor, S. Giusepponi, A. Hoffmann, W. Schulmeyer, R. Pippan, H. Boldyryeva, L. Veleva, D. Nguyen-Manh, F. Koch, T. Palacios, T. Ahlgren, D.E.J. Armstrong, L. Ciupinski, M.A. Yar, H. Traxler, J. Reiser, A. Muñoz, C. Domain, C. García-Rosales, M.R. Gilbert, A. Ureña, M. Sanchez, S.L. Dudarev, M. Balden, M. Rieth, M. Muhammed, M. Muzyk, N. Holstein, T. Weber, D. Blagoeva, J.H. You, H. Maier, A. De Backer, N. Baluc, J. Gibson, G. Pintsuk, M.-F. Barthe, S.M. Gonzalez de Vicente, J.B. Correia, S. Wurster, K. Nordlund, B. Gludovatz, M. Walter, H. Greuner, W. Krauss, N. Ordás, S.G. Roberts, E. Gaganidze, T. Höschen, T.P. Mishra, S. Antusch, A. Zivelonghi, *J. Nucl. Mater.* 432 (2012) 482–500.
- [2] D.E.J. Armstrong, P.D. Edmondson, S.G. Roberts, *Appl. Phys. Lett.* 102 (2013) 251901.
- [3] R.G. Abernethy, *Mater. Sci. Technol.* 0836 (2016) 1–12.
- [4] A. Xu, D.E.J. Armstrong, C. Beck, M.P. Moody, G.D.W. Smith, P.A.J. Bagot, S.G. Roberts, *Acta Mater.* 124 (2017) 71–78.
- [5] F. Ferroni, X. Yi, K. Arakawa, S.P. Fitzgerald, P.D. Edmondson, S.G. Roberts, *Acta Mater.* 90 (2015) 380–393.
- [6] M.R. Gilbert, J.-C. Sublet, *Nucl. Fusion* 51 (2011) 043005.
- [7] M.R. Gilbert, J.C. Sublet, S.L. Dudarev, *Nucl. Fusion* 57 (2017) 044002.
- [8] M. Klimenkov, U. Jäntschi, H.C. Schneider, D.E.J. Armstrong, J. Gibson, *Nucl. Mater. Energy* 9 (2016) 480–483.
- [9] Materials Research Facility: Developing Materials for the Nuclear Industry, 2018.

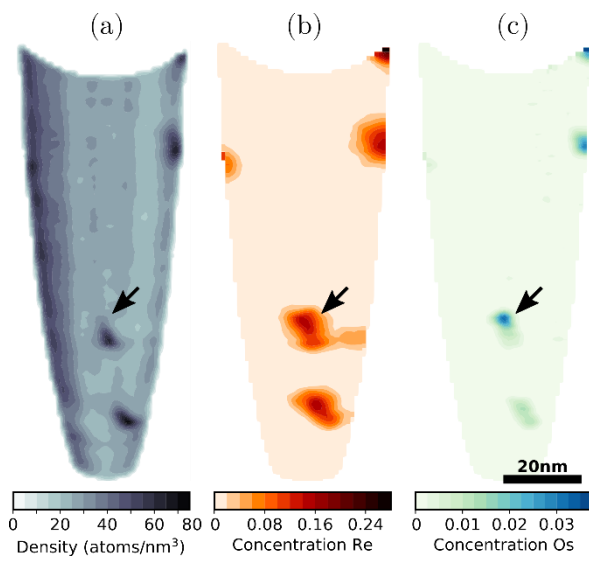
- 1  
2  
3  
4  
5  
6  
7  
8  
9  
10  
11  
12  
13  
14  
15  
16  
17  
18  
19  
20  
21  
22  
23  
24  
25  
26  
27  
28  
29  
30  
31  
32  
33  
34  
35  
36  
37  
38  
39  
40  
41  
42  
43  
44  
45  
46  
47  
48  
49  
50  
51  
52  
53  
54  
55  
56  
57  
58  
59  
60  
61  
62  
63  
64  
65
- [10] J.C. Sublet, J.W. Eastwood, J.G. Morgan, M.R. Gilbert, M. Fleming, W. Arter, Nucl. Data Sheets 139 (2017) 77–137.
  - [11] A. Xu, C. Beck, D.E.J. Armstrong, K. Rajan, G.D.W. Smith, P.A.J. Bagot, S.G. Roberts, Acta Mater. 87 (2015) 121–127.
  - [12] T. Suzudo, M. Yamaguchi, A. Hasegawa, Model. Simul. Mater. Sci. Eng. 22 (2014).
  - [13] A. Hasegawa, M. Fukuda, S. Nogami, K. Yabuuchi, Fusion Eng. Des. 89 (2014) 1568–1572.
  - [14] H.C. Eaton, H. Nordén, Scr. Metall. 17 (1983) 1043–1046.
  - [15] M.K. Miller, L. Longstreth-Spoor, K.F. Kelton, Ultramicroscopy 111 (2011) 469–472.
  - [16] P.J. Birdseye, D.A. Smith, G.D.W. Smith, J. Phys. D. Appl. Phys. 7 (1974) 1642–1651.
  - [17] M. Dagan, L.R. Hanna, A. Xu, S.G. Roberts, G.D.W. Smith, B. Gault, P.D. Edmondson, P.A.J. Bagot, M.P. Moody, Ultramicroscopy 159 (2015) 387–394.
  - [18] X. Hu, T. Koyanagi, M. Fukuda, Y. Katoh, L.L. Snead, B.D. Wirth, J. Nucl. Mater. 470 (2016) 278–289.
  - [19] T. Hwang, A. Hasegawa, K. Tomura, N. Ebisawa, T. Toyama, Y. Nagai, M. Fukuda, T. Miyazawa, T. Tanaka, S. Nogami, J. Nucl. Mater. 507 (2018) 78–86.
  - [20] X. Hu, C.M. Parish, K. Wang, T. Koyanagi, B.P. Eftink, Y. Katoh, Acta Mater. 165 (2019) 51–61.
  - [21] X. Hu, N.A.P.K. Kumar, T. Koyanagi, Y. Katoh, L.L. Snead, T. Hwang, L.M. Garrison, J. Nucl. Mater. 490 (2017) 66–74.
  - [22] C.H. Huang, L. Gharaee, Y. Zhao, P. Erhart, J. Marian, Phys. Rev. B 96 (2017) 1–17.
  - [23] J.S. Wróbel, D. Nguyen-Manh, K.J. Kurzydłowski, S.L. Dudarev, J. Phys. Condens. Matter 29 (2017) 1–5.



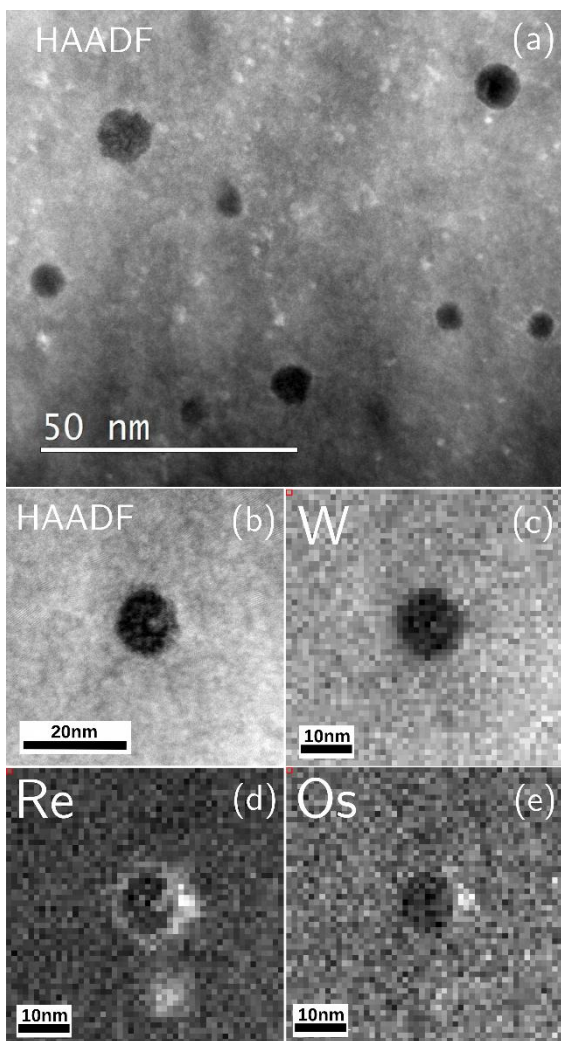
**Figure 1:** (a) Sub-section of a representative mass spectrum generated by APT analysis. Labels indicate the isotopic assignment for all reconstructed ions contributing to that peak, respectively. (b) Expected relative size of peaks if isotopes were in natural abundances. Total counts from APT for each element in the mass spectrum shown in (a) have been summed and multiplied by the natural abundance of that element.



**Figure 2:** APT reconstruction of neutron irradiated sample. (a) For visual clarity, only 1.3% of total detected W ions are shown in blue, all detected Re and Os ions are shown in red and green respectively. Red surfaces represent Re iso-concentration surfaces plotted at 5at.%Re. (b) Higher magnification of the region of interest showing Re iso-concentration surfaces at 5at.%Re and individual Os atoms as green squares. (c) Region of interest showing transparent Re iso-concentration surfaces and green Os iso-concentration surfaces at 1.5at.%Os



**Figure 3:** 5nm cross section of the atom map shown in figure 2. (a) Density contour in number of atoms per cubic nm (b) Re concentration contour (c) Os concentration contour. Arrow indicates central precipitate with Os rich zone offset from higher density region



**Figure 4:** (a) Lower magnification and (b) high magnification High Angle Annular Dark Field (HAADF) images of voids in neutron irradiated single crystal W. EDX Maps of (c) W (d) Re and (e) Os showing the decoration of central void