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Schneider, B. Kurzan, U. Stroth, the ASDEX
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Empirical study of gradient lengths ratio η_e in the near SOL region in ASDEX Upgrade tokamak

H. J. Sun^{1,2}, E. Wolfrum², T. Eich², A. Kallenbach², P. Schneider², B. Kurzan², U. Stroth^{2,3} and the ASDEX Upgrade Team

¹ United Kingdom Atomic Energy Authority, Culham Centre for Fusion Energy, Culham Science Centre, Abingdon, Oxon, OX14 3DB, UK

² Max Planck Institute for Plasma Physics, Garching, Germany

³ Physik-Department E28, Technische Universität München, Garching, Germany

E-mail: sunhj@ipp.mpg.de

Abstract. A survey of the correlation between temperature and density gradient lengths (λ_{Te} and λ_{ne}) has been performed for the near SOL regions in the ASDEX Upgrade tokamak. In common with previous studies, the near SOL temperature and density gradient length ratio, η_e , is found to be close to unity across a wide range of plasmas. However, it is found that it is possible for both the temperature and density gradients to vary separately without varying the other greatly. The empirical results show that the correlation between temperature and density profiles in the near SOL is related to the divertor regimes. In the linear/sheath limited regime, the density and temperature profiles are correlated as $T_{e,u} \propto n_{e,u}^{-2/3}$, giving $\eta_e = \lambda_{ne}/\lambda_{Te} \approx 2/3$. In the high recycling/ conduction limited regimes, which is the most common case for ASDEX Upgrade plasma, λ_{ne} and λ_{Te} are correlated due to a common dependence on q_{95} or, equally, I_p . A weak dependence of η_e on separatrix temperature is observed in this regime. Since the separatrix temperature depends only weakly on the power density flowing through the SOL and the variation is very small, a small variation of η_e is observed. For ASDEX Upgrade, it lies between 1 and 2 for most of discharges in the conduction limited regime. In the detached regime, convection could be important for the cross field particle transport and filaments appear, may trigger a regime transition leading to strongly enhanced perpendicular particle fluxes, resulting in a very flat density profiles while temperature profile remain unaffected, resulting in very high η_e , for ASDEX Upgrade, up to 6.

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1 Introduction

The electron density and temperature decay lengths in the scrape-off layer (SOL), λ_{ne} and λ_{Te} , set crucial boundary conditions for the heat- and particle-handling in magnetically confined plasmas. These decay lengths control the penetration of recycling hydrogen neutrals and edge-produced impurities into the plasma. How edge temperature and density profiles correlate eventually determines the characteristics of edge pressure. The density and temperature gradient length ratio, $\eta_e = \lambda_{ne}/\lambda_{Te}$, is also important in transport physics and determining turbulence characteristics.

A previous study on ASDEX Upgrade (AUG), found a fairly constant η_e close to 2 from pedestal across to the SOL for a series of Type I ELMy H-mode discharges [1]. For a larger data set of AUG ELMy H-modes, η_e in the pedestal region was observed to be $\eta_e = 1.9 \pm 0.6$ [2]. These studies [1] [2] were carried out with graphite plasma facing components (PFCs). A more recent study, for a database of AUG Type I ELMy H-mode plasma with tungsten PFCs, found η_e to range from $\eta_e = 1 - 2$ in both the pedestal and SOL regions, with an average value of $\eta_e \approx 1.4$ [3]. A study of η_e in the SOL for a data set of JET ELMy H-mode plasma found an average value of $\eta_e \approx 1.5$ [4]. The study of a large data set of DIII-D plasmas, including different high confinement plasma scenarios, found that η_e in the middle of the pedestal ranged from $\eta_e = 1 - 3$ [5]. For a Tore Supra data set of Ohmic limiter discharges [6], the temperature gradient length in the SOL region is found to be generally larger than the density gradient length, with a rough relation of $\lambda_{T_e} \approx 1.4(\pm 1)\lambda_{n_e}$, giving $\eta_e \approx 0.7$.

Combining previous studies, η_e is observed to be not far from unity in the edge plasma, suggesting that the density and temperature profiles may be closely coupled. However, η_e varies in the range $\eta_e \approx 0.4 - 3$, across the single machine datasets and between different machines. The confidence intervals for the measurements of SOL gradients for some of the studies are relatively large, resulting in confidence intervals in η_e comparable to the observed scatter. Hence, it is useful and important to document the empirical parametric dependences of η_e and find out whether the temperature and density profiles are always coupled.

Recent improvements of the AUG Thomson Scattering (TS) system [7] enable the detailed study of characteristics in the near SOL profiles. The purpose of this paper is to study the relation between edge T_e and n_e profiles in the SOL regions. The rest of this paper is organized as follows: The experimental methods and datasets are briefly introduced in section 2. In section 3, at first, the dependences on plasma global parameters are tested; then, the correlations with local parameters are checked. In section 4, the examples where the temperature and density profiles are decoupled are presented. In section 5, the results are summarized and physics implications are discussed.

2 Experimental methods and database

In ASDEX Upgrade, electron temperature and density profiles can be obtained simultaneously for one discharge by a vertical TS system [3] [7]. The system is equipped with 4 Nd: YAG lasers at 1064 nm for core profiles and 6 lasers for the edge plasma. The system has a spatial resolution of 25 mm and temporal resolution of 80 Hz in the core and a spatial resolution of around 3 mm and temporal resolution of 120 Hz in the near SOL.

Since the 2014 AUG campaign, significant improvements of the TS system have enabled good measurements of the SOL temperature and density profiles [3]. The Type I ELMy H-mode discharges under attached divertor condition (36 in total) and L-mode attached discharges (27 in total) used in the present dataset are the same discharges as those used in Ref [8]. In H-mode plasma, during ELMs the thermocurrents measured in the divertor are strongly increased, which indicates the onset of an ELM. Before the ELM onset 1.5 ms are excluded as margin for possible errors in the ELM recognition and time delays between different diagnostics. The relevant period for profile synchronization is defined from -3.5 to -1.5 ms relative to the ELM onset time and also at least 4 ms after the previous ELM are excluded. However, the same method is not applicable for the inter-ELM profiles under the detached condition because the ELM frequency is too high. As can be seen in figure 1, both temperature (in red)

and density (in black) profiles normally exhibit a two-layer structure. Close to the separatrix, in the so-called near SOL, it has a steep exponential decay. Beyond this region, in the so-called far SOL, the profile has an exponential decay with much longer scale length. The temperature and density decay length $\lambda_{T_{e,u}}$ and $\lambda_{n_{e,u}}$ in this paper refer to the first e-folding length in the profiles, that is why they are also called midplane e-folding distance in some studies. The solid blue line fit to the log plot of near SOL data gives the gradient decay lengths in figure 1.

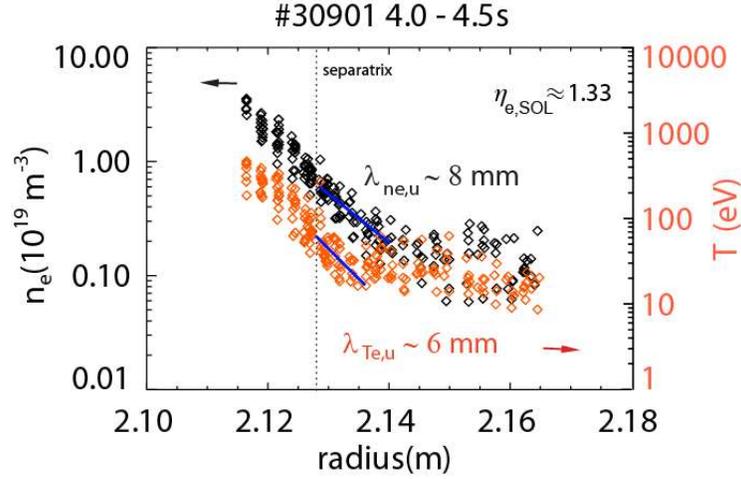


Figure 1. Log plot of electron temperature (in red) and density (in black) versus radius for AUG discharge #30901. The (solid blue) line fit to near SOL data gives the gradient decay lengths. The dashed line the separatrix indicates the separatrix position.

3 Experimental observations of the correlation between density and temperature profiles

3.1 General observation and dependences on global parameters

A previous study, with a database of Type I ELMy H-mode plasmas, found the near SOL density and temperature profiles are closely correlated with η_e in the range from $\eta_e = 1 - 2$, with an average value of $\eta_e \approx 1.4$ [3]. To increase the range of the plasma parameters for the parametric dependence study, L-mode discharges under attached divertor conditions are also included in this paper. Figure 2 shows a plot of $\lambda_{T_{e,u}}$ versus $\lambda_{n_{e,u}}$ in the near SOL for both H-mode and L-mode discharges in the present study. Most of η_e have values in the range 1–2, as shown by the solid black lines. However, there are data from L-mode discharges which lie below one, i.e. $\lambda_{T_{e,u}} > \lambda_{n_{e,u}}$. In the next section, it will be shown that these discharges associate with very low separatrix density. The previous studies from both IR camera and TS system measurements show that the SOL characteristic temperature and power decay lengths have a strong plasma current, I_p , dependence [9] [3] [8]. Detailed analysis in both H- and L-mode plasma from TS measurement suggests that a safety factor, q_{95} , dependence may be more appropriate to describe the temperature gradient length scaling [3] [8]. Do these global plasma parameters still play an important role in determining the gradient length ratio? Since this AUG dataset contains only a limited range in B_T and the majority of the dataset for both L-mode and H-mode is comprised of discharges around 2.5 T. Thus, I_p and q_{95} are strongly correlated. Since q_{95} is directly related to important SOL parameters, such as connection length, only the correlation with q_{95} will be discussed in the following sections.

Plots of η_e against q_{95} are shown in figure 3. For H-mode plasma, there is no correlation between η_e and q_{95} . For the L-mode plasma, for the discharges with $q_{95} < 6$, there is no correlation between η_e and q_{95} . However, η_e for the discharges with $q_{95} > 6$ ($I_p \approx 600$ kA) clearly drops below unity, giving a negative trend between η_e and q_{95} . These discharges with higher q_{95} are the same ones mentioned above which will be discussed more in the next section.

As observed in the previous study of SOL temperature profile [8], $\lambda_{T_{e,u}}$ has a strong q_{95} dependence, $\lambda_{T_{e,u}} \propto q_{95}$ for both H-mode and L-mode plasmas. If $\eta_e = \lambda_{n_{e,u}}/\lambda_{T_{e,u}}$ has different dependences on global parameters, this must be due to the density profile. As can be seen from figure 4(a) and (b), $\lambda_{n_{e,u}}$ also has a strong q_{95} dependence, $\lambda_{n_{e,u}} \propto q_{95}$, for H-mode plasma. For L-mode plasma with $q_{95} < 6$, there is a similar positive trend between $\lambda_{n_{e,u}}$ and q_{95} , while the discharges with $q_{95} > 6$ show no such correlation. The correlations between $\lambda_{n_{e,u}}$ and q_{95} are consistent with the observed relation between η_e and q_{95} . In particular, the lack of a dependence of η_e on q_{95} , for H-mode plasma and L-mode plasma with $q_{95} < 6$, results from a co-dependence of $\lambda_{n_{e,u}}$ and $\lambda_{T_{e,u}}$ on q_{95} .

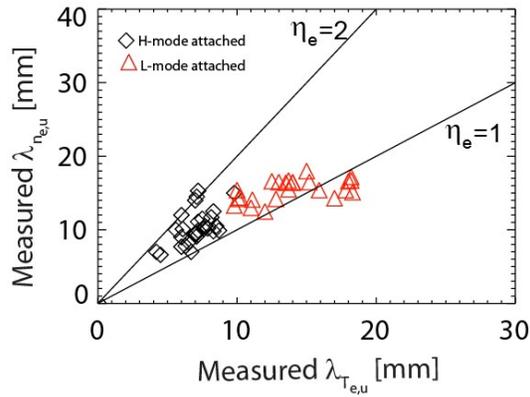


Figure 2. $\lambda_{n_{e,u}}$ versus $\lambda_{T_{e,u}}$ for both H-mode and L-mode discharges. Upper line corresponds to η_e of 2 and lower line to η_e of 1.

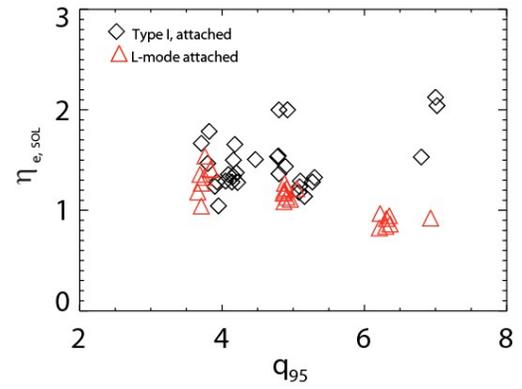


Figure 3. η_e against safety factor q_{95} in the near SOL region, for the AUG SOL dataset.

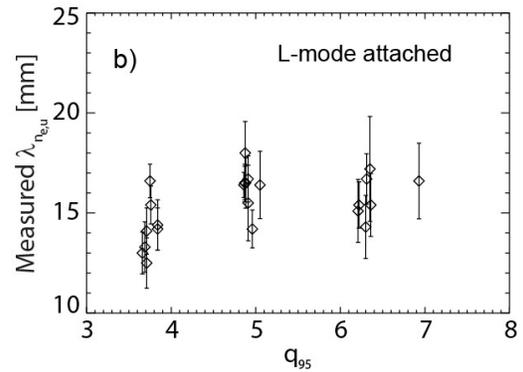
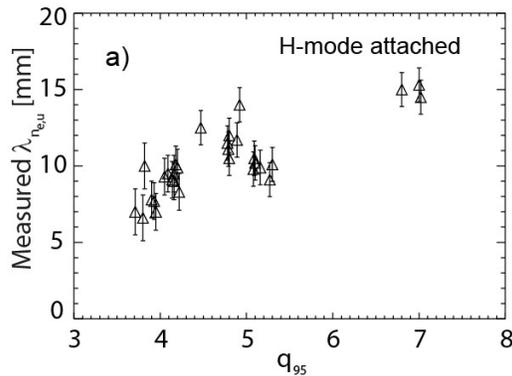


Figure 4. $\lambda_{n_{e,u}}$ versus q_{95} (a) for H-mode plasma; (b) for L-mode plasma.

3.2 Dependences on local parameters

In this section, the correlation between η_e and the separatrix local parameters will be examined. In Figure 5(a), for H-mode attached plasmas, η_e is seen not to correlate with $n_{e,sep}$. For example, η_e with $n_{e,sep} \approx 1.5 \times 10^{19} m^{-3}$ lies in almost the same range as that with $n_{e,sep} \approx 5.5 \times 10^{19} m^{-3}$. For L-mode plasma, the discharges with smallest $\eta_e \leq 1$ are those with lowest density. As discussed in the previous section, these discharges with the lowest density do not follow the generally observed positive trend between $\lambda_{n_{e,u}}$ and q_{95} . At low density, it is likely that the divertor is in the linear/sheath limited regime. In this regime, the upstream and target temperatures, T_u and T_t , are approximately the same and scale as $T_u \approx T_t \propto n_u^{-\frac{2}{3}}$ [10] [11], giving $\eta_e \approx \frac{2}{3}$, which is consistent with the values for these very low density discharges. In a previous study [12], η_e in the pedestal region is observed to lie mainly between 1 and 3 across the entire data set from AUG, JET and DIII-D, as shown in figure 7.43 in Ref [12]. However, there are several discharges that fall outside of this range, with very low pedestal density, $n_{e,ped} \approx 1 - 1.5 \times 10^{19} m^{-3}$, the gradient length ratio is also nearly $\eta_e \approx \frac{2}{3}$. This suggests that these discharges are also in the linear regime. A positive trend of higher separatrix temperature $T_{e,sep}$ being associated with a larger η_e in the near SOL region is seen, figure 5(b), similar to the observation for the pedestal region [12]. In Tore Supra [6], a rough relation, $\lambda_{Te} \approx 1.4 \lambda_{ne}$, is found between the density and temperature decay length in the SOL, equivalent to $\eta_e \approx 0.7$. Across this Tore Supra database, the separatrix temperature is about 30 eV. Thus, the lower η_e in the Tore Supra study appears consistent with the trend observed in figure 5(b). The dependence of η_e on $T_{e,sep}$ could be explained that, when cross-field transport remains the same, the reduction of parallel heat transport due to the lower temperature will broaden λ_{Te} , thus leading to a smaller η_e . Unlike the result from pedestal region, there is no obvious correlation between η_e and the SOL collisionality, $v_{SOL,e}^* \approx 10^{-1} n_{e,u} L / T_{e,u}^2$ [13], over the analyzed dataset, figure 5(c). In the study of the pedestal region [12], a negative correlation between η_e and the collisionality, $v_{ped,e}^*$, was identified for all analyzed data. The fact that there is correlation observed between η_e and local collisionality in the pedestal region, but not in the SOL region, could be due to the change of the particle transport channel in the SOL region.

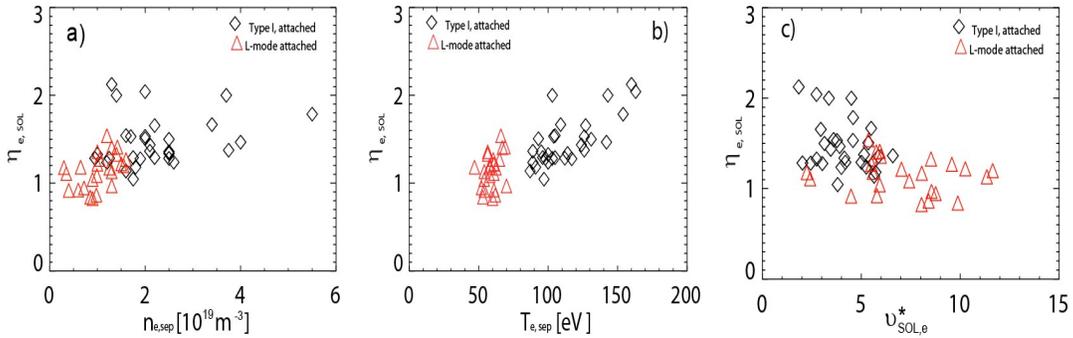


Figure 5. η_e in the near SOL plotted against a) separatrix density, $n_{e,sep}$, b) separatrix temperature, $T_{e,sep}$, (c) SOL collisionality, $v_{SOL,e}^*$, for the AUG SOL dataset.

3.3 Studies where temperature and density gradient lengths change separately

The results from previous sections show that the measured η_e in the near SOL region is not far from unity for both H-mode and L-mode plasmas. Even when the plasma transits from the L-mode regime to

the H-mode regime, the temperature and density profiles become steeper simultaneously and there is no sudden increase of η_e from L-mode to H-mode. This gives an impression that a coupling between the density and temperature profiles exists, in the sense that the gradient length of one cannot be changed alone. To study whether this is the case, this section presents a study of two plasma regimes which have been identified as having a relatively wider range in one of the gradient lengths whilst having a relatively small range in the other.

An alternative improved confinement regime to the H-mode is the so-called improved energy confinement regime (I-mode) [14] [15]. In I-mode plasma, an increase of the edge temperatures with a concomitant steepening of the edge gradients builds a temperature pedestal similar to that of the H-mode, but the density profile almost does not change, with an edge density profile very similar to that of the L-mode. Thus, the transition to I-mode plasma provides a good way to investigate the correlation of the temperature profile with the gradient length ratio, for constant plasma density.

An example of an AUG I-mode discharge, #30865, is described in detail in Ref [15]. In this discharge, the ECRH heating power was increased every 500ms, while the density remained almost constant. Around $t \approx 3.18$ s, the electron temperature and pressure gradients slightly steepen and the confinement increases slightly at constant heating power. This spontaneous confinement transition is considered as a transition to a weak I-mode. With additional heating power, at $t = 3.5$ s, $H_{98(y,2)}$ increases above 0.6 and a clear I-mode develops: the electron temperature gradients strongly increase, while the density gradient remains almost unchanged. Detailed related physics for these three phases are presented in Ref [15]. Here, only the change of temperature and density profiles in L-mode and strong I-mode phases are discussed to demonstrate how η_e can change only due to a change of the temperature profile. The evolution of the temperature and density profiles from L-mode to strong I-mode is illustrated in figure 6 and figure 7. The L-mode phase has $\eta_e \approx 1.27$, in the I-mode regime, η_e increases, to $\eta_e \approx 2$. This I-mode discharge demonstrates how changes in the temperature profile affect the gradient length ratio. As the temperature profile becomes much steeper, while the density profile remains almost unchanged, the gradient length ratio η_e increases.

Assuming Spitzer-Harm transport along the field lines, the calculated separatrix temperature, $T_{e,sep}$ is about 50 eV for the L-mode phase and around 80 eV in the strong I-mode phase. Although similar positive trends between η_e and local electron temperatures are observed in the I-mode regime, given the same $T_{e,sep}$, η_e in the I-mode plasma is higher than that in typical L- and H-mode plasma. This could be due to the fact that in I-mode, a weakly coherent mode, which is located very close to the separatrix, increases dominantly particle transport [15] [16] [14].

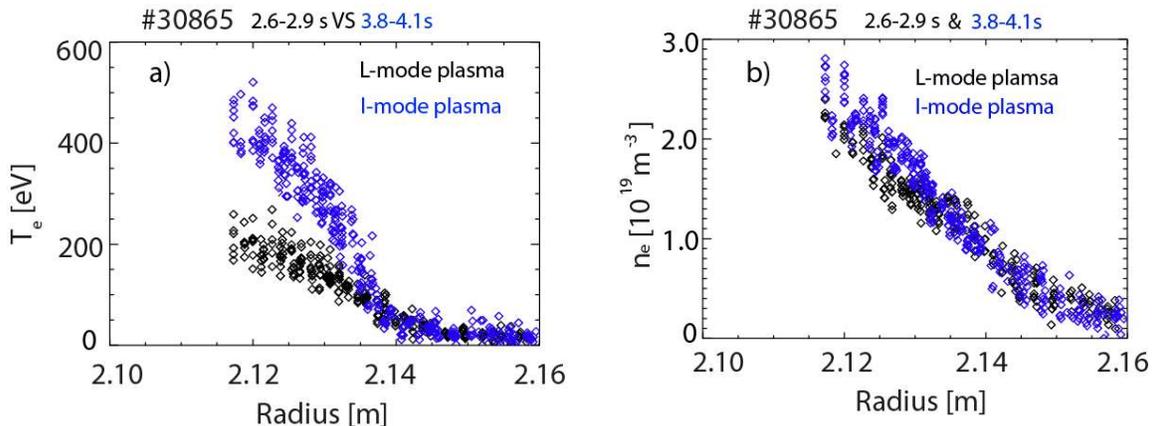


Figure 6. (a) electron temperature T_e and (b) electron density against major radius at the midplane for the L-mode (black), and strong I-mode (blue) phases of AUG discharge #30865.

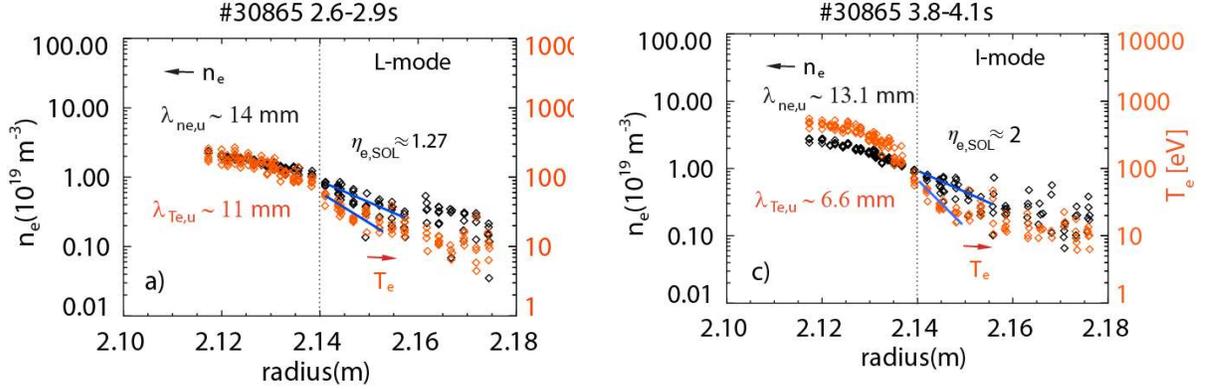


Figure 7. Log plot of electron temperature and density versus radius in the edge region for L-mode plasma (black), I-mode plasma (red) of the same AUG discharge as figure 6. The black dashed lines indicate the position of separatrix.

The next case is that of the transition of a H-mode plasma from attachment to detachment and then its approach to the H-mode density limit. In this case, the temperature gradient length is relatively constant compared to the change in the density gradient length. Figure 8 shows the characteristic time traces of an H-mode density limit discharge. Around $t \approx 2$ s, the plasma transits from the L-mode regime to the H-mode regime, due to the increase of heating power. During the H-mode regime, due to fueling without the cryo-pump, both the divertor neutral density and the electron density increases dramatically, around $t \approx 3$ s, the divertor is detached, and the density continues increasing until it reaches the H-mode density limit around $t \approx 4$ s, when it triggers the H-L back transition shortly before the plasma collapses. Temperature and density profiles from four different phases of this discharge are shown in figure 9. In the L-mode regime, with $n_{e,sep} \approx 1.5 \times 10^{19} \text{ m}^{-3}$, the temperature and density profiles have very similar gradient lengths, giving $\eta_{e,SOL} \approx 1$. After transition into the H-mode regime, both temperature and density gradient lengths decreases, while λ_{T_e} decreases more than λ_{n_e} , η_e increases to $\eta_{e,SOL} \approx 1.33$, and separatrix density also increases, $n_{e,sep} \approx 3.1 \times 10^{19} \text{ m}^{-3}$. By the H-mode regime under detached divertor region, $t = 3.0-3.2$ s, since the divertor neutral pressure continues to increase, the separatrix density has increased correspondingly [17], $n_{e,sep} \approx 4.4 \times 10^{19} \text{ m}^{-3}$, and density profile becomes much flatter, η_e increases much more, $\eta_{e,ped} \approx 4.4$. Just before the H-L back transition, $t = 3.8-3.95$ s, $n_{e,sep}$ has increased up to $n_{e,sep} \approx 6 \times 10^{19} \text{ m}^{-3}$, the pedestal top density also increases, whilst the pedestal and SOL temperature have decreased. Due to the drop of $T_{e,sep}$, the temperature profile in the SOL also becomes slightly flatter, $\lambda_T \approx 11 \text{ mm}$. In this phase of the discharge, η_e is around 6. As can be seen from the evolution of profiles, the significant increase of η_e is caused primarily by the flattening of density profile. This feature of density profile flattening in the SOL region under detached divertor condition was observed in earlier studies of L-mode plasmas, where the phenomenon

has been called the L-mode high density transition (HDT) [18]. When the Greenwald density fraction f_{GW} is larger than a certain value, the density profile in the SOL changes: flattening its gradient and giving rise to a ‘shoulder’. In a recent study [19], filaments appear as a result of the interchange instability which enhances perpendicular particle transport, as high collisionality in the SOL may trigger a regime transition leading to strongly enhanced perpendicular particle fluxes. It seems that these convective structures do not significantly enhance the heat transport, as shown in figure 9(c), suggesting that conductive transport remains the dominant cross-field heat transport mechanism. Only when the separatrix temperature decreases because of the increase of density does the slowing of parallel transport flatten the temperature profile in near SOL, figure 9(d).

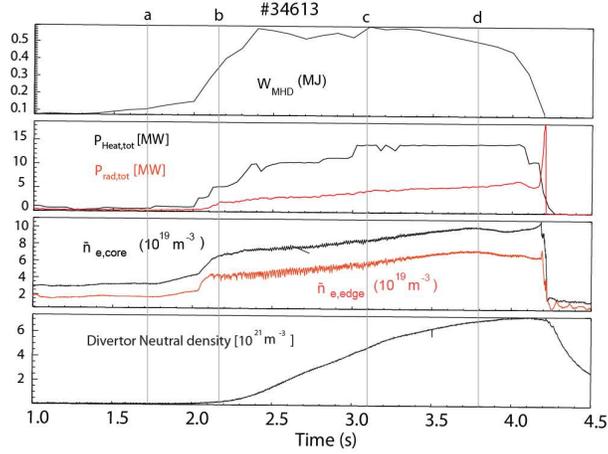
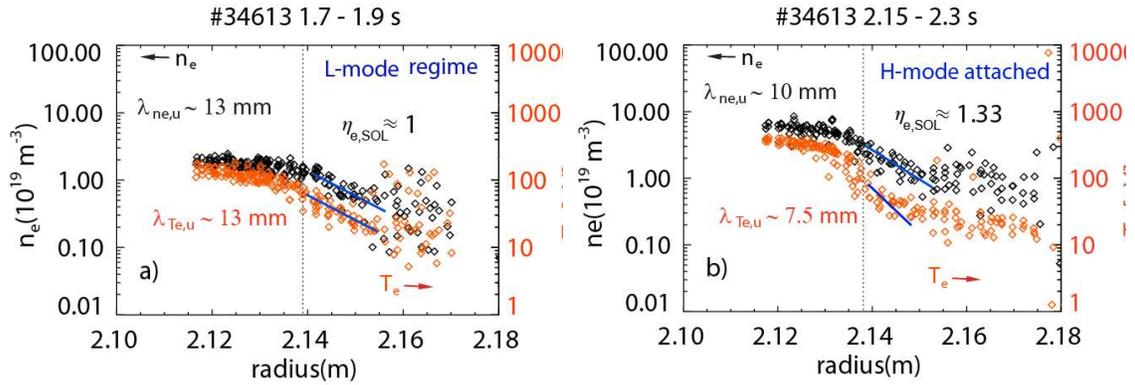


Figure 8. Time traces for AUG discharge #34613: plasma stored energy (W_{MHD}), total heating power and total radiation ($P_{Heat,tot}$ and $P_{rad,tot}$), line integrated edge and core electron density ($\bar{n}_{e,edge}$ and $\bar{n}_{e,core}$), divertor neutral density.



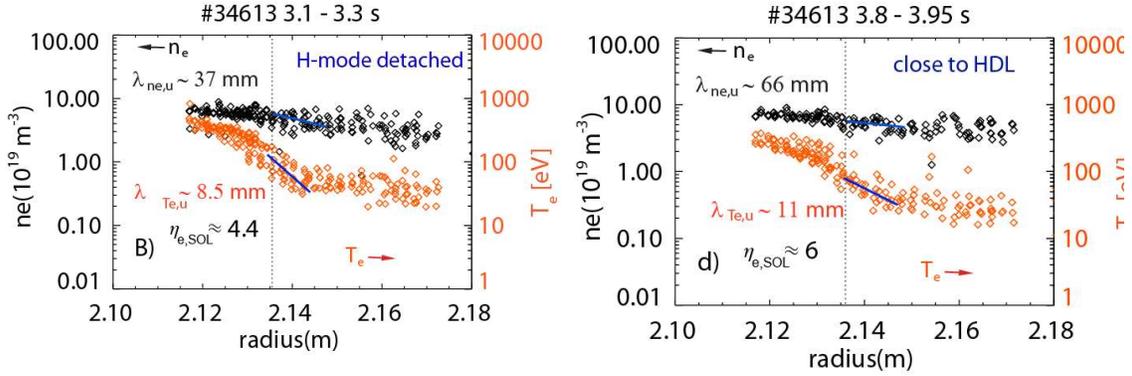


Figure 9. Log-linear plot of the electron temperature and density profiles, for the same AUG discharge as in figure 8, (a) in the L-mode regime; (b) at the beginning of H-mode regime, $n_{e,sep} \approx 3.1 \times 10^{19} m^{-3}$; (c) in the middle of H-mode regime, $n_{e,sep} \approx 4.4 \times 10^{19} m^{-3}$; and (d) at the end of H-mode, just before H-L back transition, $n_{e,sep} \approx 6.0 \times 10^{19} m^{-3}$.

4 Summary

A survey of the correlation between temperature and density gradient length for the near SOL region has been performed for both H- and L-mode regimes in the ASDEX Upgrade tokamak. For both H- and L-mode plasmas, the majority of the gradient length ratio η_e lies between 1 and 2. However, there are data from L-mode discharges which lie below one. These low η_e L-mode discharges are found to be associated with very low density and are more likely in linear/sheath limited divertor regime. In this regime, the difference between the upstream and target temperatures is small and, by assuming constant pressure along a flux tube, the upstream temperature and density can be shown to be closely correlately as, $T_u \propto n_u^{-\frac{2}{3}}$ [10] [11]. Thus, in the linear divertor regime, η_e is expected to be $\eta_e \approx \frac{2}{3}$, which is consistent with the values of these low-density L-mode discharges in this paper. In a previous study, η_e in the pedestal region was observed to lie mainly between 1 and 3 across the entire data set from AUG, JET and DIII-D, as shown in figure 7.43 in Ref [20]. However, there were also several discharges that fall outside of this range, with very low pedestal density, and the gradient length ratio for these is also $\eta_e \approx \frac{2}{3}$. This suggests that, in the linear regime, the correlation between upstream temperature and density profiles also has influence into the pedestal region. Excluding those very low density L-mode discharges, the majority of data should be in the high recycling/conduction limited regime [3] [8]. In the conduction limited regime, there is no obvious correlation between η_e and global plasma parameters, q_{95} or I_p . This could be explained by both λ_{ne} and λ_{Te} having similar dependence on these global parameters. A weak positive trend of higher separatrix temperature $T_{e,sep}$ being associated with a larger η_e in the near SOL region is observed, this can be explained that by assuming that the cross-field transport depends only weakly on $T_{e,sep}$. In this case, increased $T_{e,sep}$ will result in reduced λ_{Te} and so increased η_e due to the positive temperature dependence of the Spitzer-Harm dominated parallel heat transport. In Tore Supra [6], a rough relation of $\lambda_{Te} \approx 1.4 \lambda_{ne}$ is found between the density and temperature decay lengths in the SOL, equivalent to $\eta_e \approx 0.7$. Across this Tore Supra database, the separatrix temperature is about 30 eV. Thus, the lower η_e in the Tore Supra study appears consistent

with the trend observed in AUG. No obvious correlation with separatrix density and collisionality is observed.

For most data, η_e is not far from unity, suggesting that, changing the gradient length of one would affect the other and it is not likely to control these density and temperature profiles separately. However, there are regimes and conditions where the temperature and density profile length scales, and presumably the heat and particle transport, are not coupled. During the transition from the L-mode to the I-mode regime, a relatively large increase of η_e is associated with the temperature profile becoming much steeper while the density profile remains unchanged. During the I-mode phase, a dependence of η_e on temperature alone is also observed: larger η_e being associated with higher local temperature. However, with the same $T_{e,sep}$, $\eta_{e,SOL}$ in I-mode regimes are larger than those in similar H- and L-mode regimes. Under detached divertor condition, the high density transition (HDT) occurs, when the collisionality is sufficiently high, filaments appear as a result of interchange instability which enhance perpendicular particle transport, may trigger a regime transition leading to strongly enhanced perpendicular particle fluxes, resulting in a very flat density profiles [19]. The experimental observation on AUG shows that these convective structures have small influence on heat transport, and so η_e can reach a very high value, up to 6 for the H-mode case studied, whilst the temperature gradient length remains almost constant, only slightly broadens until the separatrix temperature drops. Taken together, the I-mode transition and HDT demonstrate that the electron density and temperature gradients in the SOL can vary independently. Thus, it has been demonstrated that the general correlation between them is due to co-dependence on other parameters and not due to a direct link. However, it remains the case that across the wide database of L-mode and H-mode discharges, sufficiently below the density limit, lie in the range $1 < \eta_e < 2$ in the ASDEX Upgrade tokamak. The divertor, SOL and confinement regimes in this study is summarized in Table 1 along with the SOL temperature and density profile behavior and the associated η_e .

TABLE 1: Divertor, SOL and confinement regimes in this study of AUG plasma and the associated η_e .

Divertor/SOL regime	Confinement mode	SOL temperature and density profiles	η_e
Attached/Sheath limited	L-mode	Temperature and density closely coupled as $T_{u,e} \propto n_{u,e}^{-\frac{2}{3}}$	$\eta_e \approx 2/3$
Attached/Conduction limited	L/H-mode	Both λ_{ne} and λ_{Te} depend on q_{95} or equivalently I_p .	$1 < \eta_e < 2$
	I-mode	Temperature pedestal similar to that of the H-mode, with an edge density profile very similar to that of the L-mode.	$\eta_e > 2$
Detached	L/H-mode	After shoulder formation in the near SOL density profile.	$\eta_e \gg 2$

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