



UKAEA-CCFE-PR(23)05

S. Von Tiedemann, D Collins, M. Gilbert, I. Kodeli

Nuclear Data Uncertainty Propagation and Implications for Radioactive Waste Management of Fusion Steels

Enquiries about copyright and reproduction should in the first instance be addressed to the UKAEA Publications Officer, Culham Science Centre, Building K1/0/83 Abingdon, Oxfordshire, OX14 3DB, UK. The United Kingdom Atomic Energy Authority is the copyright holder.

The contents of this document and all other UKAEA Preprints, Reports and Conference Papers are available to view online free at <u>scientific-publications.ukaea.uk/</u>

Nuclear Data Uncertainty Propagation and Implications for Radioactive Waste Management of Fusion Steels

S. Von Tiedemann, D Collins, M. Gilbert, I. Kodeli

This is a preprint of a paper submitted for publication in Fusion Engineering and Design

1 Highlights

Nuclear Data Uncertainty Propagation and Implications for Radioactive Waste Management of Fusion Steels

- ⁴ Sophia O. von Tiedemann, David M. Collins, Mark R. Gilbert, Ivan A. Kodeli
- Nuclear fusion activation calculations were performed using FISPACT-II
- Sensitivity and uncertainty analysis carried out with XSUN-2022 code package
- Large parts of reactor exceeded UK low-level activity limits for centuries
- Uncertainties in nuclide production larger in rear of blanket than in first wall
- Nuclear data uncertainties affected decay time by several years

Nuclear Data Uncertainty Propagation and Implications for Radioactive Waste Management of Fusion Steels

¹² Sophia O. von Tiedemann^{a,*}, David M. Collins^a, Mark R. Gilbert^b, Ivan A. Kodeli^b

^a University of Birmingham, School of Metallurgy and Materials, Edgbaston, Birmingham, B15 2TT, United Kingdom

^bUnited Kingdom Atomic Energy Authority, Culham Centre for Fusion Energy, Culham Science Centre, Abingdon, OX14 3DB, United Kingdom

13 Abstract

Predictions of material activity in commercial fusion conditions predominantly rely 14 on computational methods, due to a lack of data on long-term effects of high-energy 15 neutron irradiation on structural steels. Consequently, this could result in a bias 16 due to uncertainties in nuclear data used. This work focused on modelling neutron 17 activation of four structural steels in a fusion reactor environment after 20 years 18 of operation. Eurofer, F82H and G91, were assessed as candidate in-vessel ma-19 terials, whereas SS316L(N)-IG was solely modelled in the vacuum vessel. Activa-20 tion calculations were performed using the inventory code FISPACT-II using inputs 21 from Monte-Carlo transport simulations performed with OpenMC. The study em-22 ployed a one-dimensional reactor model with a Helium-Cooled Pebble Bed (HCPB) 23 tritium-breeding blanket design. With the XSUN-2022 code package, a nuclear data 24 sensitivity and uncertainty analysis on production cross-sections of relevant radio-25 nuclides was carried out. Eurofer and F82H steels exhibited significantly higher 26 resistance to neutron activation than G91 and SS316L(N)-IG. At 100 years after 27 shutdown, none of the steels reached UK low-level waste (LLW) activity levels in 28 the first wall. In the rear of the back-support structure (BSS) of the reactor blan-29 ket, all assessed steels reached LLW levels within approximately 30 to 45 years of 30 reactor shutdown. It was found that the vacuum vessel (SS316L(N)-IG) would not 31 be classifiable as LLW for several centuries. Dominant radio-nuclides for each ma-32 terial were identified with FISPACT-II to carry out the uncertainty analyses. The 33 calculated uncertainties were too small to affect the waste disposal options for the 34 first wall within 100 years, but the time-to-reach LLW for BSS and vacuum vessel 35 steel could be uncertain by up to approximately 3 and 6 years, respectively. 36

37 Keywords: Nuclear data, Radioactive waste, Fusion steels, Uncertainty

³⁸ propagation, Sensitivity studies

Preprint submitted to Fusion Engineering and Design

^{*}Corresponding author

Email address: s.o.vontiedemann@bham.ac.uk (Sophia O. von Tiedemann)

³⁹ 1. Introduction

Amid the growing global energy demand and pressure to move away from fossil fuels, 40 nuclear fusion is becoming an increasingly attractive energy source. Where power 41 from nuclear fission produces high-level radioactive waste (HLW), nuclear fusion 42 is anticipated to produce only intermediate- and low-level waste (ILW and LLW), 43 making it more sustainable and favourable over traditional fission power plants. Al-44 though the fusion reaction of tritium and deuterium does not directly create any 45 radioactive products, it results in the emission of high-energy (14 MeV) neutrons. 46 Upon interaction with surrounding materials in the reactor wall, these neutrons can 47 lead to activation and the production of significant volumes of radioactive waste 48 (RW) through transmutation, as well as extensive damage in the material structure 49 through atomic displacement [1]. Structural and other in-vessel materials are antic-50 ipated to be the major source of RW from fusion. 51

52

Due to a lack of existing experimental data on the long-term effects of material expo-53 sure to such high-energy neutron irradiation, predictions of resulting RW activities 54 rely heavily on computational models and approaches. Such methods, however, may 55 be biased due to uncertainties in the nuclear data used. For its fusion programmes, 56 the UK aims to meet LLW criteria (less than 12 MBq/kg of β/γ activity) for most 57 of the RW at 100 years after permanent reactor shutdown. However, it has been 58 predicted that some parts of future fusion reactors are likely to result in significant 59 volumes of waste that would be classified as ILW, even after 100 years after end-of-60 life (EOL), requiring costly geological disposal [2, 3, 4]. 61

62

To mitigate neutron activation of structural steels used in fusion environments as 63 much as possible, reduced activation ferritic/martensitic (RAFM) steels - such as 64 Eurofer [5] and F82H [6] - have been developed for several decades. The application 65 of RAFM steels aims to uphold the necessary physical properties delivered by con-66 ventional structural steels, while reducing neutron activation and hence the amount 67 of RW. To achieve this, the use of alloving elements known to be susceptible to ac-68 tivation (such as Ni, Cu, Nb, Mo) is reduced as much as possible, using less critical 69 elements, such as V, W and Ta, instead [7]. 70

71

Bailey et al. [8] previously found RAFM steels to activate much less than conven-72 tional steels. Of the non-reduced activation FM steels, G91 was least prone to activa-73 tion, which will partly be subject to study here. Another relevant structural steel is 74 the austenitic stainless steel SS316L(N)-IG (hereon referred to as SS316), which will 75 be used extensively for in-vessel structures of ITER [9, 10] – it is also the primary 76 nuclear steel being considered for the vacuum vessel of the future EU DEMOnstra-77 tional power plant (EU-DEMO) [2]. However, due to its high Ni-content, SS316 is 78 highly susceptible to neutron activation, and LLW classification of such a VV under 79

⁸⁰ UK criteria is challenging within 100 years of EOL [11, 12, 13].

81

As existing predictions on waste classifications and activity levels after reactor EOL 82 are mostly based on modelling and simulation, it is crucial to provide correspond-83 ing uncertainty and sensitivity data. Sensitivity data provides insight to which 84 factors of the model have the largest effect on the quantity of interest, whereas 85 uncertainties provide information on the accuracy and hence reliability of a result. 86 This information can subsequently be used to calculate necessary safety margins for 87 quantities such as safe reactor operating times, shielding requirements as well as ac-88 tivity levels of produced RW within the reactor lifetime. Such known uncertainties 89 can then be accounted for in the estimation of operational costs and the necessary 90 handling/disposal of RW. Although previous studies have investigated the effects of 91 uncertainties in nuclear data [16, 17], upon which the majority of activation studies 92 are based, these analyses are generally separated, complicating the direct utilisation 93 of sensitivity and uncertainty results. 94

95

Using computational methods, this study aims to model and compare the neutron 96 activation of four structural steels, which are being considered for application in 97 future fusion power plants. The steels of interest are the FM steel G91 (T2), the 98 austenitic stainless steel SS316L(N)-IG, as well as two RAFM steels: Eurofer and 90 F82H. Subsequently, an independent sensitivity and uncertainty analysis was carried 100 out to study the impact of calculation uncertainty on the activation and subsequent 101 waste classifications. The objective of the work is to exemplify a rigorous method-102 ology by which uncertainties can be included in predictions of the neutron-induced 103 response of fusion materials, which can be applied to subsequent engineering design 104 applications. 105

106 2. Methodology

107 2.1. Transport Simulations

Neutron transport simulations were performed using the Monte Carlo code OpenMC 108 (with ENDF/B-VII.1) [14] to obtain a neutron energy flux spectrum for each steel at 109 a given position. Simulations were based on a simplified, spherical, one-dimensional 110 reactor model, the cross-section of which is shown in Figure 1. An isotropic neu-111 tron source with an average energy of 14.1 MeV was positioned at the centre of 112 the reactor. The blanket configuration and material compositions were based on a 113 helium-cooled pebble bed (HCPB) design [15] for EU-DEMO, summarised in Table 114 1. The standard homogenised DEMO HCPB configuration consists of an armour 115 and first wall (FW), followed by layers for the breeding module (BM), backplates 116 and back support structure (BSS). Note that this model does not reflect the actual 117 engineering design, which would be considerably more complex, but rather approx-118 imates the variation in material through the thickness of the blanket using mixed 119 average compositions (following table 1). This study compares the activation of 120 Eurofer, F82H and G91 in the FW and the outermost layer (5 cm) of the BSS. In 121 all cases, SS316 was assumed for the VV. In the reactor model, the neutron spec-122 tra were tallied in cell tallies, where cells were defined as concentric spheres with 123 a maximum layer width of 5 cm. A total of three simulations were run (one for 124 each option of in-vessel structural material), with 10^{10} neutron histories each. This 125 was judged to be sufficient to ensure good statistical coverage for all the tallies of 126 interest (based on prior experience with similar simulations). 127

Neutron source	Armour First Wall	Breeding Module		Backplates	Back-support Structure		Vacuum Vessel	
	Component	Armour	First Wall	Breeding Module	Backplates	BSS	Vacu Ves	ium sel
	Thickness (cm)	0.2	2.5	52	8.5	60	6	

Figure 1: 1D reactor model employed in neutron transport simulations using OpenMC.

Material (vol.%)	Armour	FW	BM	Backplates	BSS	VV
Tungsten	100	-	-	-	-	-
Eurofer/F82H/G91	-	65	10	41	61	-
Beryllium	-	-	37	-	-	-
Li_4SiO_4	-	-	15	-	-	-
Helium	-	35	38	59	39	-
SS316	-	-	-	-	-	100

Table 1: Employed reactor wall configuration as of the HCPB design.

128 2.2. Inventory Simulations

The above transport simulations were followed by a series of inventory calculations 129 with FISPACT-II (version 5.0) [16] to simulate activity as a function of time and 130 identify the dominant radio-nuclides in each material. The elemental compositions 131 of all evaluated steels are summarised in Table A.9 in the appendix. The FISPACT-132 II simulations used the TENDL-2017 database of nuclear reaction data and the 133 OpenMC-calculated neutron flux spectra, where the spectra were converted from 134 their original neutron cm/source neutron units to neutrons/cm²s units assuming a 135 first wall neutron loading of 2 MWm⁻². Irradiation was simulated for 4 hours per 136 day for 20 years, and subsequent $(\beta + \gamma)$ -activities were evaluated from shut-down 137 up to 200 years after EOL, with the dominant (highest activity) nuclides identified 138 from the evolving inventory as a function of time. 139

140 2.3. Sensitivity and Uncertainty Analysis

The sensitivity and uncertainty analysis (SUA) was carried out separately from any 141 inventory calculations, using the deterministic code package XSUN-2022 [17], includ-142 ing TRANSX-2.15 [18, 19] for the preparation of multi-group nuclear cross-sections, 143 the discrete-ordinate (S_N) transport code PARTISN-5.97 [20, 21], and SUSD3D 144 [22, 23] for the final nuclear data sensitivity and uncertainty calculations. The 145 XSUN-2022 code system involves a complete set of the deterministic codes men-146 tioned above, with the internal data processing shown in Figure 2. Note that in 147 PARTISN, the Vitamin-J 175-energy group structure was used, with the same 1D-148 reactor configuration that was used as in Section 2.1. Sensitivities were calculated 149 using Generalised Perturbation Theory. Reaction rates for the production of nu-150 clides were calculated via: 151

$$RR = \sum_{g} \sigma_{g}^{D} \Phi_{g} \tag{1}$$

Here, σ_g^D is the response function for the nuclide generation reaction in the energy group g, and Φ_g is neutron flux. The corresponding uncertainties were obtained ¹⁵⁴ from the sandwich equation:

$$(\Delta RR)^2 = S^T \cdot Cov \cdot S \tag{2}$$

where S and S^T represent the sensitivity vector of the reaction rate to groupwise

cross sections and its transpose, respectively, and *Cov* is the corresponding cross section covariance matrix.



Figure 2: Work chain of XSUN-2022 code system [17].TRANSX was used to process nuclear data into transport tables compatible with deterministic codes. With PARTISN, the Boltzmann Transport Equation was solved for direct and adjoint flux. ANGELO and NJOY were used for covariance matrix processing. SUSD3D performed the SUA using first order generalized perturbation theory.

157

The activity level A(t) at a given time t, is calculated as the sum of contributions from all decaying nuclides i using their concentrations at the end of operation, viz

$$A(t) = \sum_{i} N_{0i} \lambda_i \mathrm{e}^{-\lambda_i t} \tag{3}$$

¹⁶⁰ Taking into account the uncertainty in nuclide production, this activity becomes:

$$A(t) = \sum_{i} N_{0i} (1 \pm \Delta_i) \lambda_i e^{-\lambda_i (t \pm \Delta t)}$$
(4)

where Δ_i is the uncertainty (from XSUN-2022) in the production of a given radionuclide. Here we define Δt as the delay in reaching a specified activity level (such as the UK LLW limit of 12 MBq/kg) associated with the uncertainty. FISPACT-II activity results from close to the target activity were used to interpolate (in general it is not possible to fit exactly the activity decay curves as there may be contributions from multiple nuclides with different half-lives) the range of 'time-to-target' values (min, max) and hence to obtain Δt . If i = 1, i.e. there is only one dominant nuclide in a material, the time delay Δt due to an uncertainty Δ can be calculated analytically. We have, in this case:

$$\frac{N(t)}{N_0(1\pm\Delta)} = e^{-\lambda(t\pm\Delta t)},\tag{5}$$

where $A(t) = N(t)\lambda$ and

$$\pm \Delta t = \frac{\ln(1 \pm \Delta)}{\lambda} \tag{6}$$

For every assessed steel, an independent sensitivity and uncertainty analysis was 171 conducted for each dominant nuclide contributing above 15% to the total material 172 activity, as identified from the FISPACT-II results (see Tables 2 and 3). For Eurofer, 173 F82H and G91 this was carried out in the front (FW) and the back (BSS) of the 174 reactor wall. For SS316, the sensitivity and uncertainty analysis was performed in 175 the VV. Note that any uncertainties calculated by FISPACT-II only include those 176 associated with the decay constants (λ) and transmutation cross-sections (σ) [16]. 177 The results from XSUN-2022 include uncertainties in the transmutation reaction-178 rates (response functions) as well as uncertainties propagated from transport cross-179 sections (i.e. uncertainty contributions from all nuclide cross-sections which impact 180 neutron transport through the reactor). To ensure compatibility between FISPACT-181 II and XSUN-2022, nuclear data from TENDL-2017 [24] were used for response 182 functions (i.e. SUSD3D used the same data as FISPACT-II for transmutation re-183 actions), whereas JEFF-3.3 [25], which is the reference cross section evaluation in 184 XSUN-2022, was used for transport cross-sections. For cases where the uncertainty 185 is dominated by the response function, values of uncertainties are therefore similar 186 for FISPACT-II and XSUN-2022 results. Note that uncertainties from FISPACT-II 187 are only displayed for total material activity. It is emphasised that the uncertainty 188 analysis conducted using the XSUN-2022 package only encompasses uncertainties 189 on nuclear cross-section data, not on decay data. 190

191

The individual uncertainties in nuclide production cross-sections were used to calculate a lower and upper bound for the amount of each dominant nuclide present in a material, which was then used to define a range of possible material concentrations at the end of operation. These altered compositions were used in FISPACT-II calculations to define the range in activities of each material from which the minimum and maximum 'time-to-reach LLW' was obtained via interpolation (as described above).

¹⁹⁹ 3. Results and Discussion

²⁰⁰ 3.1. Neutron Activation and Waste Categorization Results

Figure 3 shows the total activity of all in-vessel steels 100 years after EOL, as a function of distance through the outboard reactor wall. In each material case, LLW criteria were not met for large proportions of the blanket. For the RAFM steels (Eurofer and F82H), a transition from ILW to LLW activity level was observed approximately half-way through the reactor wall, whereas G91 exceeded those limits almost entirely. For all cases, this mixture of LLW and ILW material within components may complicate the decommissioning process.

208

At the FW (1.45 cm depth), the activity of Eurofer was over an order of magnitude higher than the LLW limit of 12 MBq/kg. F82H exhibited the best resistance to neutron activation throughout, but is still predicted to be activated to over six times the LLW limit at the FW, whereas G91 activity exceeded the LLW limit by more than two orders of magnitude at 100 years.

214

As expected, material activation decreased as a function of distance through the reactor wall due to decreasing neutron fluxes. Activation of Eurofer resulted in ILW at 100 years until the backplate; from approximately 60 cm depth, activity was below the LLW limit at 100 years. For F82H, LLW at 100 years was achieved after about 54 cm (within the BM). G91 performed much worse, with activity only falling below the LLW limit after approximately 115 cm (BSS).

221

In the outermost 5 cm of the BSS, all in-vessel steels categorise as LLW at 100 years. One could argue that it may be more suitable to use RAFM steels closer to plasma-facing components, whereas the advantage of reduced activation is less apparent in the far back. Hence, the use of different (non-RAFM) steels may be more viable in those regions.

227

The total activity of the VV (SS316) (see Table 8) varied with the kind of steel used in the blanket structure. In each case, VV activity exceeded LLW criteria at 100 years after shutdown, reaching 8.05×10^7 Bq/kg, 8.03×10^7 Bq/kg and 9.85×10^7 Bq/kg for the in-vessel steel cases of Eurofer, F82H and G91, respectively. As described in sections 3.1.1 and 3.1.2, the activated nuclides responsible for elevated radioactivity vary between the FW and BSS, even within the same kind of steel.

234 3.1.1. First Wall

Table 2 summarises the dominant nuclides with highest contributions to total steel activity in the FW at 100 years after EOL. Only nuclides contributing more than 15% to total material activity at this time are listed with their relevant production pathways. The production pathway analysis was performed using a tree search algorithm in FISPACT-II, see [16] for more details.

240

The high activity of Eurofer at 100 years after permanent reactor shutdown is due to several nuclides; ^{121m}Sn, ¹²¹Sn, ⁶³Ni and ¹⁴C. In F82H, by contrast, ⁶³Ni was the only dominant nuclide identified, being responsible for about two thirds of the



Figure 3: Total activity of in-vessel materials at 100 years after reactor EOL as a function of depth through the reactor wall, moving radially outwards.

total activity (with minor contributions from a number of other radionuclides - see 244 Figure 5). This is in accordance with their respective compositions, as Eurofer con-245 tains 50 times the amount of tin compared to F82H and six times more nitrogen 246 than F82H, whereas F82H has triple the amount of nickel. ⁶³Ni is also a dominant 247 nuclide in G91, in addition to 91 Nb, which make up approximately 27% and 37%, 248 respectively, of total G91 activity 100 years after EOL. Niobium is an element known 249 to cause activity-related problems and is therefore commonly minimised in RAFM 250 steels; ⁹¹Nb has a half-life of 680 years. With 0.1 wt.%, the Nb-content in G91 is 20 251 times higher than in Eurofer, and 2000 times that of F82H. 252

253

The production pathways in Table 2 show the nuclear reaction cross-sections on which the sensitivity and uncertainty analyses were performed for the FW; the results are presented in Section 3.2.

257

Figure 4 displays the total material activity for Eurofer in the FW from EOL to 259 200 years after shutdown, including relative contribution to activities (top graph)

Steel	Dominant	Contribution	- Ualf lifa	Production	Pathway
Steel	Nuclide	to Activity	nall-life	Pathway	Percentage
Furafor	121m C p	26 50%	11 yoorg	120 Sn $(n,\gamma)^{121m}$ Sn	26.4%
Eurorer	511	20.370	44 years	$^{122}Sn(n,2n)^{121m}Sn$	62.9%
	121 C n	20 607	27 hours	$^{120}\mathrm{Sn}(\mathrm{n},\gamma)^{121}\mathrm{Sn}$	81.4%
	511	20.070	27 nours	$^{122}{ m Sn}({ m n},{ m 2n})^{121}{ m Sn}$	7.7%
	63NI;	10 107	101 woowa	$^{62}\mathrm{Ni}(\mathrm{n},\gamma)^{63}\mathrm{Ni}$	36.8%
		10.1/0	101 years	$ m ^{63}Cu(n,p)^{63}Ni$	56.5%
	$^{14}\mathrm{C}$	15.7%	5705 years	$^{14}N(n,p)$ ^{14}C	99.9%
гооц	63N;	66 907	101 woowa	62 Ni $(n,\gamma)^{63}$ Ni	52.6%
Г 02П		00.070	101 years	$ m ^{63}Cu(n,p)^{63}Ni$	35.0%
C01	91 N L	27 007	680 waawa	${}^{92}Mo(n,np){}^{91}Nb$	83.9%
G91	ND	37.070	000 years	${}^{92}Mo(n,2n){}^{91}Mo(\beta+){}^{91}Nb$	14.6%
	63NI;	26 707	101 woowa	$^{62}\mathrm{Ni}(\mathrm{n},\gamma)^{63}\mathrm{Ni}$	46.5%
		20.170	101 years	${}^{63}{ m Cu(n,p)}{}^{63}{ m Ni}$	42.9%

Table 2: Summary of all radio-nuclides produced through neutron activation in the FW, contributing a minimum of 15% to the overall material activity at 100 years after EOL. The pathway percentage indicates the amount of radio-nuclide produced though the given nuclear reaction.

of individual dominant nuclides. At EOL, dominant nuclides are ⁵⁵Fe as well as
other nuclides, such as ⁵⁴Mn and ¹⁸²Ta, originating from the base and main alloying
elements. Due to their relatively short half-lives, the significance of their activity
decreases with time, whereas the relative contributions of other nuclides increase.
Activity levels 100 years after EOL were approximately an order of magnitude above
the LLW limit, due to the combined activity of ^{121m}Sn, ¹²¹Sn, ⁶³Ni and ¹⁴C, contributing between approximately 16 - 27% each (see Table 2).

267

FW activity of F82H is shown in Figure 5 for the same time period. At EOL, dominant nuclides are ⁵⁵Fe ⁶⁰Co, ³H, which cross-over with ⁶³Ni at about 60 years. The half-life of ⁶³Ni is approximately 100 years, and is largely responsible for exceeding LLW limits at 100 years after EOL. Although the elemental composition of F82H contains less tin and nitrogen, its higher nickel content shifts the relative nuclide contribution to activity from several to just one dominant nuclide.

274

Figure 6 shows activity of G91 in the FW. Similarly to Eurofer and F82H, the main initial activity is due to ⁵⁵Fe, produced from neutron capture of ⁵⁴Fe. At about 50 years after EOL, almost two thirds of its activity are accounted for by ⁶³Ni and ⁹¹Nb. Nuclides that never contribute more than 10% of the total activity during the 200 years of decay are not plotted separately in Figures 4 to 6 but instead their activities are summed together as the "other" curve in the plots. Note that the individual uncertainties in their production cross-sections are assumed to be uncorrelated. It is apparent that, in the FW, neither Eurofer, nor F82H or G91 reach LLW activity levels within the displayed 200 years, where the activity of G91 is over two orders of magnitude above the RAFM steels. In fact, it was found that these structural steels exceeded LLW limits for over 750 (F82H) or even 1000 years (Eurofer, G91). This is problematic since, even with the use of RAFM steels, ILW disposal will most likely not be preventable.

282



Figure 4: (Top): %-Contribution of dominant nuclides to total activity of Eurofer from EOL to 200 years in the FW. (Bottom): Total activity of Eurofer at FW. Nuclides are only displayed if they contributed more than 10% at any point in the decay. The "other" curve is the sum of nuclides which do not meet this criteria; they include 54 Mn and 182 Ta among others.



Figure 5: (Top): %-Contribution of dominant nuclides to total activity of F82H from EOL to 200 years.(Bottom): Total activity of F82H at FW. Nuclides are only displayed if they contributed more than 10% at any point in the decay. The "other" curve is the sum of nuclides which do not meet this criteria; they include ^{121,121m}Sn and ¹⁰⁸Ag among others.



Figure 6: (Top): %-Contribution of dominant nuclides to total activity of G91 in FW from EOL to 200 years. (Bottom): Total activity of G91 at FW. (Bottom): Total activity of F82H at FW. Nuclides are only displayed if they contributed more than 10% at any point in the decay. The "other" curve is the sum of nuclides which do not meet this criteria; they include 94 Nb, 14 C and 60 Co among others.

289 3.1.2. BSS and VV

The identified dominant radio-nuclides in the BSS are summarised in Table 3. Since 290 the BSS reaches LLW limits much earlier than 100 years after EOL, the dominant 291 nuclides are given for 20 years after EOL. Due to the drastic change in neutron 292 energies incident on the material in the BSS compared to the FW, other nuclear 293 cross-sections dominate radio-nuclide production in the BSS. The main nuclides re-294 sponsible for material activity in the BSS are 55 Fe and 60 Co in Eurofer and F82H. 295 where they make up 99.5% of material activity at 20 years after EOL in Eurofer 296 and 99.9% in F82H. G91 does not contain any cobalt, so at 20 years 97.8% of its 297 activity is solely due to ⁵⁵Fe. Figures 7 to 9 show the absolute activities as well as 298 the relative nuclide contributions for Eurofer, F82H and G91 activity in the BSS 299 from EOL to 200 years. These results were obtained from the outermost 5 cm of 300 the BSS, where all in-vessel steels met LLW the criteria within 100 years. In this 301 regime, LLW is reached within approximately 30 to 45 years after shutdown. 302 303

The activity of the VV is shown in Figure 10 from EOL to 200 years after (data 304 shown corresponds to neutron spectra obtained with Eurofer as in-vessel steel). As 305 the VV is made of austenitic stainless steel 316, the high Ni-content of 12.5 wt.%306 leads to significant activation. ⁶³Ni is the dominant nuclide in the VV, with its 307 contribution to total activity varying slightly with the in-vessel steel used. With 308 contributions of 92.5%, 92.7% and 93.7% for Eurofer, F82H and G91, respectively, 309 it is obvious that ⁶³Ni is solely responsible for the failure of the VV to meet LLW 310 requirements 100 years after EOL. Initially, activity is dominated by ⁵⁵Fe and ⁶⁰Co, 311 but as their half-lives are only 2.7 and 5.2 years, respectively, ⁶³Ni starts to dominate 312 material activity after approximately 30 years, accounting for up to 93% of total 313 VV activity at 80 years post EOL. As can be seen in the bottom part of Figure 314 10, the total activity curve follows the 63 Ni-line closely from 60 years onwards. To 315 successfully reduce activation of the VV, ⁶³Ni-production must be prevented either 316 by employment of a different material or by providing sufficient shielding of the VV, 317 which poses a variety of challenges. 318

319

Figure 11 shows the neutron flux profile across the VV of 6 cm thickness. The data is 320 presented corresponding to the 709-group energy-bin structure, as used in FISPACT-321 II, which is a high-resolution grid where the bins are approximately equidistant on 322 a logarithmic scale. The highest energy peak above 10^7 eV represents the direct 14 323 MeV neutrons from the fusion reaction. In the low and thermal energy regions, the 324 neutron flux in the VV is much higher if the in-vessel steel used is G91, compared 325 to the RAFM steels. This difference in flux is responsible for the difference in VV 326 activities between each blanket material case. The flux for Eurofer is also slightly 327 higher than for F82H in this region, which is in agreement with the resulting higher 328 VV activity after 100 years. This suggests that the RAFM steels not only suffer 329

from less neutron activation, but they render a better shielding performance against fusion neutrons than G91 [26].

Table 3: List of all radio-nuclides produced through neutron activation in the BSS, contributing a minimum of 15% to the overall material activity at 20 years after EOL. The pathway percentage indicates the amount of radio-nuclide produced through the given nuclear reaction.

Steel	$\operatorname{Dominant}$	Contribution	Half-life	Production	Pathway
Steel	Nuclide	to Activity	(years)	Pathway	Percentage
Eurofer	55 Fe	71.1%	2.74	54 Fe $(n,\gamma)^{55}$ Fe	89.5%
	60 Co	28.4%	5.27	59 Co(n, γ) 60 Co	99.9%
F82H	55 Fe	83.3%	2.74	54 Fe $(n,\gamma)^{55}$ Fe	89.3%
	60 Co	16.6%	5.27	59 Co(n, γ) 60 Co	99.9%
G91	55 Fe	97.8%	2.74	54 Fe $(n,\gamma)^{55}$ Fe	91.3%



Figure 7: (Top): %-Contribution of dominant nuclides to total activity of Eurofer in BSS from EOL to 200 years. (Bottom): Total activity of Eurofer from EOL to 200 years. Nuclides are only displayed if they contributed more than 10% at any point in the decay. The "other" curve is the sum of nuclides which do not meet this criteria; they include and 14 C, 121,121m Sn and 63 Ni among others.



Figure 8: (Top): %-Contribution of dominant nuclides to total activity of F82H in BSS from EOL to 200 years. (Bottom): Total activity of F82H from EOL to 200 years. Nuclides are only displayed if they contributed more than 10% at any point in the decay. The "other" curve is the sum of nuclides which do not meet this criteria; they include and ¹⁸⁷W, ^{108m}Ag and ⁶³Ni among others.



Figure 9: (Top): %-Contribution of dominant nuclides to total activity of G91 in BSS from EOL to 200 years. (Bottom): Total activity of G91 from EOL to 200 years. Nuclides are only displayed if they contributed more than 10% at any point in the decay. The "other" curve is the sum of nuclides which do not meet this criteria; they include and ¹⁸⁷W, ⁹⁴Nb and ⁵⁹Ni among others.



Figure 10: (Top): %-Contribution of dominant nuclides to total activity of SS316 in VV from EOL to 200 years (using data from simulations with Eurofer as the in-vessel steel). (Bottom): Total activity of SS316 from EOL to 200 years. "Other" displays the sum of any individual radio-nuclides contributing less than 10% to the total material activity at any point within the plotted time period, which include 93 Mo, 59 Ni and 93m Nb among others.



Figure 11: Neutron energy profile across the VV (SS316), shown for each in-vessel material.

332 3.2. Sensitivity and Uncertainty Results

For the dominant nuclides identified for each assessed steel in the FW, BSS and 333 VV, a sensitivity and uncertainty analysis was carried for each production pathway. 334 Table 4 lists the uncertainties for each dominant nuclide production cross-section 335 in the FW and BSS. For the same nuclear cross-section, uncertainties are generally 336 higher in the BSS than in the FW. This is the case for the ${}^{62}Ni(n,\gamma){}^{63}Ni$ reaction 337 in G91, where the uncertainty increases from 0.8% to 3.3%. In regions closer to 338 the source (FW), uncertainties are dominated by uncertainties in the transmutation 339 cross-sections, whereas in the deeper blanket regions (BSS), the contribution of 340 transport cross-sections (such as collisions with nuclides of other blanket materials) 341 increases. However, a direct comparison is not possible with the remaining data 342 shown, as the dominant nuclide and hence their production pathways (cross-sections) 343 change between front and back-end of the blanket, which is due to a change in shape 344 of the neutron flux spectrum. 345

from transport and transmutation uncertainties.
from JEFF-3.3 and are based on the total uncertainty impact on the identified reaction channel
nuclides for each steel assessed in the FW and BSS. For SUSD3D, the combined uncertainties come
Table 4: Uncertainty contributions for the main production pathways of the identified dominant

Rogion	Stool	Dominant	Production	202D2D		
rtegion	Steel	Nuclide	Cross-section	Uncer	tainty	
				$\operatorname{combined}^1$	$\mathrm{response}^3$	
FW	Eurofer	121 Sn	120 Sn(n, γ) 121m Sn	0.9%	$0\%^{2}$	
			$^{122}Sn(n,2n)^{121m}Sn$	11.1%	11.0%	
		⁶³ Ni	62 Ni $(n,\gamma)^{63}$ Ni	1.4%	0%	
			${}^{63}Cu(n,p){}^{63}Ni$	2.1%	2.0%	
		$^{14}\mathrm{C}$	$^{14}N(n,p)^{14}C$	0.4%	$0\%^{2}$	
	F82H	⁶³ Ni	62 Ni $(n,\gamma)^{63}$ Ni	0.9%	0%	
			${}^{63}Cu(n,p){}^{63}Ni$	2.2%	2.0%	
	G91	$^{91}\mathrm{Nb}$	${}^{92}Mo(n,2n){}^{91}Mo(\beta^+){}^{91}Nb$	23.8%	$23.8\%^{3}$	
			$^{92}Mo(n,np)^{91}Nb$	0.5%	$0\%^{2}$	
		⁶³ Ni	62 Ni $(n,\gamma)^{63}$ Ni	0.8%	$0\%^2$	
			63 Cu(n,p) 63 Ni	2.1%	1.9%	
BSS	Eurofer	$^{55}\mathrm{Fe}$	${}^{54}\mathrm{Fe}(\mathrm{n},\gamma){}^{55}\mathrm{Fe}$	28.5%	28.4%	
		$^{60}\mathrm{Co}$	$^{59}\mathrm{Co}(\mathrm{n},\gamma)^{60}\mathrm{Co}$	3.8%	$0\%^{2}$	
	F82H	$^{55}\mathrm{Fe}$	${}^{54}\mathrm{Fe}(\mathrm{n},\gamma){}^{55}\mathrm{Fe}$	28.1%	28.0%	
		$^{60}\mathrm{Co}$	$^{59}\mathrm{Co}(\mathrm{n},\gamma)^{60}\mathrm{Co}$	3.9%	$0\%^2$	
	G91	55 Fe	${}^{54}\mathrm{Fe}(\mathrm{n},\gamma){}^{55}\mathrm{Fe}$	28.3%	28.1%	
		⁶³ Ni	62 Ni $(n,\gamma)^{63}$ Ni	3.3%	$0\%^2$	
		$^{93m}\mathrm{Nb}$	$^{93}Nb(n,n')^{93m}Nb$	10.0%	9.0%	

¹ Uncertainty due to both transport (JEFF-3.3) and transmutation (TENDL-2017) cross-sections

³ Uncertainty due to transmutation (response) TENDL-2017 cross-sections only

 2 Covariance matrices not available in JEFF-3.3 / TENDL-2017

 3 Covariance matrices not available in JEFF-3.3 and taken from ENDF/B-VIII.0

346 TABLE FROM PROOF

Table 5: Uncertainty contributions for the main production pathways of the identified dominant nuclides for each steel assessed in the FW and BSS. For SUSD3D, the combined uncertainties come from JEFF-3.3 and are based on the total uncertainty impact on the identified reaction channel from transport and transmutation uncertainties.

Region	Steel	Dominant nuclide	ominant nuclide Production cross-section		ncertainty
5-6				Combined[a]	Response[b]
\mathbf{FW}	Eurofer	$121\mathrm{Sn}$	$120\mathrm{Sn}(\mathrm{n},\gamma)^{121\mathrm{m}}\mathrm{Sn}$	0.9%	0%[c]
			$122Sn(n,2n)^{121m}Sn$	11.1%	11.0%[3pt]
		63Ni	$62 \mathrm{Ni}(\mathrm{n},\gamma) 63 \mathrm{Ni}$	1.4%	0%
			63 Cu(n,p) 63 Ni	2.1%	2.0% [3pt]
		14C	14N(n,p)14C	0.4%	0%[c]
	F82H	63Ni	$62 \mathrm{Ni}(\mathrm{n},\gamma) 63 \mathrm{Ni}$	0.9%	0%
			$63 \mathrm{Cu(n,p)} 63 \mathrm{Ni}$	2.2%	$2.0\%[3\mathrm{pt}]$
	G91	$91\mathrm{Nb}$	$92Mo(n,2n)91Mo(\beta^+)91Nb$	23.8%	23.8%[d]
			92 Mo(n,np) 91 Nb	0.5%	0%[c] [3pt]
		63Ni	$62 \mathrm{Ni}(\mathrm{n},\gamma) 63 \mathrm{Ni}$	0.8%	0%[c]
			63 Cu(n,p) 63 Ni	2.1%	1.9%
BSS	Eurofer	55Fe	$54 \text{Fe}(n, \gamma) 55 \text{Fe}$	28.5%	28.4%
		$60\mathrm{Co}$	$59 \mathrm{Co}(\mathrm{n}, \gamma) 60 \mathrm{Co}$	3.8%	0%[c] [3pt]
	F82H	$55\mathrm{Fe}$	$54 \mathrm{Fe}(\mathrm{n},\gamma) 55 \mathrm{Fe}$	28.1%	28.0%
		$60\mathrm{Co}$	$59 \mathrm{Co}(\mathrm{n}, \gamma) 60 \mathrm{Co}$	3.9%	0%[c] [3pt]
	G91	$55\mathrm{Fe}$	$54 \mathrm{Fe}(\mathrm{n},\gamma) 55 \mathrm{Fe}$	28.3%	28.1%
		63Ni	$62 \mathrm{Ni}(\mathrm{n}, \gamma) 63 \mathrm{Ni}$	3.3%	0%[c]
		$^{93m}\mathrm{Nb}$	$93Nb(n,n')^{93m}Nb$	10.0%	9.0%

[a]Uncertainty due to both transport (JEFF-3.3) and transmutation (TENDL-2017) cross-sections. [b]Uncertainty due to transmutation (response) TENDL-2017 cross-sections only. [c]Covariance matrices not available in JEFF-3.3/TENDL-2017. [d]Covariance matrices not available in JEFF-3.3 and taken from ENDF/B-VIII.0.

The calculated uncertainties were propagated to evaluate their effect on the average 347 time taken for each assessed steel to reach LLW limits. As described in Section 348 2.3, values for minimum and maximum time to LLW were calculated for the BSS. 349 The FW activities vastly exceeded UK LLW limits far beyond 100 years after EOL, 350 as summarised in Table 6. F82H exhibited the lowest overall activity in the FW, 351 but was still found to require over 750 years to reach the same activity as the LLW 352 limit, far off the desired 100 years. However, Eurofer and G91 were both found to 353 exceed 1000 years to reach the UK LLW limit of 1.2×10^7 Bq/kg, in agreement with 354 previous works [3, 8]. 355

356

Table 6:	Approximate tin	ne taken for	activated in	n-vessel stru	ctural steels	in the FW	to decay to
UK LLW	waste activity li	mits. The a	ctivity level	at 100 years	s after EOL	is also prov	ided.

Stool	Activity 100 years	Time to LLW
Steel	after EOL (Bq/kg)	(years)
Eurofer	1.35E + 08	>1000
F82H	7.33E + 07	>750
G91	1.49E + 09	>1000

The minimum, mean and maximum time taken to meet LLW levels for the BSS 357 were interpolated using FISPACT-II, with the results listed in Table 7. As F82H 358 exhibited the lowest activity in the BSS and throughout the blanket, the corre-359 sponding time to LLW was the shortest. The maximum time to LLW calculated 360 for F82H from the corresponding uncertainties was about 33 years - still more than 361 two years before the earliest possible time predicted for Eurofer to meet LLW levels. 362 However, a mean time of 36 years to LLW for Eurofer is well before the 100-year 363 aim. According to the results of this study, the outermost part of the BSS could 364 reach LLW within 30 years of EOL. This gives a 10-year advantage compared to 365 using G91, which would reach LLW at the earliest within 39.6 years, but may take 366 up to 43 years. Note, that this only assesses the outermost 5 cm of the BSS, and 367 uncertainty values are based on dominant nuclide production only. 368 369

Table 7: Approximate time taken for activated steels in the outermost 5 cm of the BSS to decay to LLW waste activity limits. Minimum and maximum required times were calculated from minimum and maximum amounts of dominant nuclides present due to nuclear data uncertainties.

Steel	Minimum time to LLW (years)	Mean time to LLW (years)	Maximum time to LLW (years)
Eurofer	35.2	36.3	37.3
F82H	29.5	31.5	32.9
G91	39.6	40.2	43.0

VV activity levels at 100 years and uncertainties in ⁶³Ni production in SS316 for each 370 in-vessel case are listed in Table 8, along with the corresponding time uncertainties 371 to reach LLW. FISPACT-II results showed that although in each assessed case the 372 VV contained only SS316, employing RAFM steels for the in-vessel structures had 373 a measurable effect, as previously discussed and shown in Figure 11. The activity 374 of SS316 reached with G91 as the blanket structural material was roughly 12-fold of 375 the VV activity behind the assessed RAFM steels used in the blanket. For Eurofer 376 and F82H used in the blanket structure, the VV would take an excess of 250 years 377 to satisfy LLW criteria, compared to over 500 years for G91. The uncertainty in 378 ⁶³Ni-production in SS316 varied for each shielding material employed, translating 379 to an uncertainty in time taken to reach the LLW limit of 0.3 years for Eurofer, 380 3.9 years for F82H and 5.5 years for G91. Not only would using G91 throughout 381 the reactor blanket result in a much higher activity of SS316, but the corresponding 382 time uncertainty is also higher, making G91 unfavourable in that regard. However, 383 the intended 100 years to LLW will be exceeded significantly by the VV for each 384 shielding material case, hence classification as ILW would come into effect either 385 way, requiring different measures of disposal. Thus, the use of RAFM steels would 386 not yield any valuable advantage regarding the activity of the VV, at least for the 387 present modelling. 388

389

Table 8: Summary of the activity of SS316 in the VV 100 years after EOL for each in-vessel case and corresponding time taken to reach UK LLW levels, the uncertainty of the cross-section producing the dominant nuclide (63 Ni) and corresponding uncertainty in time to reach LLW.

Shielding	VV Activity	Nuclide Production	Uncontainty	Time to	Δt to LLW
Material	$({ m Bq/kg})$	Cross-section	Uncertainty	LLW (years)	(years)
Eurofer	8.05E + 07	62 Ni $(n,\gamma)^{63}$ Ni	0.21%	>250	0.3
F82H	8.03E + 07	$^{62}\mathrm{Ni}(\mathrm{n},\gamma)^{63}\mathrm{Ni}$	2.76%	>250	3.9
G91	$9.85E{+}07$	${}^{62}\mathrm{Ni}(\mathrm{n},\gamma){}^{63}\mathrm{Ni}$	3.85%	>500	5.5

Since this study employed a very simplified, one-dimensional reactor model utilising 390 homogenised material cells in the neutron transport calculations, the authors em-391 phasise that the calculated activities, uncertainties and time-scales are subject to 392 those simplifications. Uncertainties were calculated only for the production of radio-393 nuclides deemed "dominant"; other sources of uncertainty in the model or approach 394 were not investigated in this study and should be addressed separately. Whereas 395 this model is representative of a cross-section through the blanket and vacuum ves-396 sel, other parts of the reactor, such as the divertor, were not considered. Hence a 397 more realistic reactor model may lead to different results including material activity 398 and time-to-LLW. Since the feasibility of this method of combining activation cal-399 culations with nuclear data uncertainty propagation has now been demonstrated. 400 a similar investigation may be carried out in the future with more mature reactor 401

402 models and nuclear data.

403 4. Conclusion

Neutron transport simulations and material inventory calculations were performed to obtain the activity of the RAFM steels Eurofer, F82H and the FM steel G91 for in-vessel use as well as for SS316 in the VV. Subsequently, a series of uncertainty analyses were conducted to investigate the effect of nuclear data uncertainties on the overall prediction of time required for activated materials to reach UK LLW limits after reactor end-of-life. Overall, it can be concluded from this study that:

- RAFM steels were activated significantly less than G91, but none of the assessed in-vessel steels met LLW activity levels in the FW and for large parts of the blanket far beyond 100 years. At the rear of the blanket (BSS), Eurofer, F82H and G91 all achieved LLW limits within 45 years of EOL. Even with the use of Eurofer or F82H, it will not be possible to decommission the entire blanket as LLW within 100 years of EOL.
- In the VV, the activity of SS316 was more than one order of magnitude above the LLW limit due to its high Ni-content. Therefore, the VV will not meet LLW requirements for centuries.
- The make-up of dominant nuclides in the FW varied with each material composition, whereas activity in the BSS was mainly caused by ⁵⁵Fe and ⁶⁰Co, due to neutron flux softening.
- The mean time-to-LLW for the FW exceeded 750 years for F82H and 1000 years for Eurofer and G91. In the outermost 5 cm of the blanket (BSS), this was approximately 36, 32 and 40 years for Eurofer, F82H and G91, respectively.
- Uncertainties in the production cross-sections of dominant nuclides were propagated to estimate error margins for the calculated time needed to reach the LLW limit. This resulted in a time uncertainty of no more than three years in the BSS and up to six years in the VV.
- Nuclear data uncertainties are only a small part of a large set of uncertainty contributors affecting quantities such as reactor lifetime, activity, dose levels as well as operational and decommissioning costs. However, the uncertainties found in the selected nuclear data alone were sufficient to potentially shift the point at which steels meet the UK LLW activity limit by several years.
- The feasibility of a rigorous methodology to perform independent uncertainty analyses on nuclear data has been demonstrated. The authors recommended

that similar analyses be carried out regularly in the future, when covariance 437 matrix data become more mature in the available nuclear data libraries and 438 parameters (material composition, irradiation scenario, etc.) are more refined. 439 It is highlighted that this method is translatable to the modelling of other 440 quantities or radiological predictions directly and indirectly related to radio-441 nuclide production. As existing literature on nuclear data uncertainties and 442 their effects is scarce, further studies could generate valuable understanding 443 relevant to the planning and implementation of nuclear fusion as an established 444 energy source. 445

446 5. Acknowledgements

This work was partially carried out within the framework of the EUROfusion Consortium, funded by the European Union via the Euratom Research and Training Programme (Grant Agreement No 101052200 — EUROfusion) and from the EPSRC Energy Programme (grant number EP/W006839/1). Views and opinions expressed are however those of the authors only and do not necessarily reflect those of the European Union or the European Commission. Neither the European Union nor the European Commission can be held responsible for them.

454 Appendix A. Sample Appendix Section

⁴⁵⁵ Elemental compositions of the evaluated steels are summarised below in Table A.9.

Flomont	Composition (wt.%)						
Element	Eurofer	F82H	G91 (T2)	SS316L(N)-IG			
Fe	Base	Base	Base	Base			
Al	0.01	0.01	0.02	-			
Ag	-	0.002	-	-			
As	0.05	0.002	0.01	-			
В	0.002	0.0003	0.001	0.001			
С	0.11	0.1	0.12	0.03			
Co	0.01	0.005	-	0.05			
Cr	9	8	9.5	18			
Cu	0.01	0.01	0.1	0.3			
Mn	0.4	0.1	0.5	2			
Mo	0.005	0.001	1.05	2.7			
Ν	0.03	0.005	0.07	0.08			
Nb	0.005	0.00005	0.1	0.01			
Ni	0.01	0.03	0.2	12.5			
Ο	0.01	0.005	-	-			
Р	0.005	0.005	0.02	0.025			
\mathbf{S}	0.005	0.002	0.005	0.01			
Sb	0.05	0.0005	0.03	-			
Sn	0.05	0.001	0.01	-			
Si	0.05	0.1	0.4	0.5			
Ta	0.12	0.04	-	0.01			
Ti	0.02	0.05	0.01	0.1			
V	0.2	0.2	0.25	-			
W	1.1	2	0.05	-			
Zr	0.05	-	0.01	-			

Table A.9: Elemental compositions for Eurofer [27], F82H [28], G91 (T2) [8] and SS316 [2]

456 References

[1] M. Hernández-Mayoral, M. Caturla, 8 - microstructure evolution of irradiated

- 458 structural materials in nuclear power plants, Woodhead Publishing (2010) 180, 225 doi:https://doi.org/10.1522/0781845600056.2.180
- 459 189–235doi:https://doi.org/10.1533/9781845699956.2.189.
- 460 URL https://www.sciencedirect.com/science/article/pii/B9781845695118500081

[2] M. Gilbert, T. Eade, C. Bachmann, U. Fischer, N. Taylor, Activation, decay heat, and waste classification studies of the European DEMO concept, Nuclear

Fusion 57 (4) (2017) 046015. doi:10.1088/1741-4326/aa5bd7. 463 URL https://iopscience.iop.org/article/10.1088/1741-4326/aa5bd7 464 [3] M. Gilbert, T. Eade, T. Rey, R. Vale, C. Bachmann, U. Fischer, N. Taylor, 465 Waste implications from minor impurities in European DEMO materials, Nu-466 clear Fusion 59 (7) (2019) 076015. doi:10.1088/1741-4326/ab154e. 467 URL https://iopscience.iop.org/article/10.1088/1741-4326/ab154e 468 [4] R. Pampin, S. Zheng, S. Lilley, B. Na, M. Loughlin, N. Taylor, Activation 469 analyses updating the ITER radioactive waste assessment, Fusion Engineer-470 ing and Design 87 (7-8) (2012) 1230–1234. doi:10.1016/j.fusengdes.2012.02.110. 471 $\mathrm{URL}\,\mathtt{https://linkinghub.elsevier.com/retrieve/pii/S0920379612001846}$ 472 [5] B. van der Schaaf, F. Tavassoli, C. Fazio, E. Rigal, E. Diegele, R. Lindau, 473 G. LeMarois, The development of eurofer reduced activation steel, Fusion 474 Engineering and Design 69 (1) (2003) 197–203, 22nd Symposium on Fusion 475 Technology. doi:https://doi.org/10.1016/S0920-3796(03)00337-5. 476 URL https://www.sciencedirect.com/science/article/pii/S0920379603003375 477 [6] H. Tanigawa, E. Gaganidze, T. Hirose, M. Ando, S. Zinkle, R. Lindau, 478 E. Diegele, Development of benchmark reduced activation ferritic/martensitic 479 steels for fusion energy applications, Nuclear Fusion 57 (9) (2017) 092004. 480 doi:10.1088/1741-4326/57/9/092004. 481 URL https://iopscience.iop.org/article/10.1088/1741-4326/57/9/092004 482 A.-A. Tavassoli, A. Alamo, L. Bedel, L. Forest, J.-M. Gentzbittel, J.-W. Rens-|7|483 man, E. Diegele, R. Lindau, M. Schirra, R. Schmitt, H. Schneider, C. Petersen, 484 A.-M. Lancha, P. Fernandez, G. Filacchioni, M. Maday, K. Mergia, N. Boukos, 485 Baluc, P. Spätig, E. Alves, E. Lucon, Materials design data for reduced 486 activation martensitic steel type EUROFER, Journal of Nuclear Materials 487 329-333 (2004) 257-262. doi:10.1016/j.jnucmat.2004.04.020. 488 URL https://linkinghub.elsevier.com/retrieve/pii/S0022311504001485 489 [8] G. Bailey, O. Vilkhivskaya, M. Gilbert, Waste expectations of fusion steels 490 under current waste repository criteria, Nuclear Fusion 61 (3) (2021) 036010. 491 doi:10.1088/1741-4326/abc933. 492 URL https://doi.org/10.1088/1741-4326/abc933 493 9 M. Fabbri, D. Leichtle, A. Martin, R. Pampin, E. Polunovskiv, Nuclear heat 494 analysis for the iter vacuum vessel regular sector, Fusion Engineering and De-495 sign 137 (2018) 435–439. doi:https://doi.org/10.1016/j.fusengdes.2018.04.022. 496 $\mathrm{URL}\,\mathtt{https://www.sciencedirect.com/science/article/pii/S0920379618303107}$ 497

- [10] G. Stankunas, S. Breidokaite, A. Tidikas, Activation analysis and 498 evaluation of radionuclide inventory decay heat for eu demo vac-499 uum vessel components, Fusion Science and Technology 77(7-8)500 (2021)791-801. arXiv:https://doi.org/10.1080/15361055.2021.1906153, 501 doi:10.1080/15361055.2021.1906153. 502
- ⁵⁰³ URL https://doi.org/10.1080/15361055.2021.1906153
- [11] H. L. SWAMI, C. DANANI, A. K. SHAW, Activation characteristics of candidate structural materials for a near-term indian fusion reactor and the impact of their impurities on design considerations, Plasma Science and Technology 20 (6) (2018) 065602. doi:10.1088/2058-6272/aaabb4.
- ⁵⁰⁸ URL https://doi.org/10.1088/2058-6272/aaabb4
- [12] P. Pereslavtsev, F. Cismondi, F. A. Hernández, Analyses of the shielding options for HCPB DEMO blanket, Fusion Engineering and Design 156 (2020)
 111605. doi:10.1016/j.fusengdes.2020.111605.
- ⁵¹² URL https://linkinghub.elsevier.com/retrieve/pii/S0920379620301538
- [13] M. R. Gilbert, T. Eade, C. Bachmann, U. Fischer, N. P. Tav-513 lor, Waste assessment of european DEMO fusion reactor designs, Fus. 514 Eng. Des. 136 (2018) 42-48, special Issue: Proceedings of the 13th 515 International Symposium on Fusion Nuclear Technology (ISFNT-13), 516 https://doi.org/10.1016/j.fusengdes.2017.12.019. 517
- [14] P. Κ. Romano, В. Forget, The openmc monte carlo particle 518 Annals of Nuclear Energy 51(2013)274 - 281.transport code, 519 doi:https://doi.org/10.1016/j.anucene.2012.06.040. 520
- ⁵²¹ URL https://www.sciencedirect.com/science/article/pii/S0306454912003283
- ⁵²² [15] D. Carloni, L. Boccaccini, F. Franza, S. Kecskes, Requirements for helium cooled pebble bed blanket and r&d activities, Fusion Engineering and Design
 ⁵²⁴ 89 (7) (2014) 1341–1345, proceedings of the 11th International Symposium on Fusion Nuclear Technology-11 (ISFNT-11) Barcelona, Spain, 15-20 September,
- ⁵²⁶ 2013. doi:https://doi.org/10.1016/j.fusengdes.2014.02.036.
- ⁵²⁷ URL https://www.sciencedirect.com/science/article/pii/S0920379614001203
- [16] J.-C. Sublet, J. Eastwood, J. Morgan, M. Gilbert, M. Fleming, W. Arter,
 Fispact-ii: An advanced simulation system for activation, transmutation and
 material modelling, Nuclear Data Sheets 139 (2017) 77–137, special Issue on
 Nuclear Reaction Data. doi:https://doi.org/10.1016/j.nds.2017.01.002.
- ⁵³² URL https://www.sciencedirect.com/science/article/pii/S0090375217300029
- [17] I. Kodeli, S. Slavič, SUSD3D computer code as part of the XSUN-2017 windows
 interface environment for deterministic radiation transport and cross-section

- sensitivity-uncertainty analysis, Science and Technology of Nuclear Installa tions 2017 (2017) 1–16. doi:10.1155/2017/1264736.
- [18] R. E. MacFarlane, TRANSX 2: A code for interfacing MATXS cross-section
 libraries to nuclear transport codes (7 1992).
- ⁵³⁹ URL https://www.osti.gov/biblio/5133321
- [19] PSR-0317 TRANSX-2.15 nuclear energy agency,
 https://www.oecd-nea.org/tools/abstract/detail/psr-0317/04, ver sion: 05/01/2018.
- [20] R. E. Alcouffe, R. S. Baker, J. A. Dahl, S. A. Turner, R. C. Ward, Partisn 5.97,
 1-d, 2-d, 3-d time-dependent, multi- group deterministic parallel neutral particle transport code, la-ur-08-07258, NEA/Data Bank CCC-0760/01 Computer
 Code Collection (11 2008).
- ⁵⁴⁷ URL https://www.oecd-nea.org/tools/abstract/detail/ccc-0760/
- 548 [21] CCC-0760 PARTISN 5.97 nuclear energy agency,
 549 https://www.oecd-nea.org/tools/abstract/detail/ccc-0760/01, ver 550 sion: 02/11/2009.
- [22] I. Kodeli, Multidimensional deterministic nuclear data sensitivity and uncer tainty code system, method and application, Nuclear Science and Engineering
 138 (2001) 45–66.
- ⁵⁵⁴ [23] NEA-1628 SUSD3D nuclear energy agency,
 ⁵⁵⁵ https://www.oecd-nea.org/tools/abstract/detail/nea-1628/03, ver ⁵⁵⁶ sion: 2008.
- ⁵⁵⁷ [24] A. J. Koning, D. Rochman, TENDL-2017, release Date: April 25, 2018. Available from https://tendl.web.psi.ch/tendl_2017/tendl2017.html (2017).
- ⁵⁵⁹ [25] The JEFF team, JEFF-3.3: Evaluated nuclear data library, release Date: ⁵⁶⁰ November, 20 (2017).
- ⁵⁶¹ URL https://www.oecd-nea.org/dbdata/jeff/jeff33/index.html
- ⁵⁶² [26] C. Liu, H. Yang, J. Zhang, J. Zhang, L. Li, Y. Qiu, D. Yao, X. Gao, Current
 ⁵⁶³ status and progress on the shielding blanket of cfetr, IEEE Transactions on
 ⁵⁶⁴ Plasma Science 46 (5) (2018) 1417–1421. doi:10.1109/TPS.2018.2792453.
- ⁵⁶⁵ [27] F. Tavassoli, Eurofer Steel, Development to Full Code Qualification, Procedia ⁵⁶⁶ Engineering (2013) 9.
- L. Tan, Y. Katoh, A.-A. Tavassoli, J. Henry, M. Rieth, H. Sakasegawa,
 H. Tanigawa, Q. Huang, Recent status and improvement of reduced-activation

- ⁵⁶⁹ ferritic-martensitic steels for high-temperature service, Journal of Nuclear Ma-
- terials 479 (2016) 515–523. doi:https://doi.org/10.1016/j.jnucmat.2016.07.054.
- URL https://www.sciencedirect.com/science/article/pii/S0022311516304755