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Abstract

This work applies the coupled JINTRAC and QuaLiKiz-neural-network (QLKNN) model on the ohmic current ramp-up phase of a JET D discharge. The chosen scenario exhibits a hollow T_e profile attributed to core impurity accumulation, which is observed to worsen with the increasing fuel ion mass from D to T. A dynamic D simulation was validated, evolving j, n_e, T_e, T_i, n_{Be}, n_{Ni}, and n_W for 7.25 s along with self-consistent equilibrium calculations, and was consequently extended to simulate a pure T plasma in a predict-first exercise. The light impurity (Be) accounted for Z_{eff} while the heavy impurities (Ni, W) accounted for P_{rad}. This study reveals the role of transport on the T_e hollowing, which originates from the isotope effect on the electron-ion energy exchange affecting T_i. This exercise successfully affirmed isotopic trends from previous H experiments and provided engineering targets used to recreate the D q-profile in T experiments, demonstrating the potential of neural network surrogates for fast routine analysis and discharge design. However, discrepancies were found between the impurity transport behaviour of QuaLiKiz and QLKNN, which lead to notable T_e hollowing differences. Further investigation into the turbulent component of heavy impurity transport is recommended.

1 Introduction

The prediction of temporal plasma evolution is essential for modern fusion devices, in order to estimate and optimize the performance of any given plasma scenario. The demand for this capability is expected to increase as the devices scale in cost, energy output, and complexity, as is the case with ITER. At the start of every tokamak plasma discharge, the plasma is initiated in a breakdown and burn-through phases, then follows a set of control trajectories to set up the desired steady-state magnetic configuration after the plasma current exists called the current ramp-up phase. Accurate modelling of this phase is particularly important in preparation of the ITER pre-full-pulse-operation (PFPO) campaigns [1]. The lessons learned from the WEST device [2] highlight the necessity and the challenges of successfully developing

the current ramp-up phase, from both an operational and a plasma physics perspective. An examination of the validity of integrated plasma models in this phase helps to identify any gaps in the existing code base and enables these tools to assist in the development and interpretation of these campaigns [3].

Integrated plasma transport models interconnect multiple independent physics models to consistently evaluate the plasma state under different concurrent phenomena. Each of these component models are responsible for computing a specific plasma phenomenon, e.g. transport fluxes, sources or sinks. This study focuses on the JINTRAC [4] model, which contains a time evolution calculation scheme which allows it to take advantage of the timescale separation between various physics phenomena to optimize and orchestrate the execution of the modules without compromising the physical plasma description.

High-fidelity physics models are often precluded from high throughput analysis workflows due to their heavy computational costs, especially when using a time evolution scheme such as JINTRAC. Within the current tokamak plasma modelling landscape, the simulation bottleneck is typically the turbulent transport calculations. Neural network (NN) surrogate models of the reduced-order gyrokinetic turbulent transport model, QuaLiKiz, aptly named QuaLiKiz-neural-network (QLKNN) [5], provided a promising method of improving this speed without sacrificing as much accuracy as other existing simplified models. The QLKNN family of models combine careful data curation and threshold-aware cost functions to more accurately capture the crucial salient features of microturbulent transport physics. This study focuses primarily on the latest version trained based on JET experimental profile data, labelled QLKNN-jetexp-15D [6], mostly to verify its applicability to a variety of JET discharges.

The coupled JINTRAC plasma transport and QLKNN turbulent transport surrogate model was applied to model a reference deuterium (D) hybrid ramp-up discharge, JET#97776. The hybrid scenario was developed to access a larger β_N than the standard H-mode scenarios by preventing the development of large (1,1) MHD instabilities, or sawtooth crashes, which occur inside the $q = 1$ surface [7].

While the plasma conditions recover between sawteeth crashes, they were also observed to trigger neoclassical tearing modes (NTM) outside the $q = 1$ surface which limit the achievable experimental β_N [8]. Since the sawteeth instabilities are strongly tied to the $q = 1$ flux surface, they can be avoided by applying strong safety factor, q, profile tailoring during the ramp-up phase [9–11] to keep $q > 1$ across the entire plasma radius. Once a desirable q profile is established, it is "frozen" in place by activating the external heating systems, where the resulting high temperatures lower the plasma resistivity and thus limiting the current diffusion [12].

Incidentally, since this JET high-β_N plasma regime is typically dominated by ion temperature gradient (ITG) driven turbulence, the q-profile tailoring also improves the turbulent transport characteristics during the high-power phase. This is both due to the inner region with low magnetic shear [13] and the suppression of microturbulent instabilities around $\rho_{\mathrm{tor}} \simeq 0.7$ with higher magnetic shear, $\hat{s}$ [14]. The turbulence in both of these regions are both conducive to electromagnetic stabilization processes [15, 16].

However, there are risks that accompany the performance benefits of this hybrid scenario. The hybrid ramp-up at JET exhibits a local decrease of the electron temperature, T_e, within the central core, referred to as *temperature hollowing*. If this effect is strong enough, the core resistivity changes sufficiently to reverse the current density gradient, ∇j, leading to a reversal of the magnetic shear, $\hat{s}$. The extrapolation of previous protium (H) experiments indicate that the temperature hollowing effect will be greater in the tritium (T) experiments [17]. This shear reversal can leading to additional MHD instabilities [18], which are observed to be reliable precursors for plasma disruptions [19, 20] if they stop rotating in the plasma or *lock*. For this reason, any detection of their presence generally triggers an early plasma termination within the tokamak control system, making the approach design more difficult.

This study also serves a secondary purpose other than facilitating a physics analysis of the hybrid ramp-up scenario. As the previous QLKNN-jetexp-15D model validation exercises [6] were performed us-

ing plasma profiles which were also used in creating the training dataset, this exercise also tests the performance of the NN on data not explicitly included within the training dataset, i.e. whether it can provide accurate estimations for scenarios that were not explicitly used to guide its development.

As the impurity transport is expected to be crucial in modelling this scenario, it is noted that the QLKNN-jetexp-15D model does not directly predict the impurity particle fluxes, $\mathbf{\Gamma}_{\mathrm{imp}}$. Instead, it assumes that the electron particle flux predictions, $\mathbf{\Gamma}_e$, scales according to the following equation:

$$\mathbf{\Gamma}_i = \frac{n_i}{n_e} \mathbf{\Gamma}_e \qquad (1)$$

where i denotes any generic ion species, including the main fuel ions. As a quasilinear formulation of turbulent contributions to impurity fluxes is still under investigation [21], this ad-hoc implementation allows for the imposition of ambipolarity as an interim constraint. Due to the existence of multiple impurity species in the simulation, this constraint does not yield a unique impurity transport configuration on its own. While the validity and utility of this ad-hoc implementation is discussed in this report, an investigation of the exact turbulent contribution to the impurity particle flux is outside the scope of this study.

Section 2 provides the details of the D hybrid discharge, JET#97776, relevant to this study and Section 3 describes the development and validation results of the reference simulation of the D discharge. Section 4 describes the extrapolation of the reference D simulation to a H and T plasma, including a discussion comparing the isotopic trends observed in the simulations against known experimental ones. Section 5 then compares the predictive extrapolation simulations against actual T experiments performed with the assistance of said simulations and a deeper examination of the ad-hoc ambipolarity constraint implementation mentioned above. Finally, a summary is provided in Section 6 and comments are made on any potential future work.

2   Reference JET D hybrid discharge

The chosen D reference discharge, JET#97776, was an ohmic test of the proposed ramp-up approach to the JET hybrid scenario in preparation of the JET tritium campaigns in 2021. The reference discharge was performed with $I_p = 2.26$ MA and $B_T = 3.43$ T inside the current flat-top phase. As this exercise focuses on the current ramp-up phase, the actual values of these parameters change over the duration of the analysis window. Figure 1 shows the time trace of the global plasma and pertinent engineering parameters for this discharge.

In addition to clearly exhibiting the hollow temperature profile under investigation, this particular discharge was chosen since the auxiliary heating was intentionally omitted. For this particular hybrid ramp-up approach, both the NBI and ICRH systems are typically activated at $t = 7$ s. The lack of auxiliary heating allows for a more accurate depiction of the plasma state at the planned activation time of the auxiliary heating. Additionally, it allowed the current redistribution to continue until the formation of an unstable MHD $(1,1)$ mode [22], or sawtooth instability, at $t \simeq 8$ s. The onset time and radial extent of the first sawtooth crash is largely attributed to evolution of the safety factor, q, profile, specifically the $q = 1$ surface [23], which is itself determined by the current density profile, j. These experimental conditions both simplify the modelling requirements and allow an ideal environment for a simultaneous validation of the time-dependent current diffusion model.

3   Reference D simulation

Table 1 shows the most important details of the JINTRAC simulation using the QLKNN-jetexp-15D model to predict the turbulent transport coefficients. The simulated plasma start time ($t = 1.77$ s) was chosen to be near the first reconstructed equilibrium within the discharge with a diverted geometry, i.e. after X-point formation, and the end time ($t = 9$ s) was chosen to be sufficiently past the onset time of the MHD $(1,1)$ mode to ensure that it was resolved in the simulation. The input data preparation for the initial and boundary conditions of the model was done us-

Figure 1: Time traces for deuterium reference hybrid ramp-up discharge, JET#97776, over the selected analysis window. The nominal time at which the auxiliary heating systems are activated to transition into the high-power phase (vertical dashed line) and the time of the first observed sawtooth crash at $t = 7.995$ s (vertical dashed-dotted line) are indicated. The auxiliary heating systems were not activated during this discharge as it was an ohmic test pulse. The fast current ramp allows faster current penetration into the core to maintain a higher q_0 and a generally flatter q profile, while the subsequent current drop ensures a higher q_{95} right before the heating systems are activated, maintaining a higher magnetic shear, $\hat{s}$, to reduce the NTM drive at higher rational q surfaces. This approach leads to the plasma current "overshoot" feature around $t = 5$ s [9].

ing the EX2GK tool using Gaussian Process regression fits [24], with the new addition of electron cyclotron emission (ECE) diagnostics for improved electron temperature resolution in the inner core. Due to importance of the current diffusion in this scenario, the electron temperature [25], T_e, boundary condition was deemed a particularly crucial quantity to estimate accurately.

The radiative loss terms in the energy balance equations were expected to play a major role in successfully capturing the temperature hollowing effect. It was suspected that this sink term is dominated by the heavier impurity species inside the plasma core. Thus, a reasonable estimate of the initial and boundary conditions (IC and BC) for the impurity profiles were required within this simulation in addition to the T_e boundary condition. Due to the lack of a transport barrier at the edge within this L-mode plasma and the non-stationary nature of the ramp-up phase, the impurity concentration within the core at a given plasma time would be best computed using the self-consistent evolution equations in the model. It should be noted that a quantitative comparison of the impurity density evolution capability of this configuration of JINTRAC is not well validated.

In order to simultaneously follow Z_{eff}, P_{rad}, and keep the light impurity dilution within acceptable limits, i.e. $n_{\text{Be}}/n_e \lesssim 0.03$, it was necessary to include three impurity species within the simulation: one light, one mid-range and one heavy impurity species. Due to the specific materials of the JET plasma-facing components, these were selected as beryllium (Be), nickel (Ni) and tungsten (W), respectively. The initial impurity concentrations and profiles were estimated by adjusting them until the fixed-boundary equilibrium calculation module, ESCO [4], converged while using its recommended settings. The time-dependent impurity BC was then adjusted until the time evolution of the effective charge, Z_{eff}, and the total radiated power, P_{rad}, agreed reasonably with the Bremsstrahlung optical spectroscopy and bolometer measurements, respectively. Figure 2 shows the comparison of the simulated time traces to the experimental 0D values for P_{rad} and Z_{eff}. The effective charge measurement after $t = 7.5$ s implies a pure hydrogenic plasma, within the uncertainty of the mea-

Table 1: Summary table of most pertinent JINTRAC settings of the D reference ramp-up simulation. All listed components are internally evaluated in time within the JINTRAC model, yielding a self-consistent solution.

	JET#97776
Description	Ohmic ramp-up
Predicted quantities	j, T_e, T_i, n_e, n_{Be}, n_{Ni}, n_{W}
# of radial grid points	101
Plasma time[1]	$1.77 - 9.0$ s
Max. plasma time step	0.001 s
Sim. boundary (ρ_{tor})	0.9
QuaLiKiz region (ρ_{tor})	$0.03 - 0.9$
Equilibrium model	ESCO – fixed boundary
Impurity trans. model	SANCO
Radiation model	ADAS cooling factors
ADAS year	Be: 89, Ni: 96, W: 50
Neoclassical trans. model	NCLASS
Turbulent trans. model	QLKNN-jetexp-15D
NN part. trans.[2]	1 – see Equation (1)
QuaLiKiz E×B option	0 – no E×B suppression

[1] The reference time, $t = 0$, in the JET data system is when the magnetic coils start ramping up, instead of the usual time of plasma breakdown. These two events are typically 40 s apart at JET.

[2] This option determines the relation used to estimate the impurity particle transport coefficients from quantities predicted by the QLKNN model, with option 1 corresponding to Equation (1).

surement. Although Z_{eff} is expected to be low in this phase of the discharge, it was deemed more reasonable to assume a target of $Z_{\mathrm{eff}} \simeq 1.15$ when deciding the impurity densities at those times[1].

Figure 3 show the time trace comparison of the volume-averaged electron temperature, $\langle T_e \rangle$, the volume-averaged electron density, $\langle n_e \rangle$, the ITER normalized internal inductance [26], $l_i(3)$, and the safety factor at the magnetic axis, q_0, from the reference simulation and their corresponding processed measurement values inside the JET data repository.

Figure 2: Comparison of simulated time traces (colored lines) against measured data (gray line) for JET#97776 from $t = 1.77 - 9.0$ s, using two different initial q profiles from the pressure-constrained equilibrium reconstruction (EFTP – blue) and the polarimetry-constrained equilibrium reconstruction (EFTF – red). The quantities shown are the line-integrated effective charge (top left), Z_{eff}, the total radiated power (top right), P_{rad}.

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[1]The Z_{eff} signal has since been reprocessed to correct the value after $t = 47$ s. However, this also increases the rest of the Z_{eff} time trace, resulting in $Z_{\mathrm{eff}} \simeq 1.75$ at $t = 42$ s. As this significantly impacts both the impurity density ICs and BCs, the absolute value of the predicted profiles may be affected but the trends and physical explanations provided in this study are not.

Good general agreement is achieved with the experimental measurements, with the minor exception of $\langle T_e \rangle$ being systematically higher. While it is not inherently concerning, the reason for this will be discussed along with the comparison of the 1D T_e profiles.

The reference simulation was executed using two initial q profiles, each coming from a different EFIT [27] equilibrium reconstruction:

- one constrained using experimental pressure profiles (EFTP);

- and one constrained using both the experimental pressure profiles and polarimetry measurements of the core magnetic field, (EFTF).

This was done to assess the relative impact of the q profile ICs on the time evolution of the simulation. The main difference between the two q profile ICs are primarily seen in q_0 but they converge as the time evolution progresses, as expected after many current diffusion times have elapsed. The effect on the global plasma parameters is negligible. Thus, no further adjustments were made to the impurity density ICs and BCs between the two q profile cases shown in this section.

Since the safety factor on the magnetic axis, q_0, drops below unity in the simulation, an additional validation metric can be performed to ensure that the reference simulation adequately models the dynamic current redistribution. Figure 4 compares the time-dependence of the q profile at the observed inversion radius and the observed sawtooth onset time, given by the vertical dashed line. As shown in the figure, the arrival of the $q = 1$ surface at $\rho_{\text{tor}} = 0.2$ within the simulation is $\sim$500 ms earlier than the first experimentally observed sawtooth crash. While further modelling with MHD codes is required to determine the actual sufficiency of the plasma condition to the formation of sawteeth instabilities, this heuristic comparison is taken to indicate that the neoclassical conductivity used within JINTRAC, combined with the improved kinetic profiles, provides an adequate description of the current diffusion within these plasma conditions.

Although the results are not shown, additional simulations attempted using only the Spitzer conduc-

Figure 3: Comparison of simulated time traces (colored lines) against measured data (gray line) for JET#97776 from $t = 1.77 - 9.0$ s, using two different initial q profiles from the pressure-constrained equilibrium reconstruction (EFTP – blue) and the polarimetry-constrained equilibrium reconstruction (EFTF – red). The quantities shown are the volume-averaged electron density (top left), $\bar{n}_e$, the volume-averaged electron temperature (top right), $\bar{T}_e$, the normalized internal inductance (bottom left), l_{i3}, and the safety factor at the magnetic axis (bottom right), q_0.

Figure 4: Time trace of the safety factor, q, value at the inversion radius of the first experimentally observed sawtooth crash, $\rho_{\text{tor}} = 0.2$. The arrival of the $q = 1$ (horizontal dashed line) surface at this location occurs at $t = 7.5$ s (dotted vertical line), which is ~ 500 ms before the first experimentally observed sawtooth crash at $t = 7.955$ s (dashed vertical line).

tivity term [28] within JINTRAC resulted in the q_0 dropping to ~ 0.5 at $t = 3$ s. This implies that the neoclassical formulation is required for accurately modelling these plasma regimes. This appears in contrast with recent current diffusion modelling results of the ramp-up phase in MAST and JET [29]. Future work is recommended in determining whether the discrepancy between this work and the previous modelling exercise came from a difference in plasma regimes or from improved data handling to better capture the T_e BC.

Figures 5 and 6 shows the kinetic profiles for the electron density, n_e, electron temperature, T_e, ion temperature, T_i, and the safety factor, q, profiles at various time slices within the simulation. These figures indicate that the density peaking and temperature hollowing effects observed in the measurements are adequately captured, along with their time evolution.

The initial q profile switch altered the depth of the T_e hollowing as the plasma evolved, but it is uncertain which q profile IC is more accurate due to the discrepancy at $\rho_{\text{tor}} = 0.3$. Since the initial q profile has only a minor global impact on the simulation and the

ESCO equilibrium boundary is taken from the EFTF calculation, the EFTF q profile run was used as the reference simulation for this study. Additionally, it is noted that the simulated q profile does not agree well with the estimate from the EFIT algorithm. This further underlines the necessity to perform a predictive current simulation to self-consistently relax the current profile before using it within an integrated modelling exercise, as is current recommended practice within the field. Also, the discrepancy between the simulated and measured $\bar{T}_e$ shown in Figure 3 can be attributed to the increased T_e around $\rho_{\text{tor}} = 0.3$, although it is unclear if this is a result of the neoclassical transport or the NN turbulent transport predictions at this moment.

3.1 Role of impurities on observed temperature hollowing

Figure 7 shows the time traces for the predicted volume-averaged impurity density concentrations for the three impurity species (Be, Ni, and W) and Figure 8 shows the density profiles for each of the 3 impurity species at the same time slices as above. It can be directly seen from these profiles that the temperature hollowing comes from the accumulation of high-Z impurity species in the central core, where the peaking is much stronger for heavier impurities than light ones. This is the expected behaviour from the neoclassical inward flux for heavy impurities in this plasma scenario [30], although it is worth noting that the turbulent transport contribution to the impurity fluxes do not seem to heavily modify this response.

Since the central temperature hollowing effect is tied to the impurity transport, two possible explanations were provided regarding the mechanism:

1. through the increased local effective charge, Z_{eff}, which lowers the neoclassical conductivity [25], preventing current penetration and reducing the ohmic heating term in the energy conservation equation;

2. and through the increased local radiated power density, Q_{rad}, which acts as a sink term in the energy conservation equation.

Figure 5: Comparison of the simulated electron temperature, T_e, profile against the time-averaged measurement data of JET#97776 for three different time slices: $t = 3$ s (left), 5 s (center), and 7 s (right). The dashed vertical lines indicates the internal simulation boundary, $\rho_{tor} = 0.9$, outside which the profiles are prescribed in the simulation.

Figure 6: Comparison of the simulated electron density, n_e, profile against the time-averaged measurement data of JET#97776 for three different time slices: $t = 3$ s (left), 5 s (center), and 7 s (right). The dashed vertical lines indicates the internal simulation boundary, $\rho_{tor} = 0.9$, outside which the profiles are prescribed in the simulation.

Note that since these two channels stem from different physics phenomena affecting the same plasma quantity, it is suspected that both are present but it is crucial to understand their relative contribution. Due to the success of the reference simulation on capturing the T_e hollowing, additional simulations were performed to attempt to determine the relative importance of these two mechanisms. The main conclusion of this separation study is that Z_{eff} only affects the central T_e as an integrated effect from the plasma edge, whereas Q_{rad} has a more local impact. This confirms that an increased T_e hollowing within the scenario reflects an increase of heavy impurity accumulation specifically within $\rho_{tor} < 0.2$, as opposed to an increase only inside the turbulence region of $\rho_{tor} \geq 0.4$. It also implies that the observed temperature hollowing is more dependent on Q_{rad} than the neoclassical conductivity. For brevity, the spe-

cific details of this separation study are both shown and discussed further in Appendix B.

4 Predict-first T simulations

Since the current ramp-up approach for this scenario has already been developed and tested in the device using the D reference discharge, this predictive study assumes it can serve as an initial guess for the engineering waveforms required in the T discharge. The integrated modelling framework can then be used to investigate the performance of the scenario in tritium and determine which factors are the most effective for configuring the associated paired experiment. Specifically, the T_e and q profiles can then be monitored to ensure that the T ramp-up approach follows the reference D discharge as closely as possible to avoid triggering the disruption mitigation system.

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Figure 7: Time traces for the volume-averaged impurity concentration for the reference D simulation of JET#97776, containing Be (left), Ni (center), and W (right).

With the reference simulation configured and its results verified, as discussed in Section 3, incremental modifications can be made to it in order to predict the outcome of future experimental scenarios. Over the course of this study, these modification include:

- switching the main ion isotope;
- adjusting the electron density, n_e, boundary condition;
- and adjusting the impurity density, n_{Be}, n_{Ni}, n_W, boundary condition.

For brevity, only the modelling results of the first two items are described in the main body of this paper due to their relevance to the main conclusions. The last item is detailed further in Appendix C.

4.1 Extrapolation to T plasma

This section examines the effects of switching the fuel mix in the experiment from 100% D to 100% T without adjusting any other simulation input. For a more complete picture of the isotope effect, this exercise was also performed using a 100% protium (H) fuel mix.

Figure 9 shows the comparison of the time traces of the global plasma parameters using the different main ion isotopes in the simulation. In general, there is no significant difference in the global plasma parameters with respect to the isotopic mass. Some minor effects are the increase of total P_{rad} with isotope mass after $t = 7$ s without a corresponding increase in global Z_{eff}, indicating a higher heavy impurity content with isotope mass, as discussed in Section 3.1.

Figure 10 shows the kinetic profiles only at $t = 7$ s, as it represents the plasma state when auxiliary heating is switched on to effectively freeze the q profile. Figure 11 shows the impurity density profiles at this same snapshot, $t = 7$ s. From these figures, the main differences between the different simulations indicate that an increase in the main fuel ion mass leads to:

- an increase in T_e hollowing;
- a decrease in T_i across the entire radial profile;
- an increase in the n_e peaking;
- a delay in the arrival of $q = 1$ surface at $\rho_{tor} = 0.2$;
- an increase in the core accumulation of heavy impurities, Ni and W;
- and an increase of n_{Ni} and n_W around $\rho_{tor} \simeq 0.5$.

While the T_e and q impacts are expected trends from previous experiments comparing hybrid ramp-up phases in H and D plasmas [17], the T_i and n_e are not. The T_i trends may have been present in the previous isotope experiments but the measurements necessary to determine this were not available due to the lack of NBI injection in these discharges.

The lower T_i within the T plasma simulation is likely an effect of the reduced electron-ion energy exchange due to the increased main ion mass. It is postulated that this lower T_i is the source of the increased temperature hollowing, as it reduces the neoclassical temperature screening effect. Under these conditions, this translates into an increase in the inward flux of impurities leading to increased core im-

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Figure 8: Snapshots of the impurity density profiles, n_{Be} (top), n_{Ni} (middle), and n_W (bottom), for the reference D simulation of JET#97776 for three different time slices: $t = 2$ s (left), 5 s (center), and 7 s (right). The dashed vertical lines indicates the internal simulation boundary, $\rho_{tor} = 0.9$, outside which the profiles are prescribed in the simulation.

Figure 10: Comparison of the electron temperature (leftmost), T_e, ion temperature (left middle), T_i, electron density (right middle), n_e, and safety factor (rightmost), q, profiles at $t = 7$ s for the isotopic extrapolation of the D reference simulation based on JET#97776, executed using a 100% fuel ion mix of D (blue), H (red) and T (purple). The dashed vertical lines indicates the internal simulation boundary, $\rho_{\rm tor} = 0.9$, outside which the profiles are prescribed in the simulation.

Figure 9: Comparison of the simulated time traces for the isotopic extrapolation of the D reference simulation based on JET#97776 from $t = 1.77 - 9.0$ s, executed using a 100% fuel ion mix of D (blue), H (red) and T (purple). The quantities shown are the line-integrated effective charge (top left), $Z_{\rm eff}$, the total radiated power (top right), $P_{\rm rad}$, the volume-averaged electron density (middle left), $\bar{n}_e$, the volume-averaged electron temperature (middle right), $\bar{T}_e$, the normalized internal inductance (bottom left), l_{i3}, and the safety factor at the inversion radius (bottom right), q at $\rho_{\rm tor} = 0.2$.

purity accumulation. This then increases the radiated power density, $Q_{\rm rad}$, in the central core which is responsible for the T_e hollowing in that region, as detailed within Appendix B. This lower core T_e drives the collisionality up, leading to a local increase in current diffusion and generating the reverse shear in the core region if the T_e hollowing is severe enough.

To further emphasize the dominant impact of the local radiated power density, $Q_{\rm rad}$, on the temperature hollowing over the neoclassical conductivity, Figure 12 compares the local ohmic heat source, $Q_{\rm ohm}$, the local radiation power density, $Q_{\rm rad}$, and the local electron-ion energy exchange source, $Q_{\rm ex}$ across the isotope scan. As mentioned in Section 3.1, one expected result of the increased core collisionality is the reduction of current diffusion via a lower local neoclassical conductivity, resulting in a lower central $Q_{\rm ohm}$. However, Figure 12 shows that the central $Q_{\rm ohm}$ actually increases with isotopic mass. Since the ohmic heating process itself occurs via collisional processes, it is suspected that the lower current diffusion is more than compensated by the increased

Figure 11: Comparison of the impurity density, n_{imp}, profiles of Be (left), Ni (center), and W (right) at $t = 7$ s for the isotopic extrapolation of the D reference simulation based on JET#97776, executed using a 100% fuel ion mix of D (blue), H (red) and T (purple). The dashed vertical lines indicates the internal simulation boundary, $\rho_{\text{tor}} = 0.9$, outside which the profiles are prescribed in the simulation.

collisional energy transfer. In this case, the temperature hollowing truly only occurs due to the significant increase in Q_{rad} incurred by the increased influx of heavy impurities.

With regards to the other simulation results, the increased delay of the $q = 1$ surface arrival with the isotope mass is an experimentally observed phenomenon [17], but the time difference is quantitatively larger within the experiments. Also, the increasing n_e peaking prediction within the simulation goes contrary to the experimental trends, but the quantitative magnitude of this trend is sufficiently small to be considered negligible. Due to the contribution of n_{Be} to n_e, it is also possible that the reversal of the trend is due to quantitative discrepancies in the impurity transport coefficients between the model and experiment. A few additional insights into these discrepancies are provided in Section 5.

4.2 Advising the development of the paired T experiment

Since both current diffusion and ohmic heating are modified by the plasma collisionality, it is suspected that the q and T_e profiles can be affected by changing the plasma density, indicated by n_e. This fact has been experimentally exploited in previous H discharges to construct a paired experiment with the same q profile as the D reference discharge. Comparing these experiments allow the separation of phenomena which are affected directly by the main ion

mass from those which are affected indirectly via another plasma parameter.

Within the previous JET H discharges, the core plasma density was controlled by injecting neutral gas into the SOL. While other methods are available for controlling the core impurity accumulation [2, 31, 32], this was ultimately chosen due to experimental constraints but remains suitable for this study due to its simplicity in modelling. By empirically extrapolating the gas flow rate necessary to produce the paired experiments in H [17], it was estimated that an increase of 20–30% was needed in the gas injection system in order to achieve the same q profile in T. This range accounts for the empirical uncertainties regarding the effect of impurities and other isotope effects. This modelling study was performed to justify and/or refine that empirical extrapolation, as well as provide key physics insights into the impurity and other isotope effects which affect that estimate.

As mentioned earlier, the integrated model settings chosen for this exercise do not include any regions beyond the internal boundary condition, including the separatrix region and scrape-off layer (SOL). Thus, the effect of the additional gas injection cannot be directly modelled but a proxy can be made by adjusting the time-dependent n_e BC ($\rho_{\text{tor}} = 0.9$) in the simulation inputs. Figure 13 shows the results of the n_e BC scan on the kinetic profiles at $t = 7$ s, including the self-consistent evolution of the density profile. Figure 14 shows the results of the n_e BC scan on the

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Figure 12: Comparison of the local ohmic heat source (left), Q_{ohm}, local radiated heat source (center), Q_{rad}, and local electron-ion energy exchange source (right), Q_{ex}, at $t = 7$ s for the isotopic extrapolation of the D reference simulation based on JET#97776, executed using a 100% fuel ion mix of D (blue), H (red) and T (purple). The plotted region ends at the internal simulation boundary, $\rho_{\mathrm{tor}} = 0.9$, due to the large non-physical values present beyond this region.

Figure 13: Comparison of the electron temperature (leftmost), T_e, ion temperature (left middle), T_i, electron density (right middle), n_e, and safety factor (rightmost), q, profiles at $t = 7$ s for the electron density BC scan of the T extrapolation simulation (blue) based on JET#97776. The dashed vertical lines indicates the internal simulation boundary, $\rho_{\mathrm{tor}} = 0.9$, outside which the profiles are prescribed in the simulation.

impurity density profiles at $t = 7$ s.

As expected from the H experiments, the higher n_e from the BC modification results in less T_e hollowing. This is due to the increase in the plasma collisionality, which adjusts the turbulent transport such that it lowers the overall density gradient in the profile. This in turn reduces the neoclassical inward flux of heavy impurities, leading to less impurity accumu-

lation and a lower Q_{rad} in the central core region. Due to the explicit connection of this mechanism to the ITG turbulence regime present in this scenario, this method of controlling T_e hollowing via gas puff is likely not universally applicable.

5 Simulation validation studies

5.1 Comparison with post-analysis T experiment

A pair of discharges were performed incorporating the results of the predictive modelling exercise discussed above. This pair of discharges are hybrid ramp-up discharges with a similar approach to the reference discharge (JET#97776) but using tritium as the main fuel. One (JET#98562) was performed at the reference gas flow rate, which ultimately led to a triggering of the disruption mitigation system, and the other (JET#98567) was performed with an increased gas flow rate targeting a ∼20% higher line-integrated density, which did not trigger the disruption mitigation system. It is noted that the exact percentage increase varied with time in the T experiment although this set a valuable target.

While the success of the predictive modelling exercise in quantitatively advising target control room parameters is already a milestone achievement for the modelling suite, it is also interesting to inspect the quality of the predicted 1D profiles. Figure 16 shows the comparison of the predicted kinetic profile from

Figure 14: Comparison of the impurity density, n_{imp}, profiles of Be (left), Ni (center), and W (right) at $t = 7$ s for the electron density BC scan of the T extrapolation simulation (blue) based on JET#97776. The dashed vertical lines indicates the internal simulation boundary, $\rho_{\text{tor}} = 0.9$, outside which the profiles are prescribed in the simulation.

Figure 15: Comparison of the local ohmic heat source (left), Q_{ohm}, local radiated heat source (center), Q_{rad}, and local electron-ion energy exchange source (right), Q_{ex}, at $t = 7$ s for the electron density BC scan of the T extrapolation simulation (blue) based on JET#97776. The plotted region ends at the internal simulation boundary, $\rho_{\text{tor}} = 0.9$, due to the large non-physical values present beyond this region.

extrapolating the D reference to T against the experimental measurements taken from this pair of tritium discharges.

From these figures, the n_e and T_e predictions between $\rho_{\text{tor}} \simeq 0.3$ and the simulation boundary shows excellent agreement with the experimental measurements within their uncertainties. However, discrepancies beyond those uncertainties are found in the central region, from $0 \leq \rho_{\text{tor}} \lesssim 0.3$. This implies that the impurity profile predictions, and consequently the impurity transport coefficients from QLKNN, may not be fully representative of the physical scenario. A deeper investigation into the turbulent component of the impurity transport behaviour is recommended, but remains outside the scope of this study.

5.2 Comparison with QuaLiKiz model

In order to complete the QLKNN validation exercise, it is prudent to compare the performance of the reference discharge using QLKNN against the original QuaLiKiz model. However, this comparison is complicated by to the impurity transport assumption made in implementing the QLKNN within JINTRAC.

Firstly, an initial comparison was made including a modification to the QuaLiKiz implementation within JINTRAC such that it also used the ad-hoc impurity transport scaling described by Equation (1). Figures 17 and 18 show the comparison between these two simulations, where all the QLKNN kinetic profiles fall within the $\pm10\%$ profile RMS error with respect to QuaLiKiz, as quoted by Ref. [6].

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Figure 16: Comparison of the predictive T extrapolation simulations from the D reference discharge (solid lines) and recent JET T experiments (gray points). The scenario using the same absolute gas flow as the D reference (JET#98562 – top) resulted in a mitigated disruption, whereas the scenario with an increased gas flow (JET#98567 – bottom) did not. The profiles agree for $\rho_{tor} \gtrsim 0.3$, but significant discrepancies occur within the central core.

These plots confirm that the QuaLiKiz electron and main ion transport coefficients in this physical scenario are accurately replicated by the QLKNN-jetexp-15D model, effectively providing additional validation for those output variables. With this confirmed, the next step is removing the ad-hoc impurity transport modification to determine whether it is the source of the discrepancy and/or the extent of its impact on the simulation as a whole. Figure 19 shows the results of the D reference simulation with the ad-hoc ambipolarity constraint removed.

This shows that the original QuaLiKiz impurity profile predictions vary significantly from those of the ad-hoc rule. The increased impurity density in the central core, specifically the radiating species of Ni and W, explains the increased T_e hollowing. Additionally, the variations in the impurity density and its gradient may alter the turbulent transport characteristics sufficiently to change the n_e profile, especially near the simulation boundary.

A more direct and quantitative comparison of the Ni and W particle flux over the initial time steps of the simulation is provided in Figures 20 and 21 respectively. This indicates that the ad-hoc ambipolarity rule assigns an outward flux across the predictive region in the first 100 ms, which helps to establish a degree of hollowness in the impurity density profile. Then, the outward flux slowly transitions to an inward flux as the discharge progresses, resulting in a slow buildup of impurities in the core with a lower density around $\rho_{tor} = 0.2$. The neoclassical component (not shown in the plots) is then responsible for the central accumulation observed in the impurity profiles.

In contrast, the QuaLiKiz model provides a much higher inward flux from the beginning of the discharge and remains heavily inward as the discharge progresses. This leads to a notably higher density at $\rho_{tor} = 0.2$, which is then transported into the central core via neoclassical transport. This leads to a much higher impurity accumulation and a more pronounced T_e hollowing effect. While the turbulent component of impurity transport has not been fully validated within the QuaLiKiz code, the importance of its accurate prediction in the ramp-up phase is highlighted within this study. While a deeper funda-

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Figure 17: Profile comparisons of the electron temperature (left), T_e, ion temperature (center), T_i, and electron density (right), n_e, using QLKNN (blue line) and a modified QuaLiKiz implementation (red line) as the turbulent transport model. The QuaLiKiz implementation in JINTRAC was modified by replacing the computed impurity particle flux, Γ_{imp}, with a scaled electron particle flux, Γ_e, according to Equation (1).

Figure 18: Profile comparisons of the Be impurity density (left), n_{Be}, Ni impurity density (center), n_{Ni}, and W impurity density (right), n_{W}, using QLKNN (blue line) and a modified QuaLiKiz implementation (red line) as the turbulent transport model. The QuaLiKiz implementation in JINTRAC was modified by replacing the computed impurity particle flux, Γ_{imp}, with a scaled electron particle flux, Γ_e, according to Equation (1).

mental investigation into this phenomenon is outside the scope of this work, this paper strongly recommends it before these reduced models are used for extensive prediction in non-H-mode plasma regimes.

6   Conclusion

Overall, the JINTRAC integrated model, coupled with the newly-developed QLKNN-jetexp-15D model for turbulent transport predictions, was successfully validated on the deuterium hybrid ramp-up discharge, JET#97776, and further used for predictive simulations for tritium plasma operation at JET. This predictive capability of the model is verified not only via the simulation of the plasma ramp-up

regime, a novel application for the QuaLiKiz model in general, but also the via the extrapolation of the scenario to a pure T discharge, returning results which are in good agreement with isotope-dependent trends from previous H discharges. With T, the T_e hollowing increases and the onset of the (1,1) MHD sawtooth instability may be delayed as indicated through the $q = 1$ arrival time. The only trend which appears in opposition to the experimentally observed trends is the density peaking, which increases with isotope mass in the simulations but decreases in experiments. However, the relative change of the density peaking from D to T is small enough in both simulation and experiment that both fall within the GPR fit uncertainties and may be considered negligible.

Figure 19: Profile comparisons of the electron temperature (left), T_e, electron density (center), n_e, and the Be impurity density (left), n_{Ni}, using QLKNN (blue line) and the original QuaLiKiz model (red line) as the turbulent transport model.

Figure 20: Comparison of the turbulent impurity (Ni) particle flux predictions as a function of radius, ρ_{tor}, and time, t, within the JINTRAC model using QLKNN-jetexp-15D (top) and the original QuaLiKiz model (bottom) as the turbulent transport model. Negative values (blue regions) indicate an inward flux, transport towards the central core, and positive values (red regions) indicate an outward flux, transport towards the separatrix.

Figure 21: Comparison of the turbulent impurity (W) particle flux predictions as a function of radius, ρ_{tor}, and time, t, within the JINTRAC model using QLKNN-jetexp-15D (top) and the original QuaLiKiz model (bottom) as the turbulent transport model. Negative values (blue regions) indicate an inward flux, transport towards the central core, and positive values (red regions) indicate an outward flux, transport towards the separatrix.

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Aside from the demonstration of the predictive capabilities of the integrated modelling suite, this study also provides some insight into the physical process causing the increased temperature hollowing observed when moving from H to D in the previous isotopic experiments. The simulations in this study indicate that the increased main ion mass lowers T_i through a lower electron-ion heat exchange. This increases the neoclassical inward flux of the heavy impurities (e.g. Ni and W in this study) and leads to more accumulation in the central core. These heavy impurities then act as a local heat sink via radiation losses, effectively lowering the core T_e. The consequence is then a higher collisionality, increased current diffusion, and a higher q_0. While each of these individual mechanisms and correlations are well-documented within the literature of their respective specialized physics domains, the chain of interactions only became apparent through the analysis of this scenario within an integrated model. A key advantage of this approach is the capability to determine the relative strength of a given interaction compared to all the others present in a given scenario, as evidenced by this study.

The validation of the QLKNN-jetexp-15D simulations against the original QuaLiKiz model was complicated by the use of the ad-hoc ambipolarity constraint for the impurity flux, as QLKNN-jetexp-15D was not trained directly to predict turbulent impurity transport coefficients. In terms of the electron and main ion coefficients, the two models show good agreement as shown using an implementation of QuaLiKiz which uses the same ad-hoc ambipolarity constraint. However, without this constraint, the two models exhibit significant discrepancies which are attributed to the differences in the turbulent predictions of the impurity transport coefficients. A deeper fundamental study into the validity of these impurity transport coefficients within the original QuaLiKiz model are recommended, since the predicted profiles with the original model also deviated from the experimental measurements.

Finally, the increased computational speed of the QLKNN-jetexp-15D model was leveraged to assist with answering operational questions within the time constraints of the planned experimental schedule.

Each full ramp-up simulation took ∼6 h on a single CPU, allowing not only for the rapid development of a state-of-the-art validated ramp-up simulation reference but also for the execution of numerous sensitivity studies required to provide meaningful physics insights. Approximately 500 simulations were completed over the span of 4 calendar weeks for the bulk of this study. In the end, the predictive study estimated a required gas injection rate increase of 20 – 30%, corresponding to a line-averaged density increase of ∼20%, with reasonable assumptions on the impact of the gas injection on the plasma edge and SOL. This agrees the empirical trends from isotope effect studies from previous H discharges, which was extrapolated to T for an estimated gas injection increase of 30%. The effectiveness of this qualitative recipe was further corroborated by a pair of JET T discharges and provides even greater confidence in the predictive capabilities of the JINTRAC integrated model, using the QLKNN-jetexp-15D turbulent transport predictions, within this plasma regime at JET.

7 Acknowledgements

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A   Reference simulation transport coefficients

Figures 22 and 23 show the heat fluxes, q, and particle fluxes, Γ, respectively, for various species in the plasma at specific time slices of the reference simulation. It is important to restate that the turbulent contribution to the ion particle fluxes is proportional to the electron particle flux, following the ad-hoc relation described by Equation (1). Thus, all the profiles of the turbulent contributions in the reference simulation, shown in Figure 23, are the same and the only variation between the species is the neoclassical contribution. Then, it is interesting to note that this neoclassical contribution remains a major component, even for the main fuel ion.

Concerning the discrepancies observed in Section 5.2, it appears that the turbulent component of the particle flux in the original QuaLiKiz model can dominate over the neoclassical component. This ends up changing the direction of the heavy impurity particle transport, leading to different T_e hollowing dynamics in the ramp-up simulations. One possible explanation is the inadequacy of the electrostatic assumption of QuaLiKiz to model turbulent impurity transport in high-β regimes [33].

B   Separation of effective charge and radiation effects

In order to separate the effects of Z_{eff} and Q_{rad} on the temperature profiles, the time evolution of the n_e and all n_{imp} were imported from the reference simulation and fixed within a new one. This allowed an independent variation of the impurity density profiles for this investigation, where:

- Z_{eff} was isolated by modifying the Be impurity profile, since its Q_{rad} contribution is negligible at these core plasma temperatures;

- and Q_{rad} was isolated by modifying the W impurity profile, since its Z_{eff} contribution is minor due to its low absolute density in the plasma.

By investigating these effects through the adjustment of the impurity densities, the self-consistent calcula-

Figure 22: Example impurity density profile modifications applied to the prescribed integrated model inputs (top row) for separating local effective charge, Z_{eff}, effects (via Be – left column) and local radiated power density, Q_{rad}, effects (via W – right column). The impact of these modifications on the simulated time traces of total radiated power, P_{rad}, (middle row) and the line-integrated effective charge, Z_{eff}, (bottom row) are also provided for demonstrate its validity.

tion capabilities of the integrated modelling framework can be used to develop a more physically tangible intuition within the study. In this case, this is done via the SANCO [4] impurity transport module. Although the impurity density evolution equations are not solved in these simulations, the module is still used to compute the radiated power based on the charge state distribution and the ADAS cross-section library.

Figure 24 shows the modified impurity density profiles used in these simulations along with their effects on the Z_{eff} and P_{rad} time traces. The modifications were added in two ways:

- using a constant multiplication factor across the entire profile and a modification applied only in the inner core;

- and using a Gaussian multiplication factor centered on $\rho_{\mathrm{tor}} = 0$ with $\sigma = 0.067$.

These two different modification methods help to distinguish which effects are localized to the inner core region and which are effectively integrated across the entire plasma radius.

Figures 25 and 26 show the sensitivities due to changes in the local Z_{eff}, via Be profile modifications, and the local Q_{rad}, via W profile modifications, respectively. Only the profiles at $t = 7$ s are provided since the effect over the time evolution is fairly consistent.

From these figures, the temperature hollowing is affected more by the local Q_{rad} channel than the local Z_{eff} channel in the central core region. Also, the Q_{rad} contribution can be considered a local effect as the uniform multiplication factor has the same impact on the inner core temperature as the Gaussian multiplication factor. On the other hand, the Z_{eff} contribution is an accumulated global effect as the uniform multiplication factor shows a continuous divergence of the profile from the edge while the Gaussian multiplication factor has nearly no impact.

This Z_{eff} modification result also raises the possibility that the increased local effective charge, Z_{eff}, also lowers the turbulent transport coefficients within the central core. This can be done either through a stabilization of TEM modes by increasing the collisionality, or through a stabilization of ITG modes due to the impact of light impurity density and/or density peaking [34]. This is further supported by the fact that the q profiles for all the various Z_{eff} variations are nearly identical, although the plots for this are not provided in this report for brevity.

C Impurity boundary condition scans

Due to the higher mass of tritium over deuterium, an increase in the sputtering yield from the divertor and other plasma facing components is expected with the switch of the main ion [35], assuming all other factors remain the same. Since these simulations do not extend into the plasma edge or SOL regions, a

Figure 23: Example impurity density profile modifications applied to the prescribed integrated model inputs (top row) for separating local effective charge, Z_{eff}, effects (via Be – left column) and local radiated power density, Q_{rad}, effects (via W – right column). The impact of these modifications on the simulated time traces of total radiated power, P_{rad}, (middle row) and the line-integrated effective charge, Z_{eff}, (bottom row) are also provided for demonstrate its validity.

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Figure 24: Example impurity density profile modifications applied to the prescribed integrated model inputs (top row) for separating local effective charge, Z_{eff}, effects (via Be – left column) and local radiated power density, Q_{rad}, effects (via W – right column). The impact of these modifications on the simulated time traces of total radiated power, P_{rad}, (middle row) and the line-integrated effective charge, Z_{eff}, (bottom row) are also provided for demonstrate its validity.

quantitative estimate of the sputtering source and the impurity transport into the plasma core was not self-consistently calculated. As a proxy, the impurity density boundary conditions was varied to study the sensitivity of the core plasma conditions to the increased sputtering yield.

Figure 27 show the effect of an impurity density BC scan on the kinetic profiles of the tritium simulation at $t = 7$ s. The BC scan increases all the impurity BCs, i.e. for Be, Ni, and W simultaneously, by a factor, $\Delta n_{\text{imp},m}/n_{\text{imp}}$, of 1.2 and 1.5. The time traces are not shown as they do not add much value to this discussion.

As expected from the fact that the plasma profiles are not significantly altered, the amount of core impurity accumulation increases as the impurity BC increases, leading to more T_e hollowing. However, as both T_e and T_i are also slightly affected for $\rho_{\text{tor}} \lesssim 0.5$, this core increase is not necessarily linear with the BC. This effect on the temperature profiles between $0.2 \leq \rho_{\text{tor}} \lesssim 0.5$ is likely due to the increasing Z_{eff} with higher impurity content. This can increase the critical threshold of the turbulent transport, depending on the unstable modes present in the scenario, reducing the turbulent transport and allowing the energy conservation equation to support higher gradients. In spite of this, the boundary impurity density interestingly does not seem to have a strong impact on n_e, only minimially increasing it within the inner core.

Similarly to the switch from 100% D to 100% T, the gas injection rate used to modify the core density also has an impact on the impurity sputtering. This is especially true within the divertor region [36]. The increased neutral density causes the plasma to cool faster from collisional and radiative processes on its approach to the divertor plates. This lowers the incident energy of the particles striking the solid surface and lowers the impurity sputtering rate, specifically of W in the case of the JET divertor plates.

In order to capture this, the time-dependent density profiles were again prescribed and a multiplication factor was applied to all of the impurity density BCs simultaneously. As the exact relation between the gas injection rate and the sputtering yield is not quantitatively known, this study assumes a propor-

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Figure 25: Comparison of the electron temperature (left), T_e, ion temperature (center), T_i, and safety factor (right), q, profiles for the local effective charge, Z_{eff}, sensitivity study of JET#97776 at $t = 7$ s. The dashed vertical lines indicates the internal simulation boundary, $\rho_{\text{tor}} = 0.9$, outside which the profiles are prescribed in the simulation. The impact of the Gaussian modification factor is almost negligible but the effect integrated across the whole plasma is noticeable when the uniform multiplication factor is applied.

Figure 26: Comparison of the electron temperature (left), T_e, ion temperature (center), T_i, and safety factor (right), q, profiles for the local radiated power density, Q_{rad}, sensitivity study of JET#97776 at $t = 7$ s. The dashed vertical lines indicates the internal simulation boundary, $\rho_{\text{tor}} = 0.9$, outside which the profiles are prescribed in the simulation. The impact across the entire radial profile is almost negligible but the effect on the central core is roughly equivalent regardless of the uniform multiplication factor or the Gaussian multiplication factor.

Figure 28: Comparison of the electron temperature (leftmost), T_e, ion temperature (left middle), T_i, electron density (right middle), n_e, and safety factor (rightmost), q, profiles at $t = 7$ s for the combined electron and impurity density BC scan of the T extrapolation simulation (blue) based on JET#97776. The dashed vertical lines indicates the internal simulation boundary, $\rho_{\text{tor}} = 0.9$, outside which the profiles are prescribed in the simulation.

Figure 27: Comparison of the electron temperature (leftmost), T_e, ion temperature (left middle), T_i, electron density (right middle), n_e, and safety factor (rightmost), q, profiles at $t = 7$ s for the impurity BC scan of the T extrapolation simulation (blue) based on JET#97776, using a multiplication factor of 1.2 (red) and 1.5 (purple). The dashed vertical lines indicates the internal simulation boundary, $\rho_{\text{tor}} = 0.9$, outside which the profiles are prescribed in the simulation. All impurity BCs were multiplied by the same factor and was done to represent the increased sputtering yield from the SOL due to an increased main ion mass.

tional relation for the impurity density BC modifications in order to this study the sensitivity of this process, given as follows:

$$\frac{\Delta\, n_{\text{imp},c}}{n_{\text{imp}}} = -c_Y \frac{\Delta\, n_e}{n_e} \qquad (2)$$

where $c_Y = 0.5$. For example, this means a 50% increase in the n_e BC results in a 25% decrease in the n_{imp} BCs. While it is unlikely that c_Y is a constant across all density modifications and all impurity species in reality, this assumption is sufficient for this sensitivity exercise.

Figure 28 shows the results of the n_e BC scan on the base simulation using prescribed density profiles. Surprisingly, using the plasma conditions of this scenario, the effect of the combined electron and impurity density BC scan is nearly imperceptible from the results shown in Figure 13, except for a marginally lower T_i.

Figure 29: Comparison of the local ohmic heat source (left), Q_{ohm}, local radiated heat source (center), Q_{rad}, and local electron-ion energy exchange source (right), Q_{ex}, at $t = 7$ s for the combined electron and impurity density BC scan of the T extrapolation simulation (blue) based on JET#97776. The plotted region ends at the internal simulation boundary, $\rho_{\mathrm{tor}} = 0.9$, due to the large non-physical values present beyond this region.

D   Estimated gas injection modification

In order to use these simulations to provide a prediction for the required gas injection, the BC modifications performed to generate Figures 13 and 28 can be permuted to generate a grid of simulation results. Then, since the performance of the high-power phase of the hybrid scenario is dependent on the q profile when the auxiliary heating systems are switched on, its properties at $t = 7$ s can be used to compare the T simulation configuration to the reference D simulation. Specifically, the safety factor at the magnetic axis, q_0, at $t = 7$ s is used as the metric since it is the location with the greatest variability in the set of simulated q profiles.

Table 2 contains the values of q_0 at $t = 7$ s for the various permutations of the density BCs used in this predictive study. The target value, $q_0 = 1.132$, was taken from the validated reference D simulation. It is suspected that the impurity scenarios with $\Delta n_{\mathrm{imp},m}/n_{\mathrm{imp}} \geq 20\%$ and $c_Y = 0.5$ are the most realistic, resulting in the required density increase at the simulation BC to be estimated at 30–40%. However, the experimental deductions were all made relative to a line-integrated density measurement, a more meaningful comparison would be to compute the line-integrated density of the simulation via a synthetic diagnostic. This is because the self-consistent calculation of the density profile evolution within the original n_e BC scan can affect the normalized density gradients such that the whole profile is not just a scaling of the reference simulation.

Figure 30 shows the results of the synthetic diagnostic on the simulation results, as applied to the last three columns of Table 2. From these plots and the estimated range of n_e BC increase given above, the synthetic diagnostic estimates the corresponding line-integrated n_e increase to be between 20 – 30%. This value is slightly lower but agrees well with the empirical estimation from the previous H ramp-up experiments, providing an extra degree of confidence in the predictive capability of the JINTRAC and QLKNN-jetexp-15D coupled model within JET ramp-up scenarios.

E   Inclusion of flux compression via $\mathrm{d}B_0/\mathrm{d}t$

As indicated in Figure 1, the toroidal magnetic field, B_0, is still varying in time for the initial 5 s of the analysis window. This adds an additional complication to the JINTRAC modelling, as the resulting flux compression can both alter the time evolution of the pressure profile and the magnetic geometry. While this impact was assumed to be minor for the conclusions derived in this study, it was deemed prudent to generate a custom build of JINTRAC was developed to include the $\mathrm{d}B_0/\mathrm{d}t$-terms in the current diffusion equation to verify this assumption.

As shown in Figure 31, this flux compression term effectively increases the current diffusion in the earlier phases of the simulation, leading to a lower q_0

Table 2: Results for the safety factor at the magnetic axis, q_0, at $t = 47$ s for each of the boundary condition modification simulations. The highlighted values are the simulations in each column which are closest to the target value from the D reference simulation, $q_0 = 1.132$.

$\Delta\,n_{\mathrm{imp},m}/n_{\mathrm{imp}}$	0%	20%	50%	0%	20%	50%
$\Delta\,n_e/n_e$	$c_Y = 0.0$			$c_Y = 0.5$		
0%	1.383	1.474	1.625	–	–	–
25%	1.171	1.244	1.453	1.137	1.186	1.301
50%	1.046	1.093	1.189	1.012	1.030	1.073
75%	0.989	1.019	1.083	0.949	0.960	0.981
100%	0.959	0.981	1.032	0.903	0.922	0.934

Figure 30: Time trace comparisons of the electron density ratios, $n_e/n_{e,\mathrm{base}}$, for the T mass-induced sputtering, $\Delta n_{\mathrm{imp},m}/n_{\mathrm{imp}}$, BC scan for the values of 0% (left), 20% (center), and 50% (right) based on the JET#97776 D discharge. The dashed vertical lines indicates the reference time at which the auxiliary heating systems are planned to be activated, $t = 47$ s.

at $t = 7$ s and a faster approach of the $q = 1$ to the inversion radius – about 500 ms faster than the reference D simulation. The compression of the flux surfaces increases the core pressure, via both the density and temperature, while also pushing impurities inward and accelerating the core impurity accumulation, but an overall $\sim$10% higher core T_e and $\sim$5% higher core n_e at $t = 7$ s. However, this results from a $\sim$10% increase in Q_{ohm}, a $\sim$25% increase in $|Q_{\mathrm{rad}}|$, as shown in Figure 32. In this scenario, the combined impact of the increased current diffusion and the core pressure increase from the flux surface compression more than compensates for the higher radiative loss due to the likewise compression of the impurity species densities.

As the impact on the predicted profiles is mostly limited to the central core region, the simulated profile sensitivities from the $\mathrm{d}B_0/\mathrm{d}t$-terms do not have any strong impact on the conclusions made by this study. However, its contribution to the overall discrepancy in the central core is not negligible and should be considered an important inclusion in future ramp-up scenario modelling. That said, the likelihood for power reactors to employ a ramp-up scenario containing a time-dependent B_0 remains uncertain.

Figure 31: Comparison of the electron temperature (leftmost), T_e, ion temperature (left middle), T_i, electron density (right middle), n_e, and safety factor (rightmost), q, profiles at $t = 7$ s for the isotopic extrapolation of the D reference simulation based on JET#97776, executed using a 100% fuel ion mix of D (blue), H (red) and T (purple). The dashed vertical lines indicates the internal simulation boundary, $\rho_{\text{tor}} = 0.9$, outside which the profiles are prescribed in the simulation.

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Figure 32: Comparison of the local ohmic heat source (left), Q_{ohm}, local radiated heat source (center), Q_{rad}, and local electron-ion energy exchange source (right), Q_{ex}, at $t = 7$ s for the combined electron and impurity density BC scan of the T extrapolation simulation (blue) based on JET#97776. The plotted region ends at the internal simulation boundary, $\rho_{\mathrm{tor}} = 0.9$, due to the large non-physical values present beyond this region.